

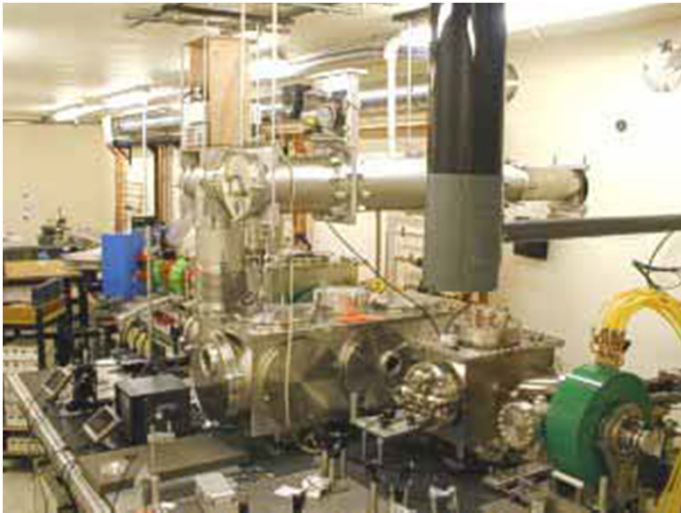
# Emittance and energy spread measurements of relativistic electrons from laser-driven accelerator

## OUTLINE

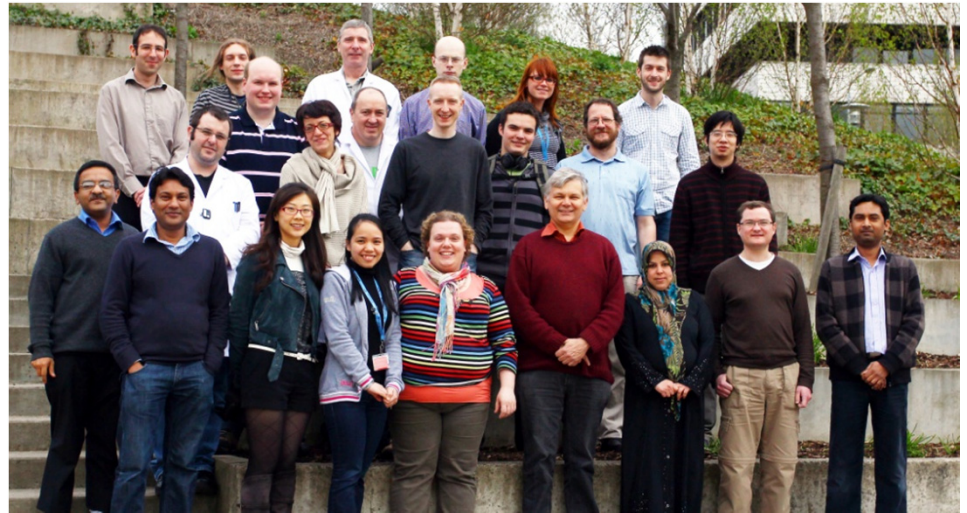
- ALPHA-X Project
- Introduction on laser wakefield accelerator (LWFA)
- LWFA as a light source
- Electron beam characterization: emittance and energy spread

## Advanced Laser Plasma High-energy Accelerator towards X-rays

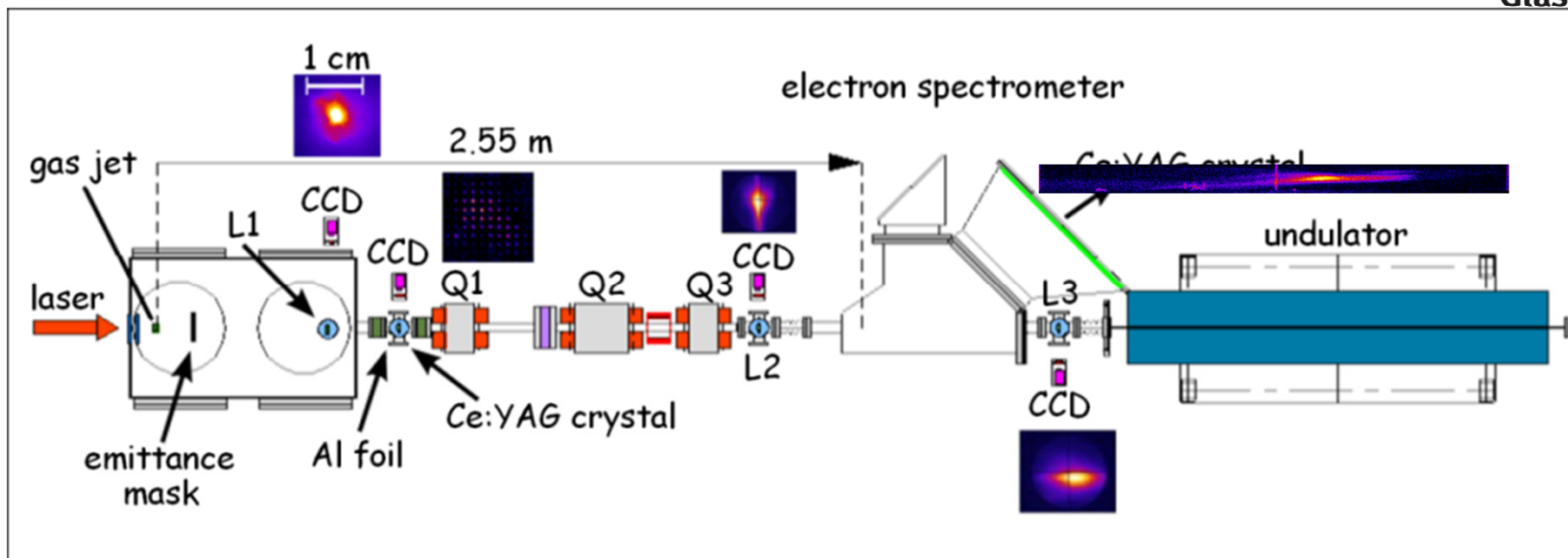
- Collaboration of several universities and research centre in UK (Strathclyde, St. Andrews, Dundee, Abertay-Dundee, Cockcroft Institute)
- Beam line is located at University of Strathclyde (Group Leader: Prof. Dino Jaroszynski)



ALPHA-X beamline (~ 10m)



ALPHA-X Group in Strathclyde



Laser:

900 mJ, 35 fs, 800 nm,  
initial  $a_0 \approx 1.0$

Plasma source:

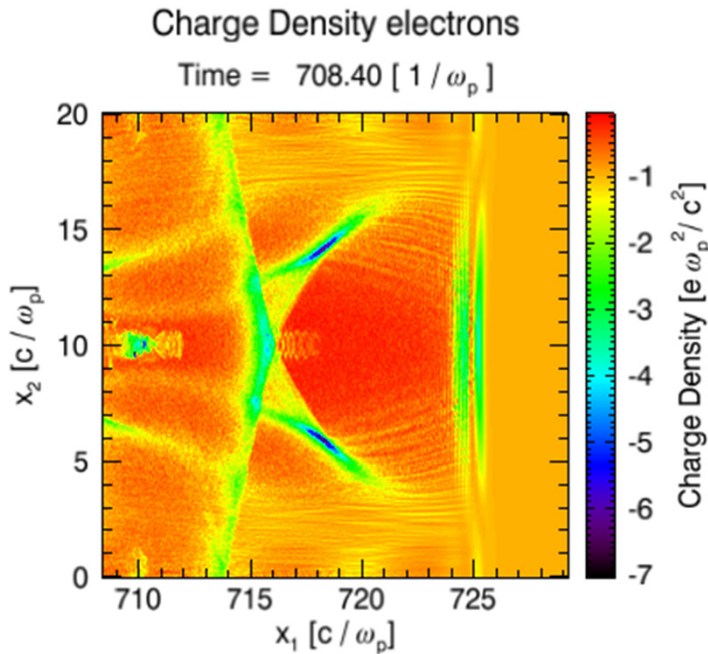
2 mm He gas jet  
 $n_e \approx 10^{19} \text{ cm}^{-3}$



Electron spectrometer

Electron characterization:  
emittance and energy spread

- Uses quadrupole lenses to aid focusing
- Can measure electron energy up to 660 MeV ( $B=1.65 \text{ T}$ )
- energy resolution is  $\sim 0.2\text{-}10\%$  (energy dependent)



- Requires  $I_{laser} \geq 10^{18} \text{ Wcm}^{-2}$  ( $a_0 > 1$  for  $\lambda=1\mu\text{m}$ )
- Relativistic self focusing leads to plasma guiding
- Ponderomotive force leads to charge separation forming the plasma wakefield
- Electrons are self-trapped at the back of the bubble and accelerated to high energy
- Maximum electron's energy at the dephasing length:

Conditions for bubble formation:

$$k_p R \approx 2\sqrt{a_0}$$

$$P > P_{crit} \approx (\tau[\text{fs}] / \lambda[\mu\text{m}])^2 \times 30 \text{ GW}$$

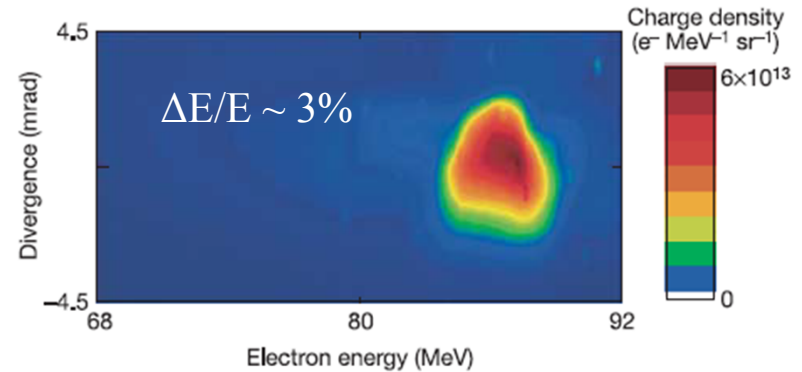
$$\gamma_{max} \cong \frac{2a_0 (\omega_{laser} / \omega_{plasma})^2}{3}$$

# Laser wakefield accelerators (LWFA)

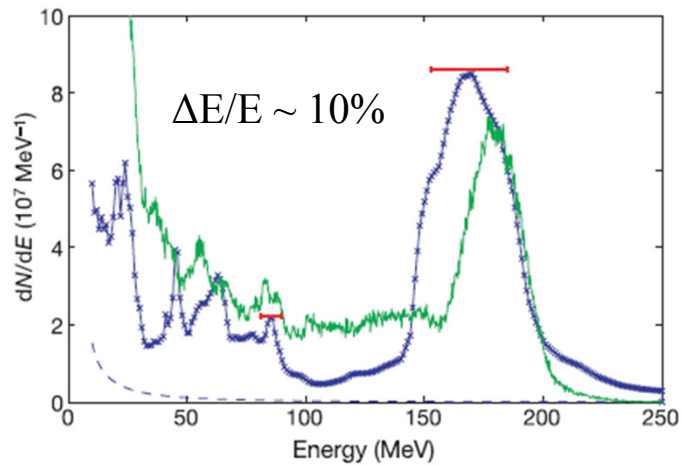


few mm-cm  
acceleration length

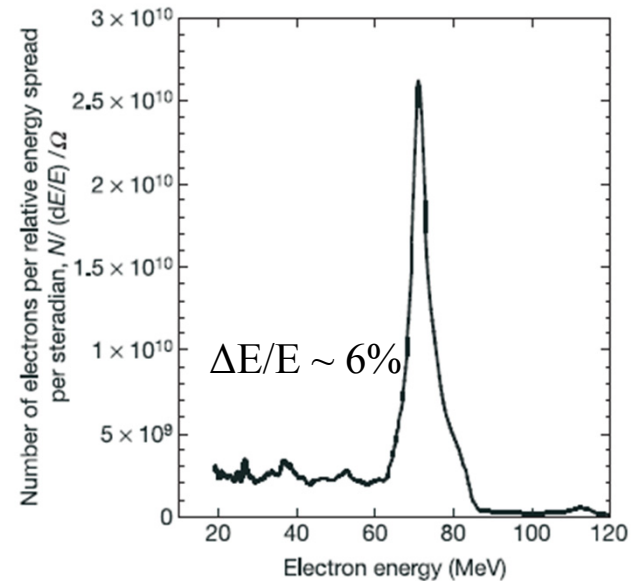
80 MeV – 1GeV  
electron energy



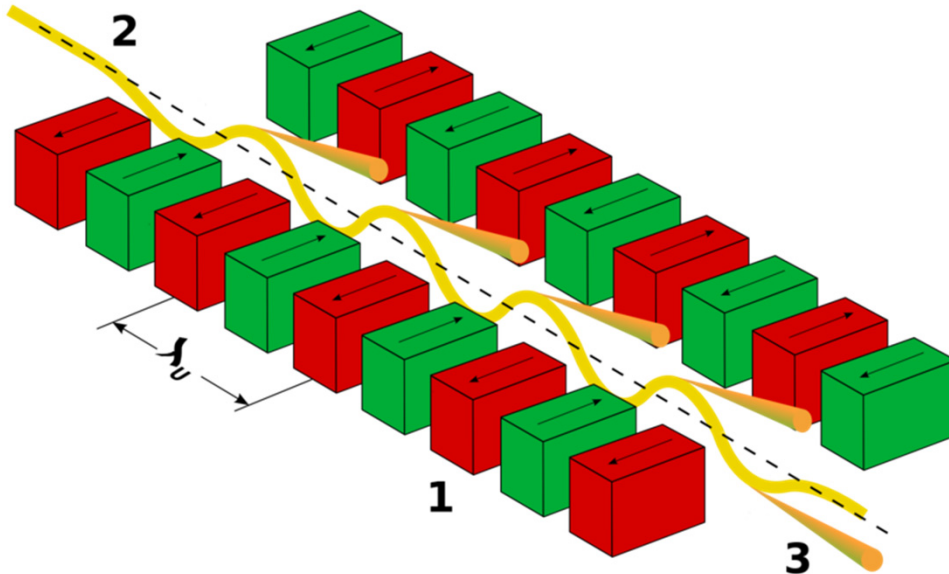
G.R. Geddes et al.



J. Faure et al.



S.P. D. Mangles et al.



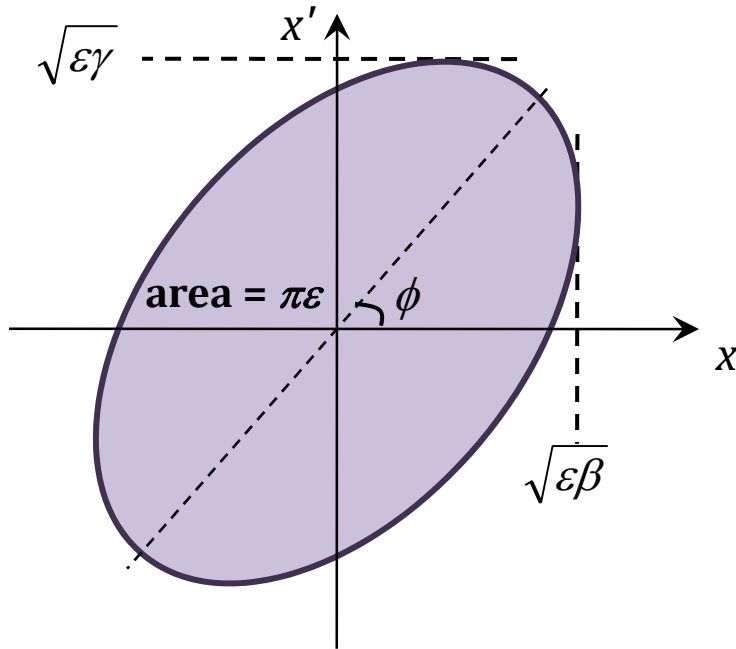
Radiation brightness:

$$B = \frac{I}{4\pi^2 \varepsilon_x \varepsilon_y}$$

Radiation coherence:

$$\varepsilon_{x,y} \leq \frac{\lambda_{rad}}{4\pi}$$

Phase space ellipse

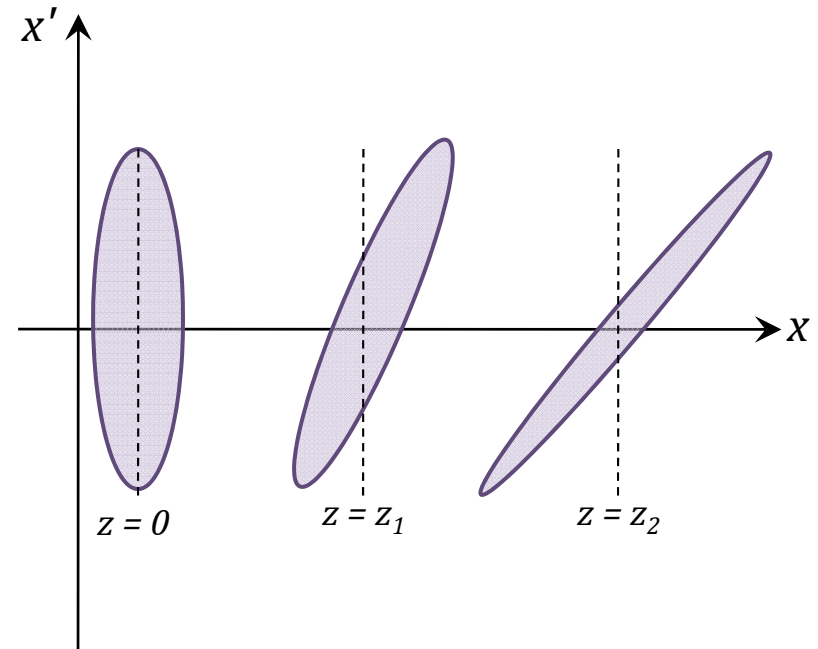


$$\gamma x^2 + 2\alpha x x' + \beta x'^2 = \varepsilon$$

$\alpha, \beta, \gamma, \varepsilon \rightarrow$  Twiss parameters

$$\tan(2\phi) = \frac{2\alpha}{\gamma - \beta} \quad \gamma = \frac{1 + \alpha^2}{\beta}$$

Transformation of phase ellipse for different locations in a drift space



The shape and orientation of the ellipse changes but the area is the same (Liouville Theorem)

Emittance as the area of trace space:

$$\int_{\text{ellipse}} dx dx' = \pi \epsilon$$

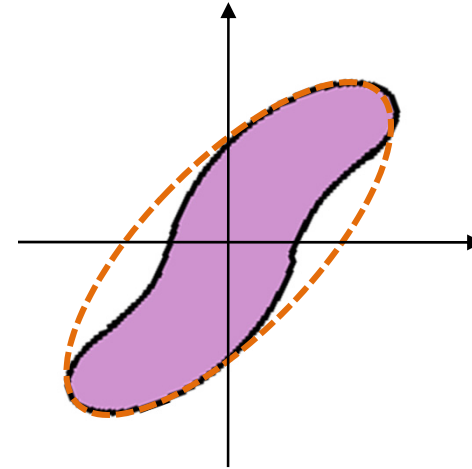
Emittance as a beam quality:

$$\epsilon_{rms} = \sqrt{\langle x^2 \rangle \langle x'^2 \rangle - \langle x x' \rangle^2}$$

rms beam  
width

rms beam  
divergence

correlation  
between x & x'  
(~zero at the  
beam waist)



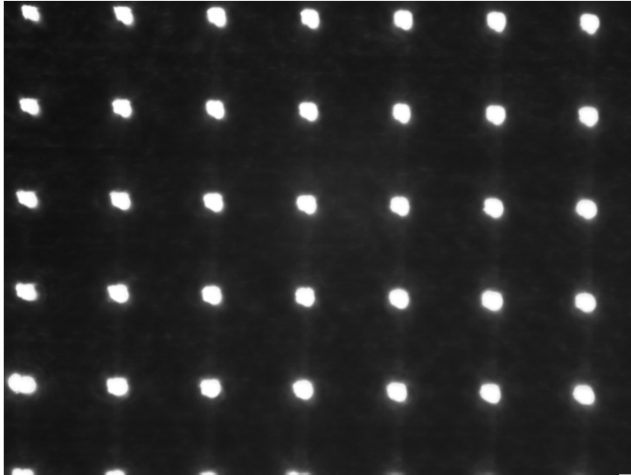
Other definitions of emittance:

$$\epsilon_{norm} = \gamma \beta_z \epsilon$$

where:  $\beta_z = v_z / c$   $\gamma = \frac{1}{\sqrt{1 - \beta_z^2}}$

$$\epsilon_{rms} = 4\epsilon$$

[Lapostelle's  
definition]



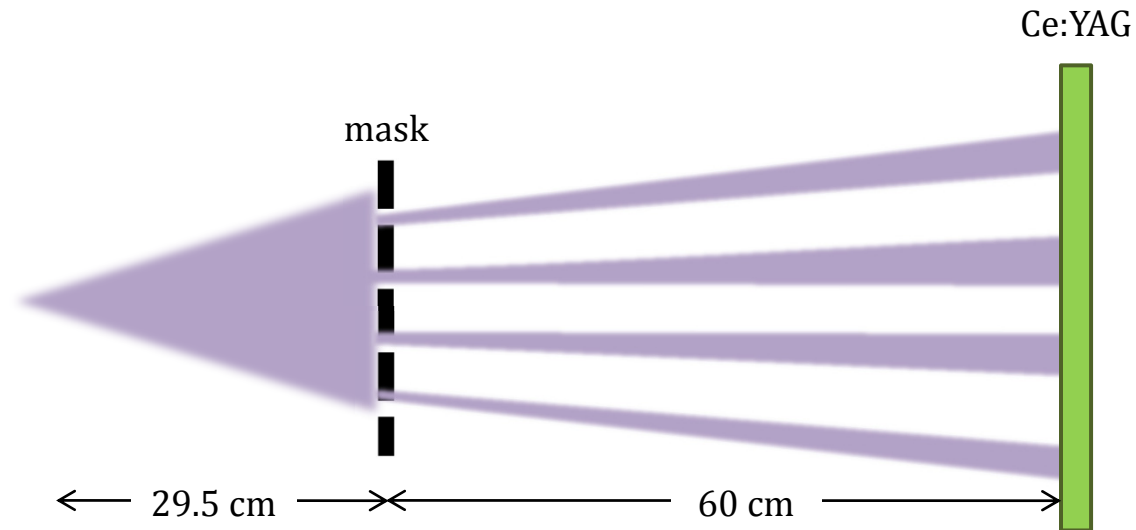
Transmission light image of the tungsten sheet

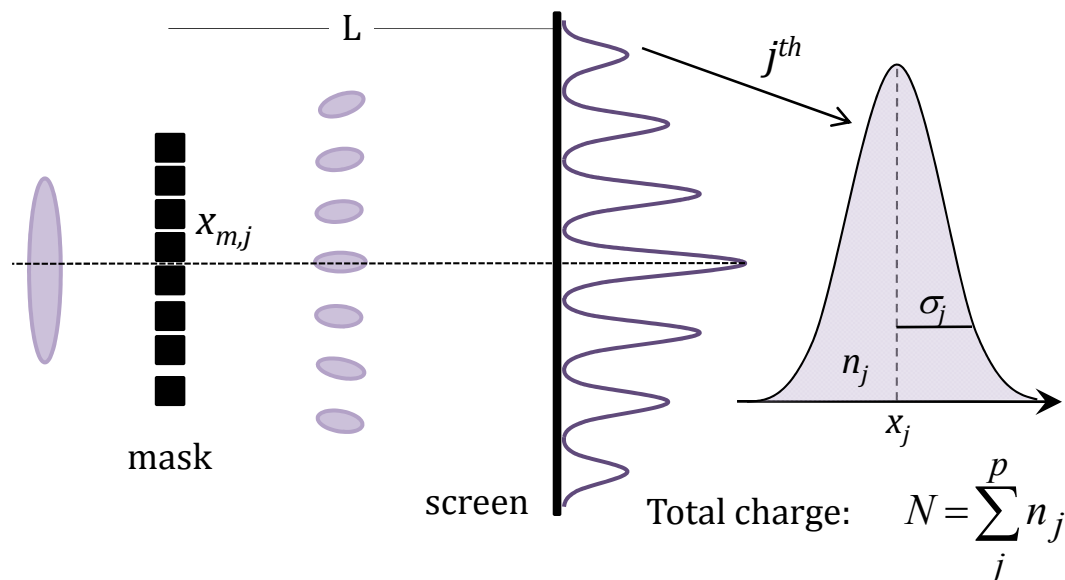
## Pepper-pot mask:

- tungsten sheet
- 27 x 27 grid
- $25 \pm 5 \mu\text{m}$  diameter
- $150 \mu\text{m}$  separation distance

## Detector:

- Ce:YAG screen
- Coupled with a 14-bit CCD camera
- $10 \mu\text{m}$  resolution





In all beamlets:

Mean position:

$$\bar{x} = \sum_j^p n_j x_{m,j}$$

Mean divergence:

$$\bar{x}' = \sum_j^p n_j \bar{x}'_j$$

In  $j^{th}$  beamlet:

Mean divergence:

$$\bar{x}'_j$$

rms divergence:

$$\sigma_{j,x'} = \frac{\sigma_j}{L}$$

$$\varepsilon_{rms} = \sqrt{\langle x^2 \rangle \langle x'^2 \rangle - \langle xx' \rangle^2}$$

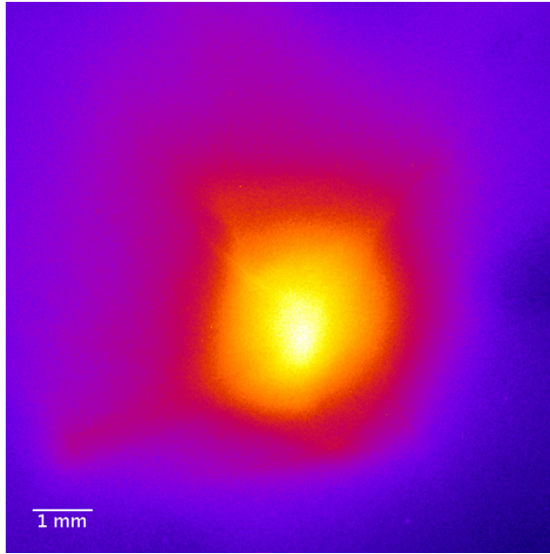
$$\langle x^2 \rangle = \frac{1}{N^2} \sum_j^p n_j (x_{m,j} - \bar{x})^2$$

$$\langle x'^2 \rangle = \frac{1}{N^2} \sum_j^p n_j [\sigma_{j,x'}^2 + (x_{m,j} - \bar{x})^2]$$

$$\langle xx' \rangle^2 = \frac{1}{N^2} \left[ \sum_j^p n_j x_j x'_j - N \bar{x} \bar{x}' \right]^2$$

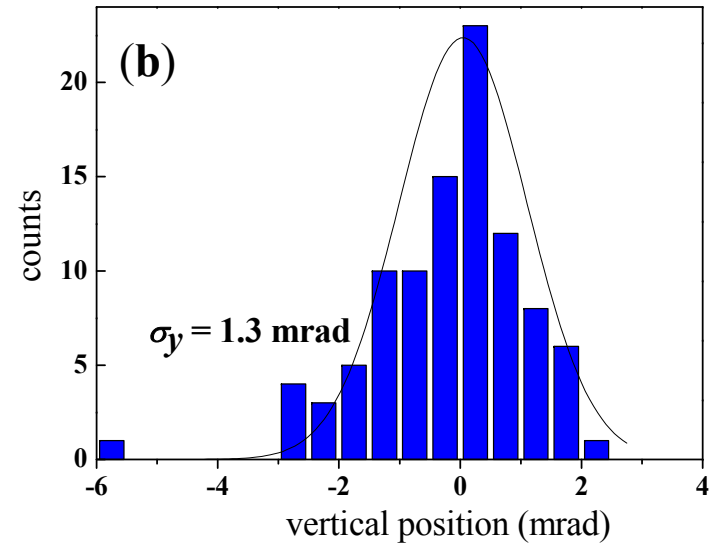
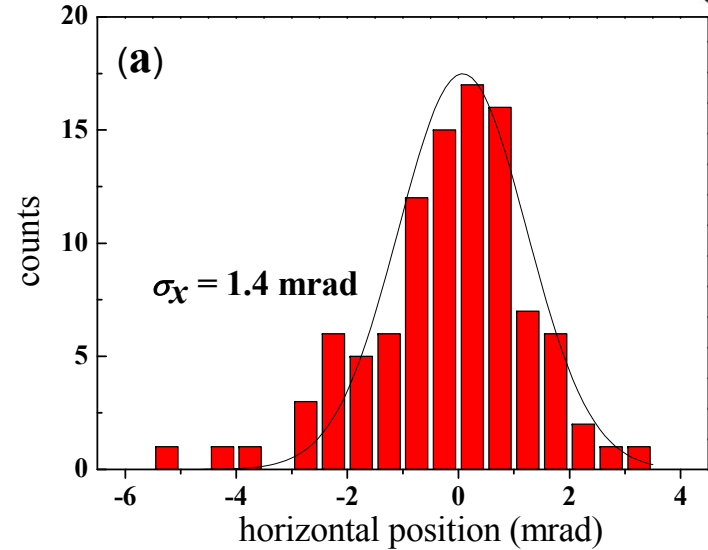
Advantages of pepper-pot method

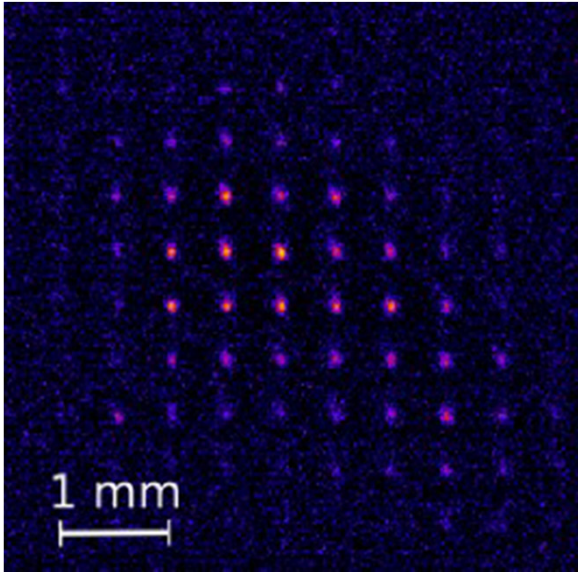
- Single shot measurement
- Less prone to space charge effect
- Can obtain x and y emittance at one measurement



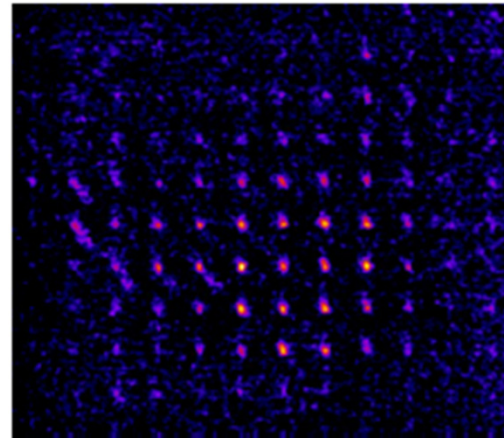
Electron beam as seen in the Ce:YAG screen

The electron beam has a good pointing stability (the deviation is less than one spot size)

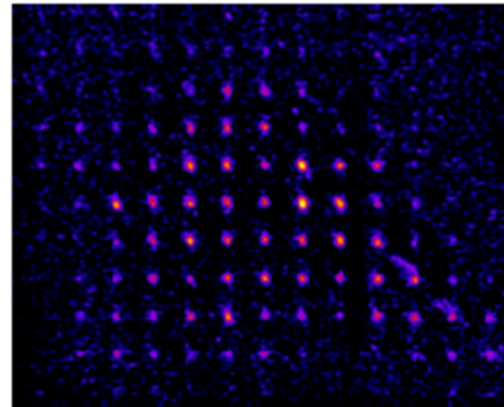




$$\begin{aligned}\varepsilon_{n,x} &= 1.9 \pi\text{-mm-mrad} \\ \varepsilon_{n,y} &= 1.8 \pi\text{-mm-mrad}\end{aligned}$$

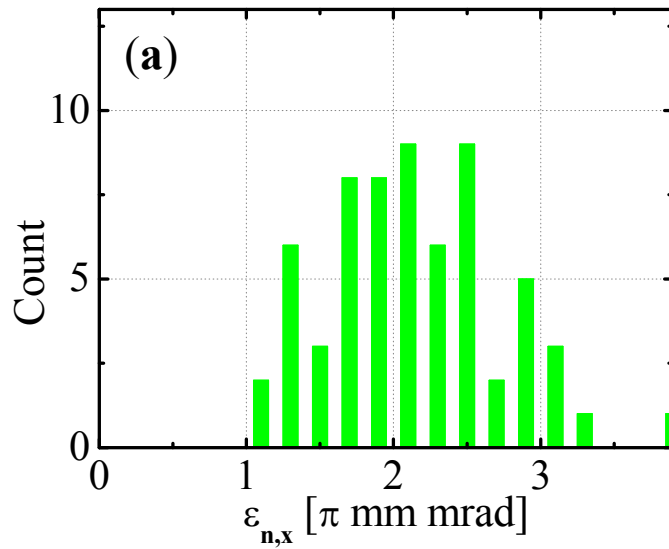


$$\begin{aligned}\varepsilon_{n,x} &= 1.3 \pi\text{-mm-mrad} \\ \varepsilon_{n,y} &= 1.5 \pi\text{-mm-mrad}\end{aligned}$$

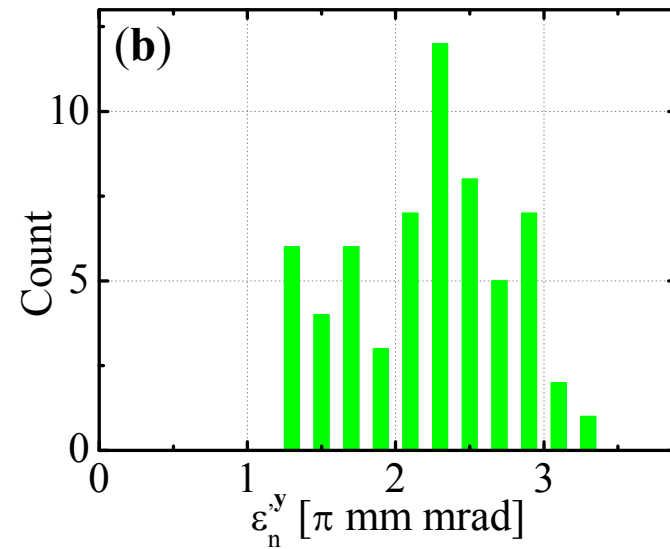


$$\begin{aligned}\varepsilon_{n,x} &= 2.2 \pi\text{-mm-mrad} \\ \varepsilon_{n,y} &= 2.5 \pi\text{-mm-mrad}\end{aligned}$$

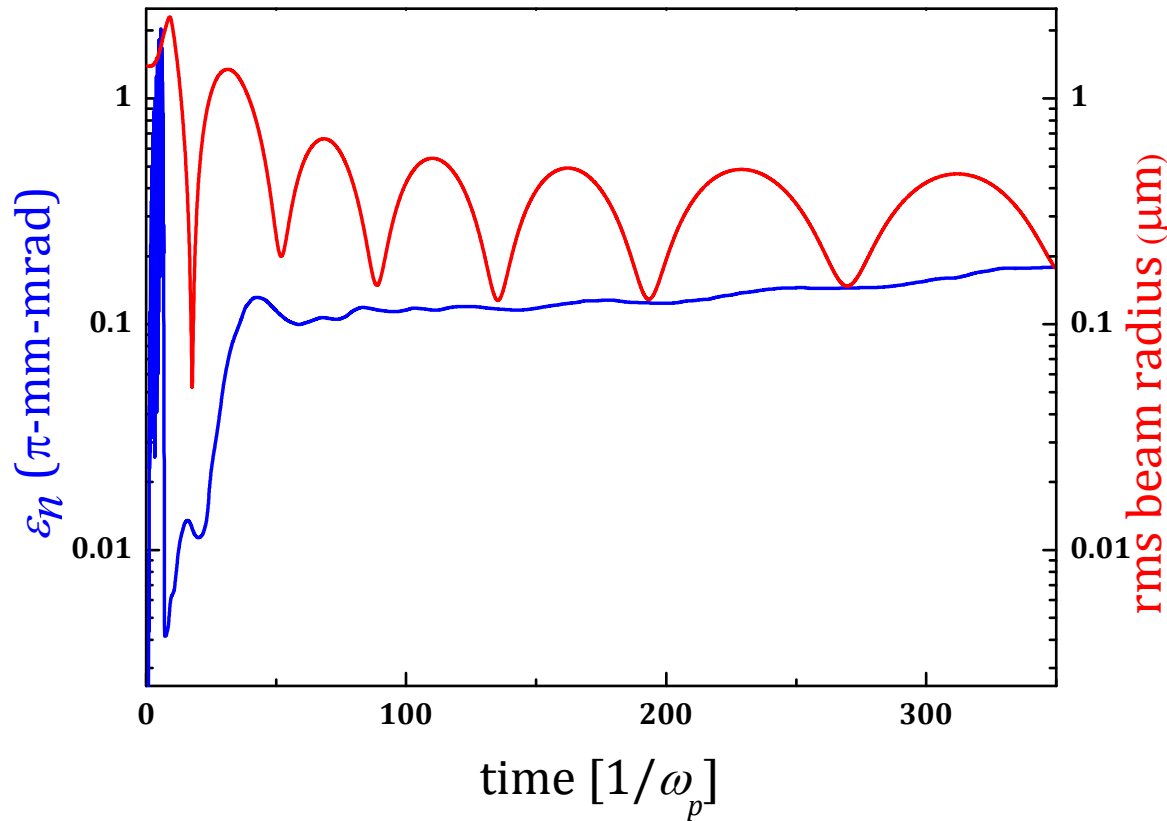
- 64 out of 400 consecutive shots were measured



$2.2 \pm 0.7 \pi\text{-mm-mrad}$

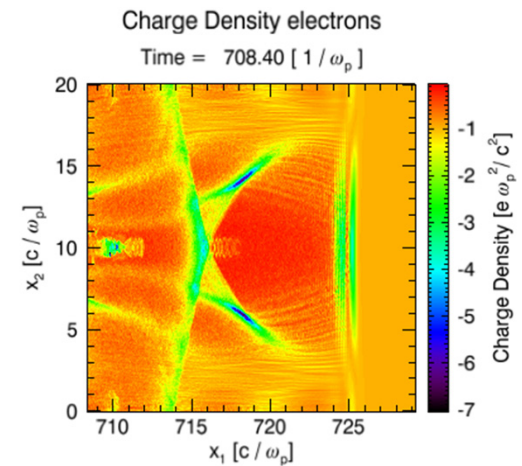


$2.3 \pm 0.6 \pi\text{-mm-mrad}$

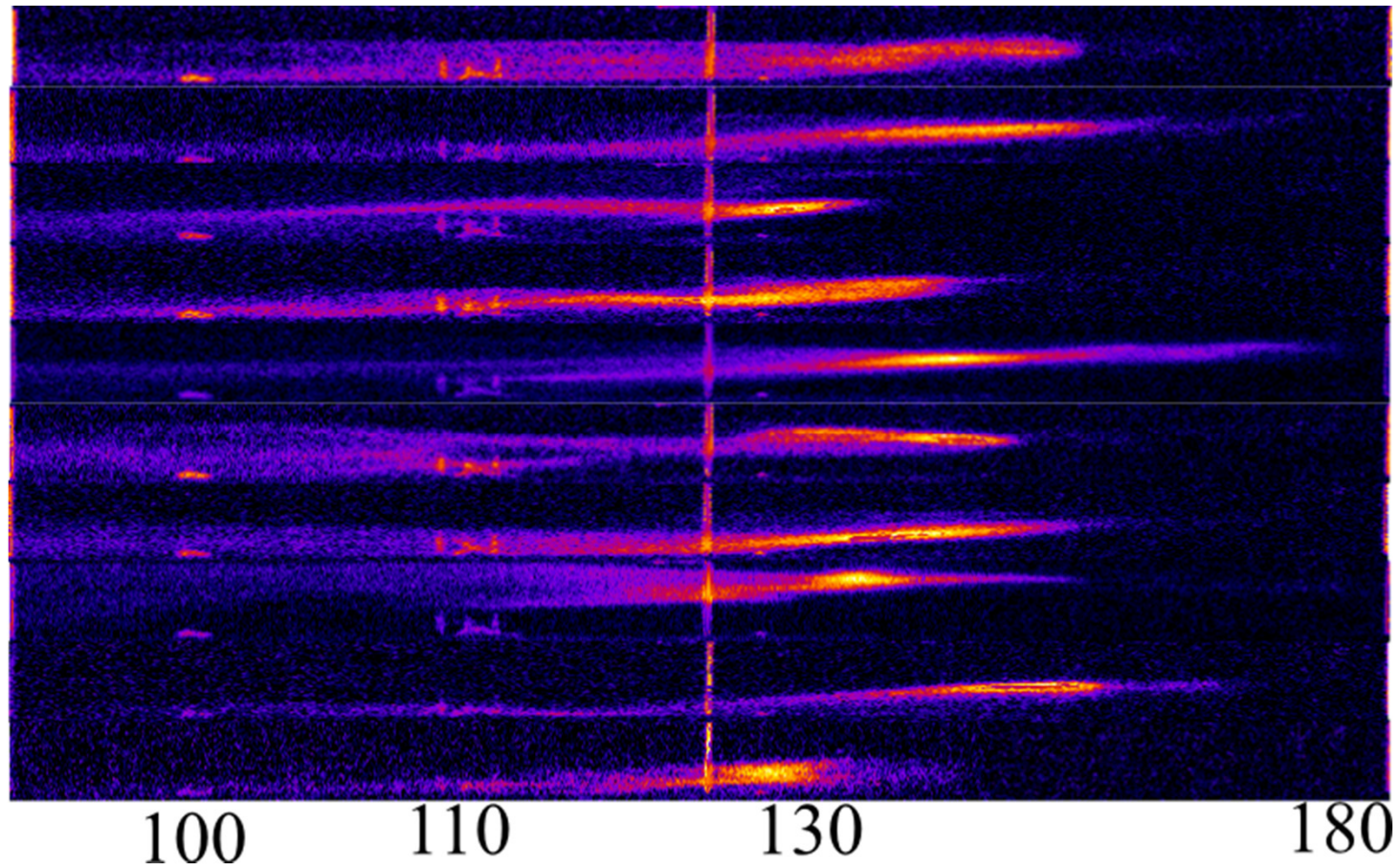


Time evolution of normalised transverse emittance and beam radius of the electron bunch

The emittance growth is dictated by the transverse forces acting on the electron before and during the capture in the plasma 'bubble.'

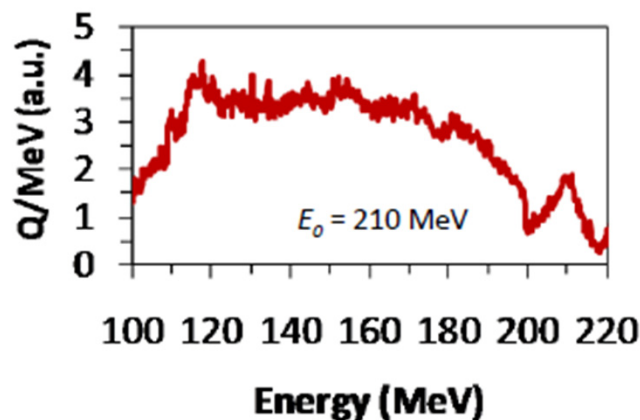
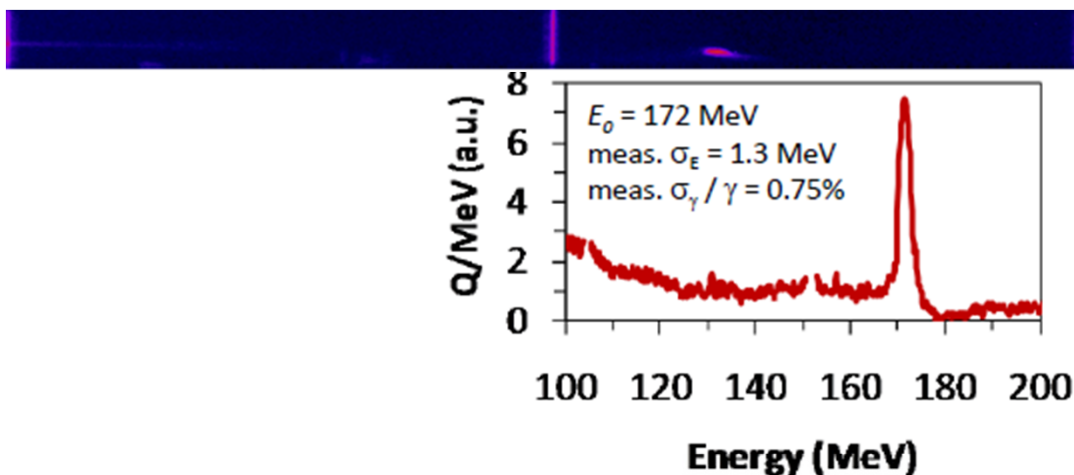
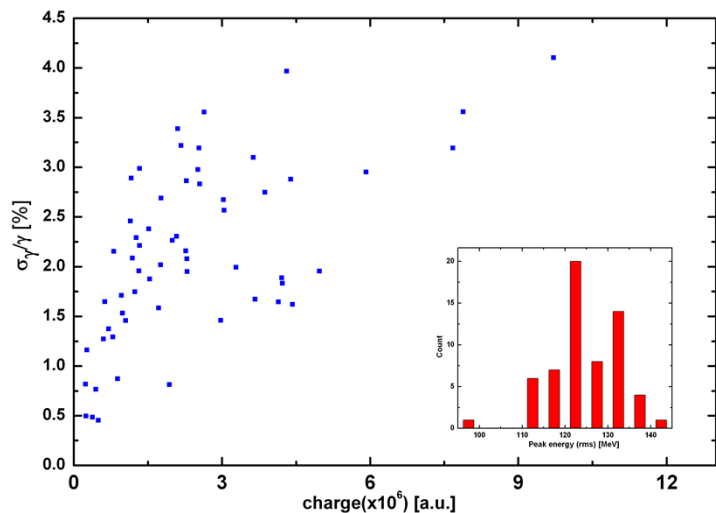
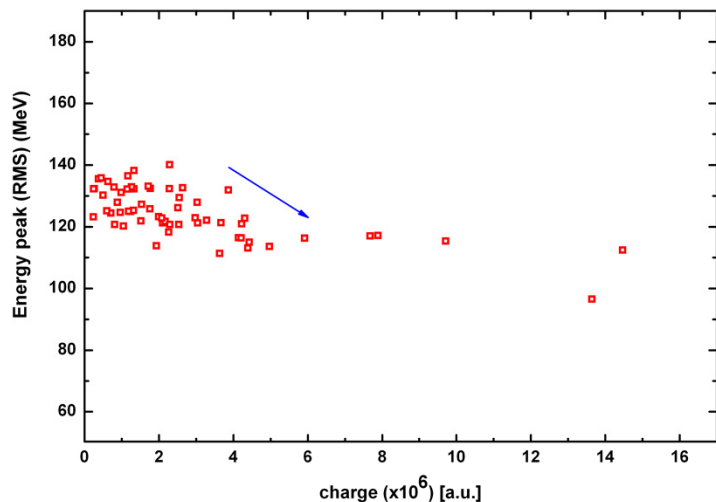


# Electron beam's energy spectra

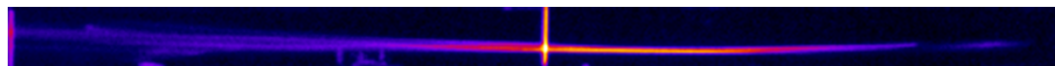


Electron's energy (MeV)

- Beam loading effect
- There seems to be scaling of energy spread and central energy to charge



M. Wiggins et al., *PPCF* 52, 124032 (2010).



- Transverse emittance and energy spread were measured to characterised the electron beam from laser wakefield accelerator.
- The normalised emittance is  $\varepsilon_{n,x,y} = 2.2 \pm 0.7, 2.3 \pm 0.6 \pi\text{-mm-mrad}$ , which is comparable to that of a linear accelerator (measurement was resolution limited).
- Central energy and energy spread shows scaling with the charge.
- The brightness for this accelerator is approximately equal to:

$$B \approx 5 \times 10^{15} \text{ Am}^{-1} \text{ rad}^{-1}$$

which is suitable for driving a compact coherent radiation FEL .