



Energy Recovery Linac Workshop

Longitudinal Stability of ERL with Two Accelerating RF Structures

Getmanov Yaroslav

Budker INP, Russia

9-13 Sep. 2013, Novosibirsk, Russia

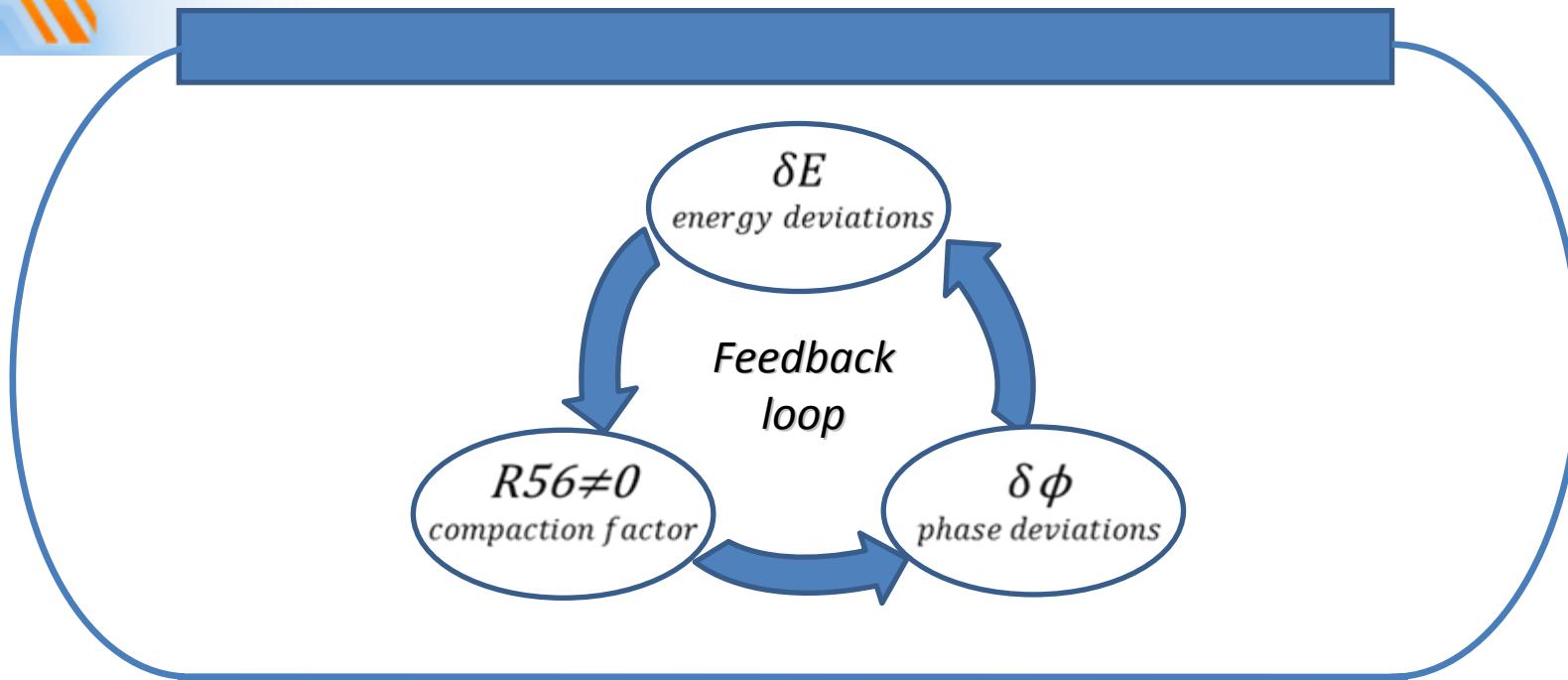


OUTLINE

- Introduction
- MARS scheme
- Theory
- Simulations
- Conclusion



Beam-cavity interaction instabilities



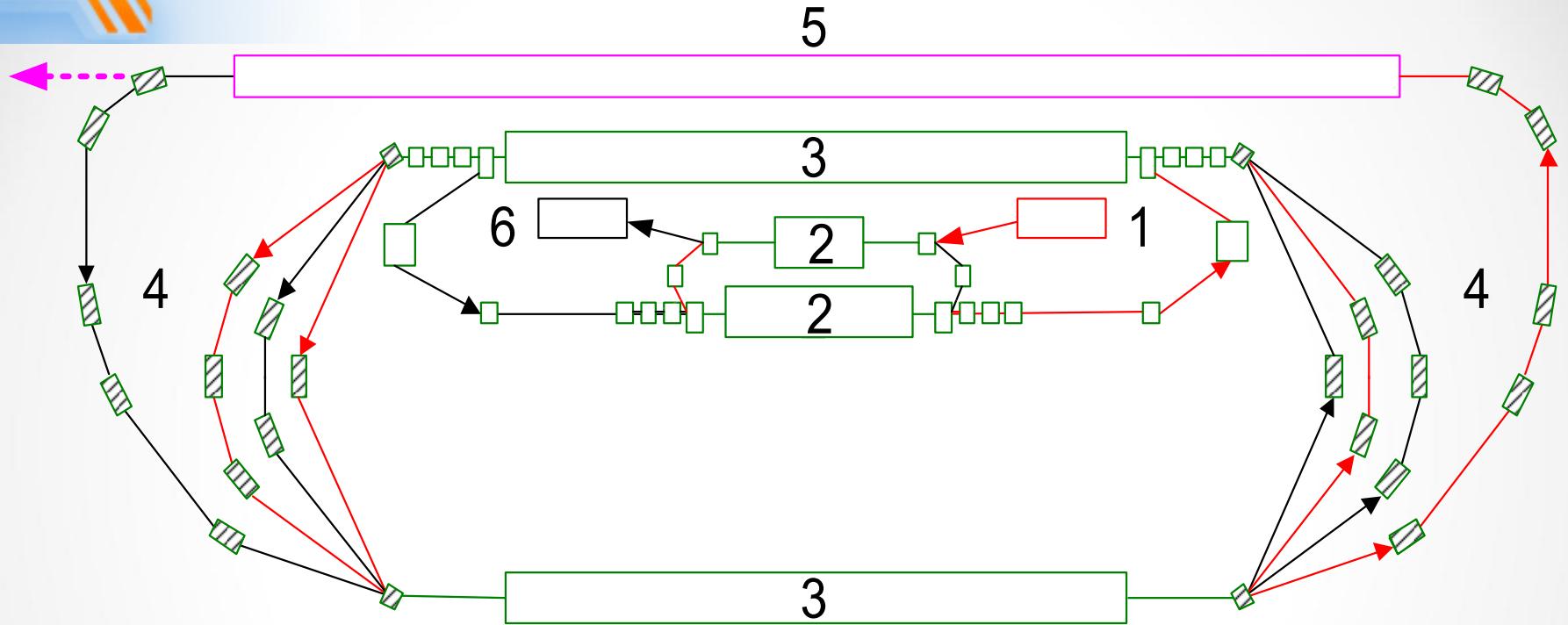
Longitudinal instability
(Interaction with fundamental accelerating mode)

An analytical expression for the threshold current of one-recirculation ERL applicable to all three instabilities can be derived and is given by

$$I_{th} \approx \frac{-2 p_r c}{e(R/Q)_m Q_m k_m M_{ij} \sin(\omega_m t_r + l\pi/2)}$$



ERL with separated linacs



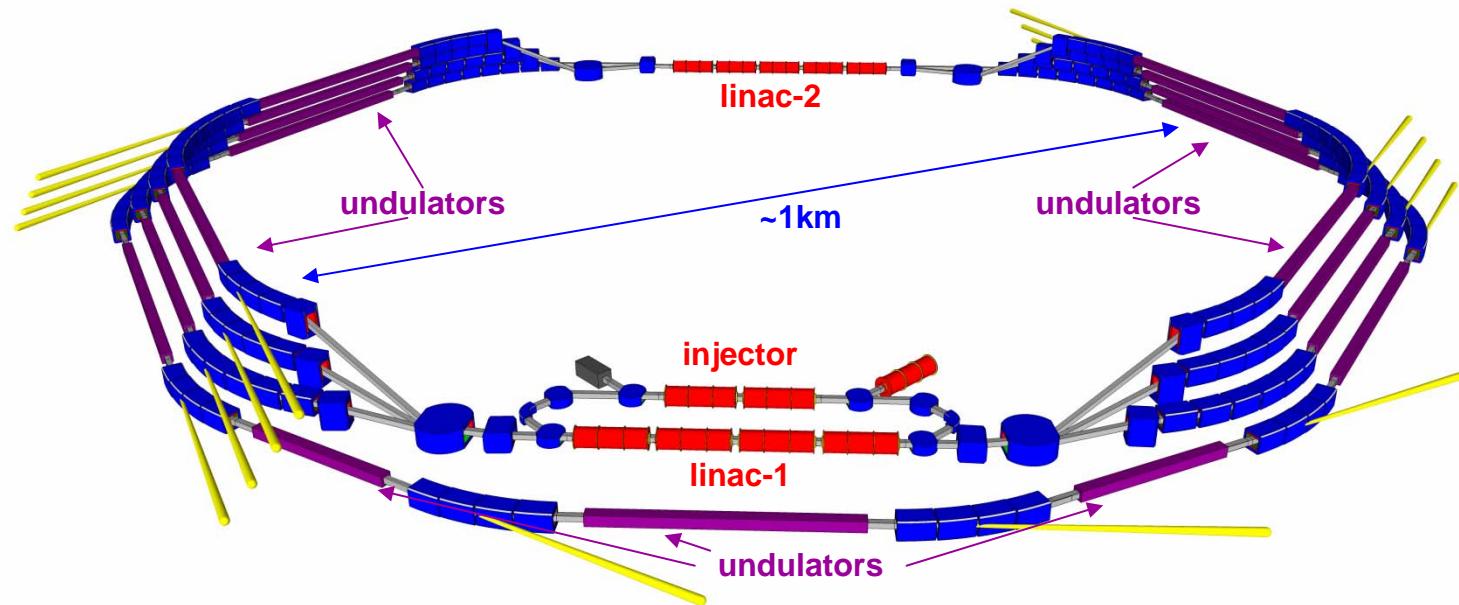
1 – injector, 2 – preliminary accelerating system, 3 – main accelerating RF structure, 4 – magnets, 5 – undulator, 6 – dump.

red arrows – accelerating bunch black arrows – used decelerating bunch



Multipass Accelerator-Recuperator Source project (MARS)

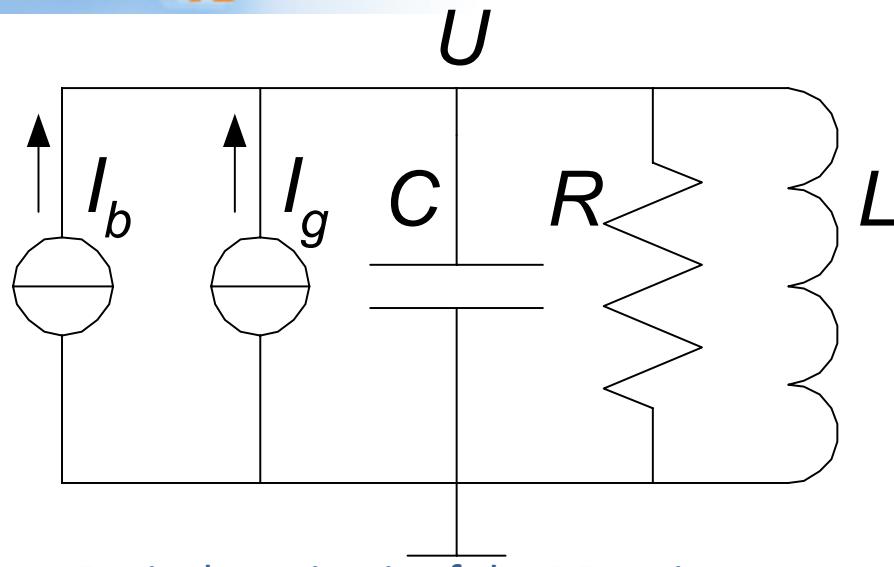
Budker INP



◆ Energy range	5.6, 3.8, 3, 1.2 GeV	◆ Bunch length	0.1 – 1 ps
◆ Average current	10 mA	◆ 7 undulators for	5.6 GeV
◆ Peak current	10 A	◆ 4 undulators for	3.8 GeV
◆ Normalized Emittance	10^{-7} m·rad	◆ 4 undulators for	3 GeV
◆ Bunch charge	10^{-11} C	◆ 4 undulators for	1.2 GeV



Single-cavity model



Equivalent circuit of the RF cavity

Effective voltage of the linac with number α :

$$\operatorname{Re}(U_\alpha e^{-i\omega t})$$

where ω – generator frequency,

U_α – complex amplitude,

$I_{b\alpha}, I_{g\alpha}$ – beam and generator currents

$$\frac{d^2U}{dt^2} + \frac{1}{RC} \frac{dU}{dt} + \frac{1}{LC} U = \frac{1}{C} \frac{d}{dt} (I_b + I_g)$$

$$\frac{2}{\omega} \frac{dU_\alpha}{dt} = \frac{i\xi_\alpha - 1}{Q_\alpha} U_\alpha + \rho_\alpha (I_{b\alpha} + I_{g\alpha})$$

where $\omega_\alpha = \frac{1}{\sqrt{L_\alpha C_\alpha}} = (1 - \frac{\xi_\alpha}{2Q_\alpha})\omega$ resonance frequency

$Q_\alpha = R_\alpha / \sqrt{L_\alpha / C_\alpha} \gg 1$ loaded quality of RF cavity

$$\rho_\alpha = R_\alpha / Q_\alpha = \sqrt{L/C}, R_\alpha$$

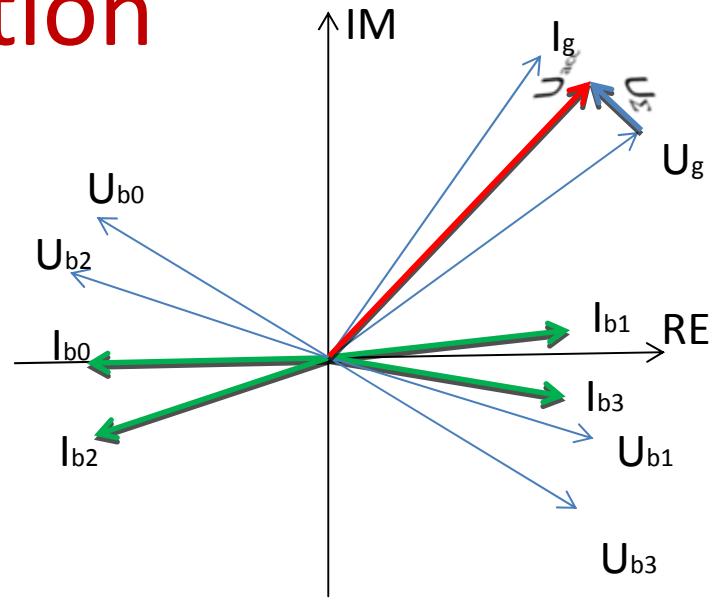
characteristic and the loaded shunt impedances for the fundamental (TM_{010}) mode



Linearization

Linearization near the stationary solution

$$U_{0\alpha} = \frac{R_\alpha}{1 - i\xi_\alpha} [I_{b\alpha}(\mathbf{U}_0) + I_{g\alpha}]$$



complex diagram of linac cavity

where $I_{b\alpha}, I_{g\alpha}$ complex amplitudes of beam and generator currents, gives

$$\frac{2}{\omega} \frac{d\delta U_\alpha}{dt} = \frac{i\xi_\alpha - 1}{Q_\alpha} \delta U_\alpha + \rho_\alpha \sum_\beta \left(\frac{\partial I_{b\alpha}}{\partial \operatorname{Re} U_\beta} \operatorname{Re} \delta U_\beta + \frac{\partial I_{b\alpha}}{\partial \operatorname{Im} U_\beta} \operatorname{Im} \delta U_\beta \right)$$

$$\lambda \delta \mathbf{U} = \mathbf{M} \delta \mathbf{U}$$



Conductivity matrix of ERL with 2 linacs

Matrix of conductivity

$$\mathbf{M} = \begin{pmatrix} -\frac{1}{Q_1} + \rho_1 \frac{\partial \operatorname{Re} I_{b1}}{\partial \operatorname{Re} U_1} & -\frac{\xi_1}{Q_1} + \rho_1 \frac{\partial \operatorname{Re} I_{b1}}{\partial \operatorname{Im} U_1} & \rho_1 \frac{\partial \operatorname{Re} I_{b1}}{\partial \operatorname{Re} U_2} & \rho_1 \frac{\partial \operatorname{Re} I_{b1}}{\partial \operatorname{Im} U_2} \\ \frac{\xi_1}{Q_1} + \rho_1 \frac{\partial \operatorname{Im} I_{b1}}{\partial \operatorname{Re} U_1} & -\frac{1}{Q_1} + \rho_1 \frac{\partial \operatorname{Im} I_{b1}}{\partial \operatorname{Im} U_1} & \rho_1 \frac{\partial \operatorname{Im} I_{b1}}{\partial \operatorname{Re} U_2} & \rho_1 \frac{\partial \operatorname{Im} I_{b1}}{\partial \operatorname{Im} U_2} \\ \rho_2 \frac{\partial \operatorname{Re} I_{b2}}{\partial \operatorname{Re} U_1} & \rho_2 \frac{\partial \operatorname{Re} I_{b2}}{\partial \operatorname{Im} U_1} & -\frac{1}{Q_2} + \rho_2 \frac{\partial \operatorname{Re} I_{b2}}{\partial \operatorname{Re} U_2} & -\frac{\xi_2}{Q_2} + \rho_2 \frac{\partial \operatorname{Re} I_{b2}}{\partial \operatorname{Im} U_2} \\ \rho_2 \frac{\partial \operatorname{Im} I_{b2}}{\partial \operatorname{Re} U_1} & \rho_2 \frac{\partial \operatorname{Im} I_{b2}}{\partial \operatorname{Im} U_1} & \frac{\xi_2}{Q_2} + \rho_2 \frac{\partial \operatorname{Im} I_{b2}}{\partial \operatorname{Re} U_2} & -\frac{1}{Q_2} + \rho_2 \frac{\partial \operatorname{Im} I_{b2}}{\partial \operatorname{Im} U_2} \end{pmatrix}$$

the characteristic equation

$$\lambda^4 - S_1 \lambda^3 + S_2 \lambda^2 - S_3 \lambda + S_4 = 0$$

Stable conditions are given by Lienard-Chipard criterion (4 equations)



Threshold current

Additional constraints

Energy recuperation

$$\begin{cases} \operatorname{Re} \left[U_{01} \sum_{n=0}^{N-1} (e^{-i\varphi_{2n}} + e^{-i\varphi_{4N-2n-1}}) \right] = 0 \\ \operatorname{Re} \left[U_{02} \sum_{n=0}^{N-1} (e^{-i\varphi_{2n+1}} + e^{-i\varphi_{4N-2n-2}}) \right] = 0 \end{cases}$$

Longitudinal focusing

$$e \operatorname{Im} [U_{0\alpha(n)} e^{-\varphi_n}] < 0$$

Assume that equilibrium phases are equal during acceleration. In this case

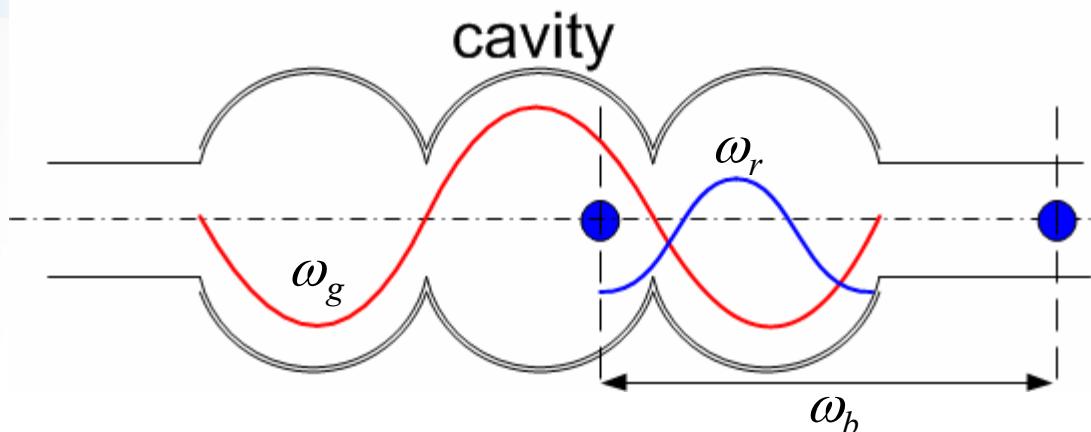
$$\varphi_{2n} - \arg(eU_{01}) = \Phi_1, \varphi_{2n+1} - \arg(eU_{02}) = \Phi_2$$

If the Liénard-Chipart criterion is satisfied, the threshold current is

$$I_{th} = \frac{1/Q_1 + 1/Q_2}{e\rho_1 \sin(2\Phi_1) \sum_{n=0}^{N-1} \sum_{k=0}^{N-1} S_{4N-2n-1, 2k} + e\rho_2 \sin(2\Phi_2) \sum_{n=0}^{N-1} \sum_{k=0}^{N-1} S_{4N-2n-2, 2k+1}}$$



Simulations: wakes



There are three different frequencies in system: generator, gun and cavity resonance

Additional energy gain by electron $\Delta E = -eV_{||}(t)$ where induced voltage $V(t)$

for bunch distribution $\lambda(t)$ is $V_{||}(\tau) = \int_0^{+\infty} G_{||}(t)\lambda(t-\tau)dt$
 $G(t)$ is the wake function

$$G_{||}(t) \cong \frac{\omega_r R_s}{Q} \cos(\omega_r t) \exp\left(-\frac{\omega_r}{2Q} t\right) H(t)$$

For point bunches with charge q , the induced voltage is

$$V_{||}(t) = q \frac{\omega_r R_s}{Q} \cos(\omega_r t) \exp\left(-\frac{\omega_r}{2Q} t\right) H(t)$$



Phases of the induced fields

Phase of induced field by bunch 2
at the pass bunch 1.

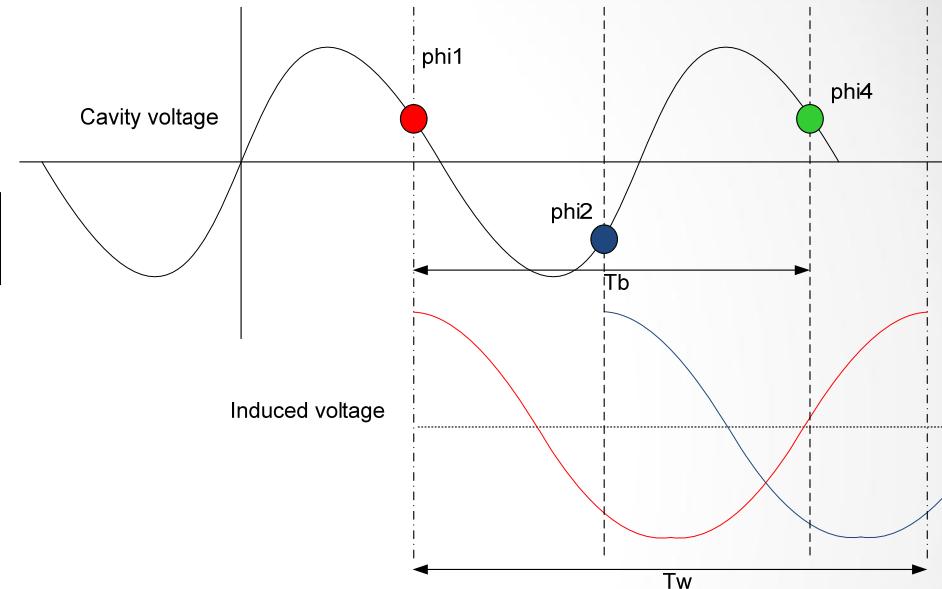
$$\phi = \frac{\omega_r}{\omega_b} \left[(\phi_{1ref} - \phi_{2ref}) + (\delta\phi_{1N_1} - \delta\phi_{2N_2}) \right]$$

Additional energy gain by electron

$$\begin{aligned}\Delta E_{n+1} &= -W_n(t_{n+1}) i^*(t_{n+1}) T = \\ &- W_n(t_n) \exp\left(-\frac{\omega_r}{2Qc} \Delta L\right) \exp\left(i \frac{\omega_r}{\omega_b} 2\pi\right) \exp(-i\phi)\end{aligned}$$

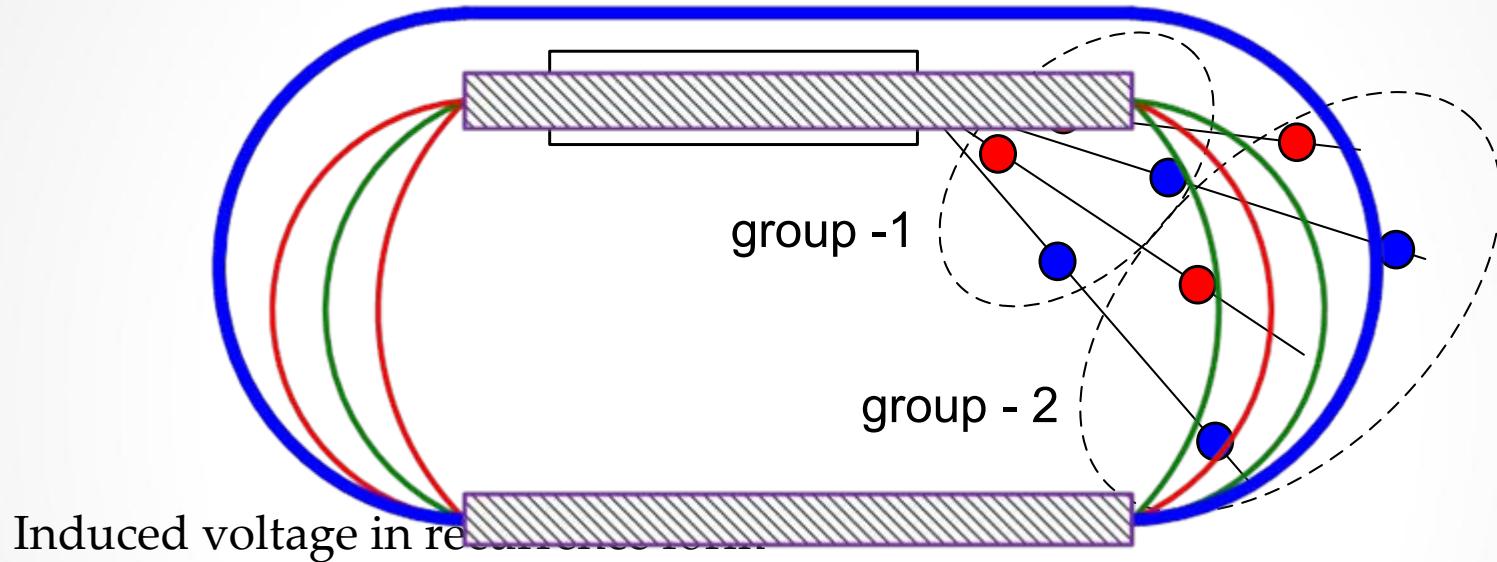
where

$$W_n(\phi_1) = q \frac{\omega_r R_s}{Q} \exp\left(-i \frac{\omega_r}{\omega_b} (\phi_{2ref} + \delta\phi_2)\right) \quad i(\phi_1) = \frac{e}{T} \exp\left(-i \frac{\omega_r}{\omega_b} (\phi_{1ref} + \delta\phi_1)\right)$$





simulations



$$W_n = W_{n-1} \exp\left(-\frac{\omega_r}{2Qc} \Delta L\right) \exp\left(i \frac{\omega_r}{\omega_b} 2\pi\right) + q \frac{\omega_r R_s}{Q} \sum_{m=0}^{Nb-1} \exp\left(-i \frac{\omega_r}{\omega_b} (\phi_{mk} + \delta\phi_{mk})\right)$$

energy gain

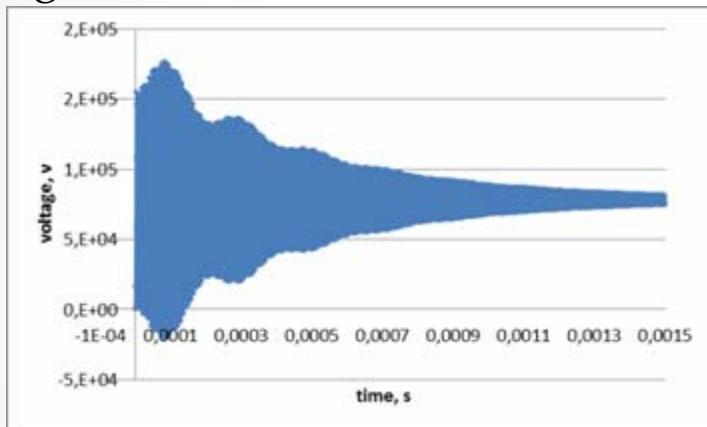
$$\Delta E_{n+1} = -W_n \exp\left(-\frac{\omega_r}{2Qc} \Delta L\right) \exp\left(i \frac{\omega_r}{\omega_b} 2\pi\right) i^* T$$



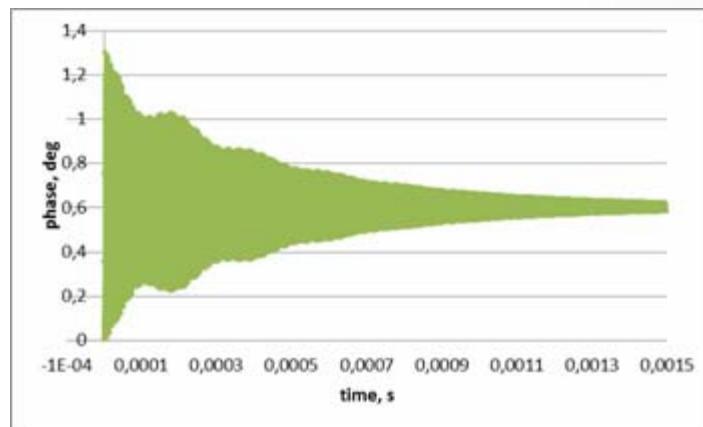
Examples

stable operation

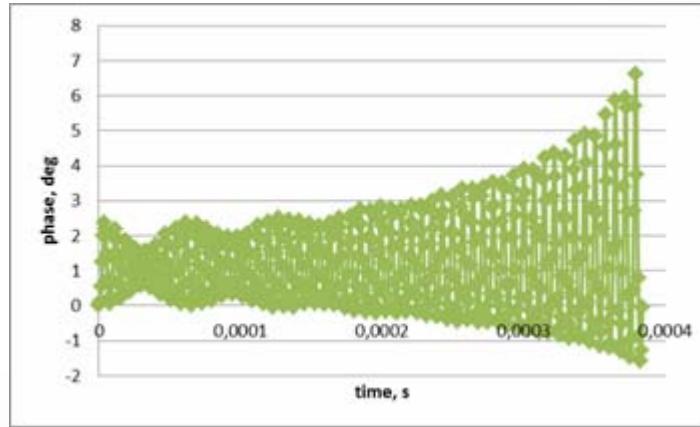
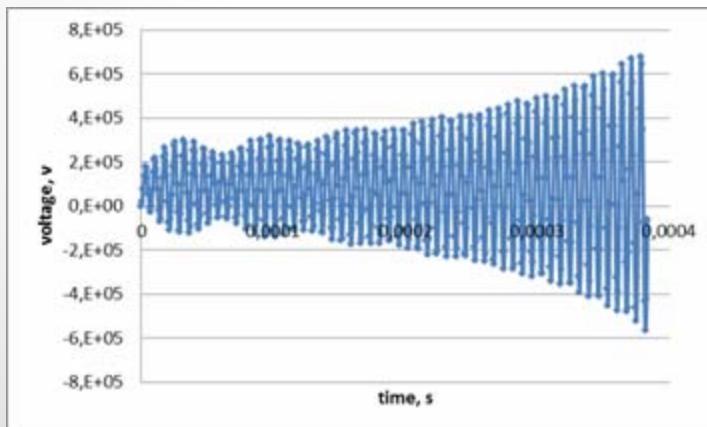
Voltage to time on 2-d linac



Phase to time on 2-st linac

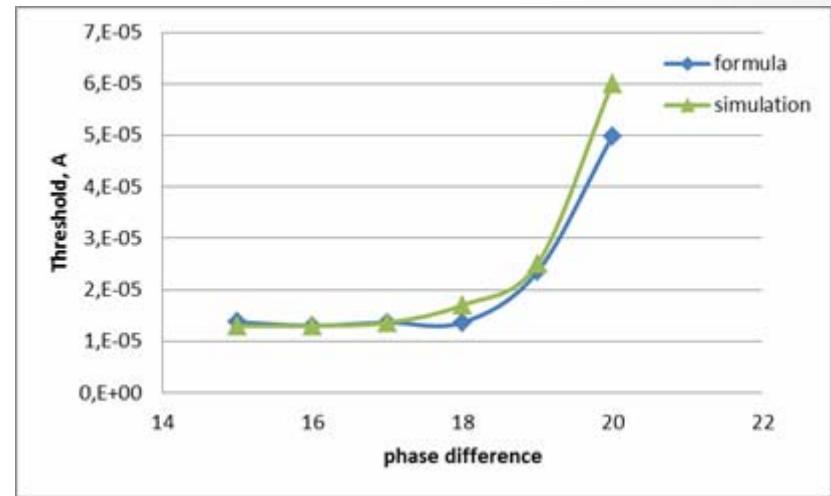
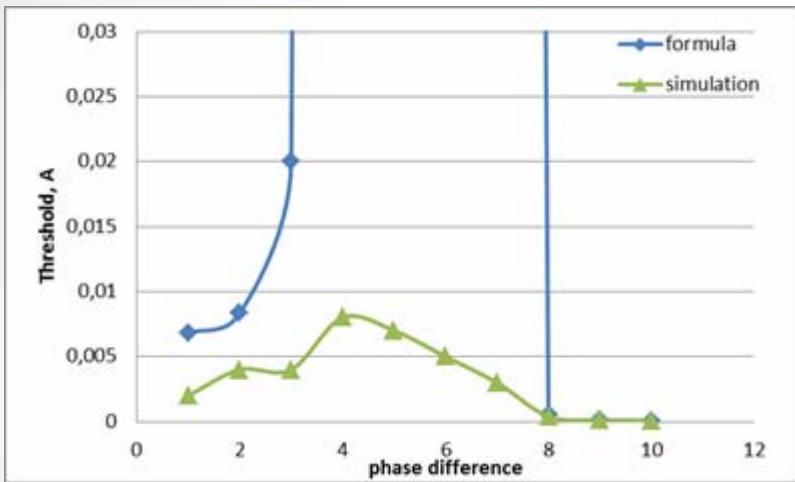


unstable operation





Threshold current



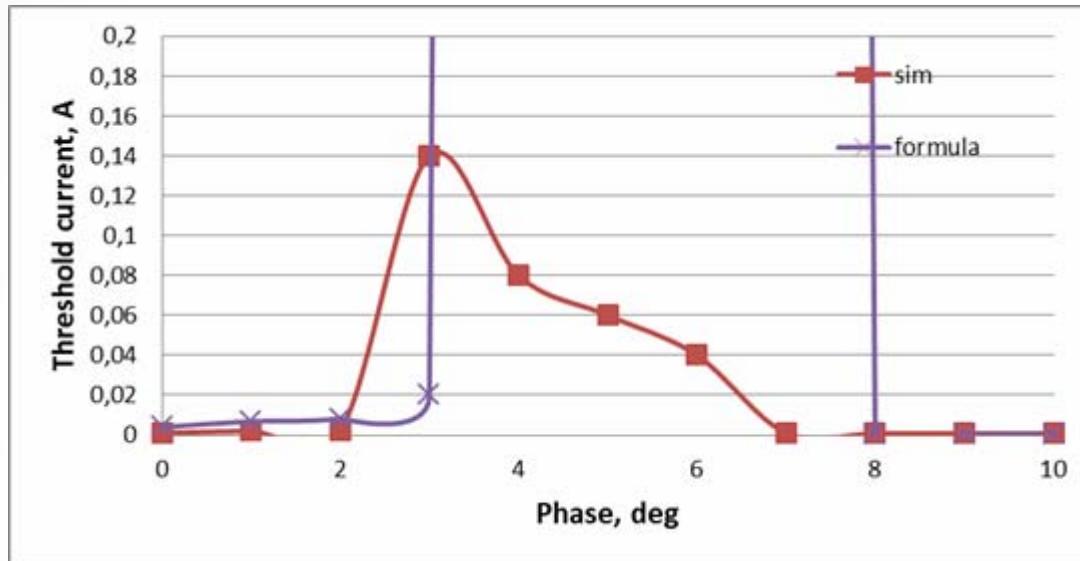
The threshold current dependence on phases with equal detunings of cavities $\xi_{1,2} = 100$

Accelerating phases are equal:

$$\varphi_{2n} - \arg(eU_{01}) = \Phi_1 = \varphi_{2n+1} - \arg(eU_{02}) = \Phi_2$$



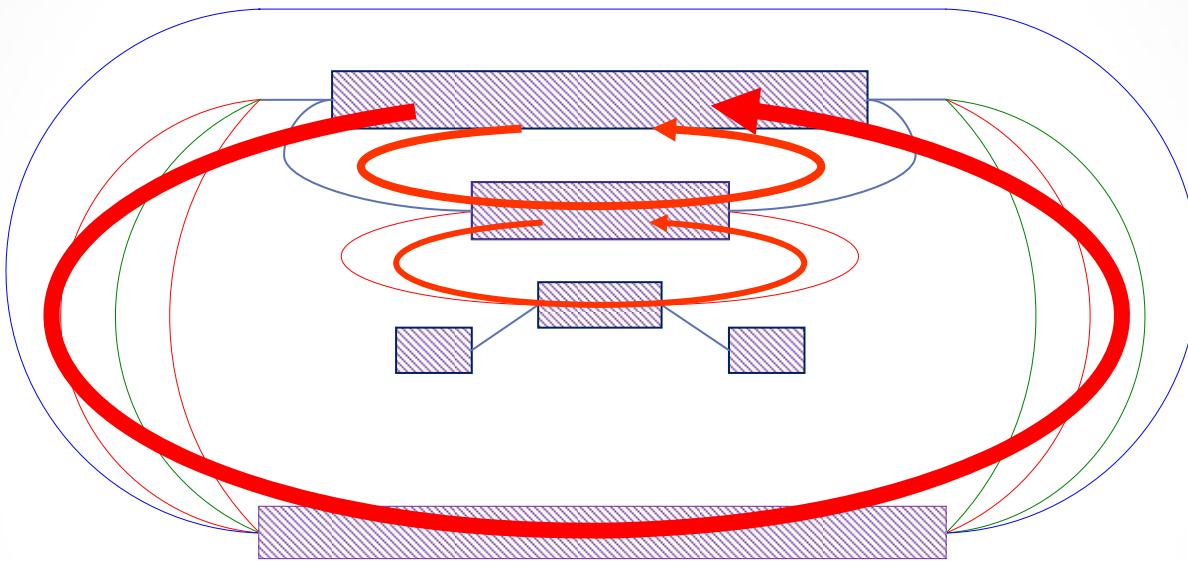
Threshold without preinjection



The threshold current dependence on phases with equal detunings of cavities $\xi_{1,2} = 1000$



Preinjection

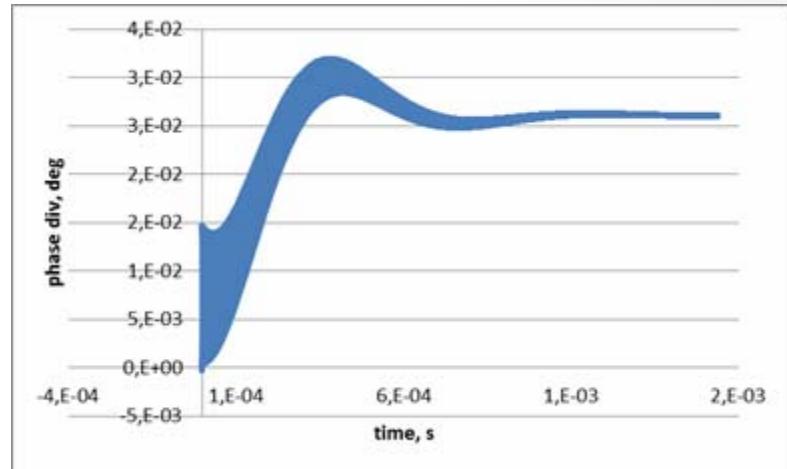
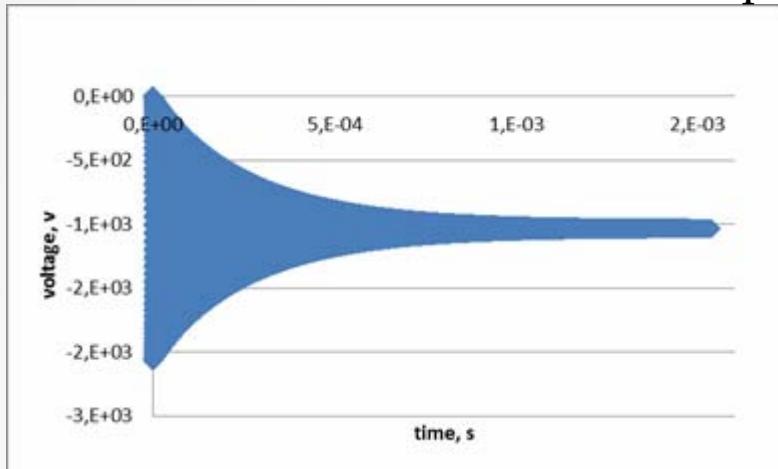


There are two additional feedback loops at the facility with preinjection.



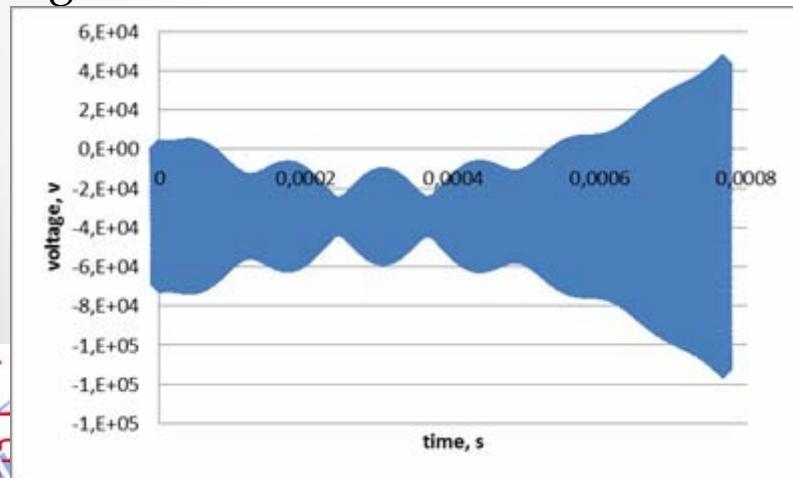
Examples

Stable operation

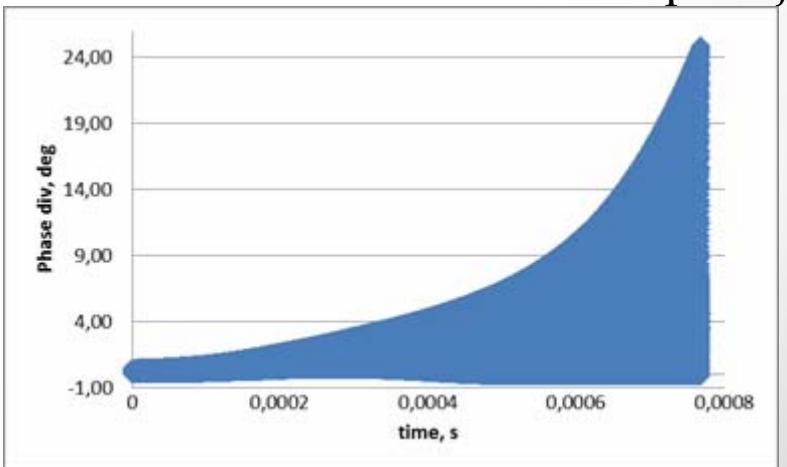


Unstable operation

Voltage to time at the main linac

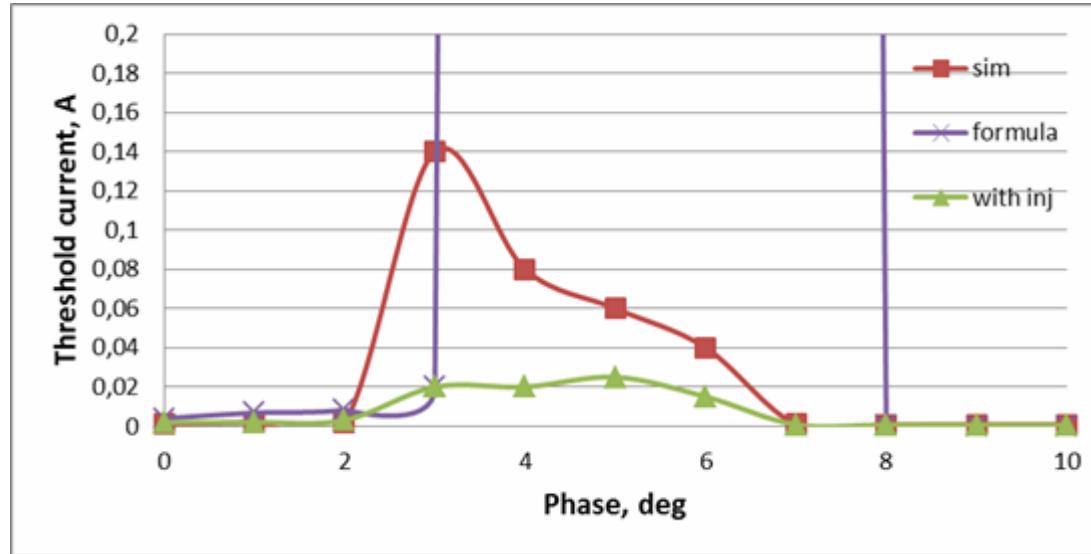


Phase deviation to time at the preinj





Threshold current with injection



The threshold current dependence on phases with equal detunings of cavities $\xi_{1,2} = 1000$



Conclusion

- The criterion of the longitudinal stability for the ERL with two accelerating structures is derived
- Numerical simulations were made for the project of the light source based on multiturn ERLs.
- The simulated threshold current is lower than the theoretical lower limit of the threshold current.
- To increase the threshold current, it necessary to develop a proper feedback system.



Thank you for your attention!



Liénard-Chipart criterion

The characteristic equation

$$\lambda^4 - S_1\lambda^3 + S_2\lambda^2 - S_3\lambda + S_4 = 0$$

where

$$S_1 = \sum_{1 \leq k \leq 4} A \begin{pmatrix} k \\ k \end{pmatrix} = \sum_{1 \leq k \leq 4} M_{kk} = \text{Tr}(\mathbf{M}) \quad S_2 = \sum_{1 \leq k < l \leq 4} A \begin{pmatrix} k & l \\ k & l \end{pmatrix} \quad S_3 = \sum_{1 \leq k < l < m \leq 4} A \begin{pmatrix} k & l & m \\ k & l & m \end{pmatrix}$$

$$S_4 = A \begin{pmatrix} 1 & 2 & 3 & 4 \\ 1 & 2 & 3 & 4 \end{pmatrix} = |\mathbf{M}| \quad \text{are the sums of main minors of the matrix } \mathbf{M}.$$

For polynomial with real coefficients $f(z) = a_0 z^n + a_1 z^{n-1} + \dots + a_n, a_0 > 0$

Positivity of roots requires one of the following conditions:

1) : $a_n > 0, a_{n-2} > 0 \dots; \Delta_1 > 0, \Delta_3 > 0, \dots$

2) : $a_n > 0, a_{n-2} > 0 \dots; \Delta_2 > 0, \Delta_4 > 0, \dots$

3) : $a_n > 0, a_{n-1} > 0, a_{n-3} > 0 \dots; \Delta_1 > 0, \Delta_3 > 0, \dots$

4) : $a_n > 0, a_{n-1} > 0, a_{n-3} > 0 \dots; \Delta_2 > 0, \Delta_4 > 0, \dots$

Hurwitz determinant:

$$\Delta_i = \begin{vmatrix} a_1 & a_3 & a_5 & \dots & \dots \\ a_0 & a_2 & a_4 & \ddots & \ddots \\ 0 & a_1 & a_3 & \ddots & \ddots \\ 0 & a_0 & a_2 & a_4 & \dots \\ \dots & \ddots & \ddots & \dots & a_i \end{vmatrix} \bullet^{21}$$