

Wir schaffen Wissen – heute für morgen

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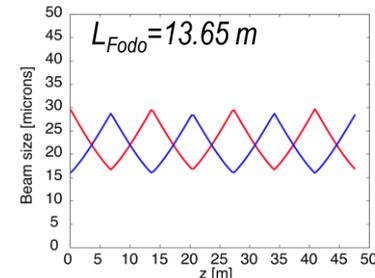
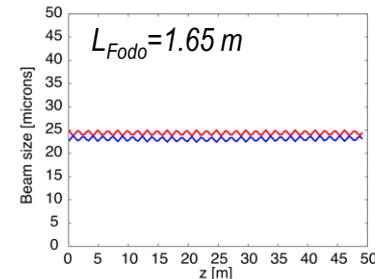
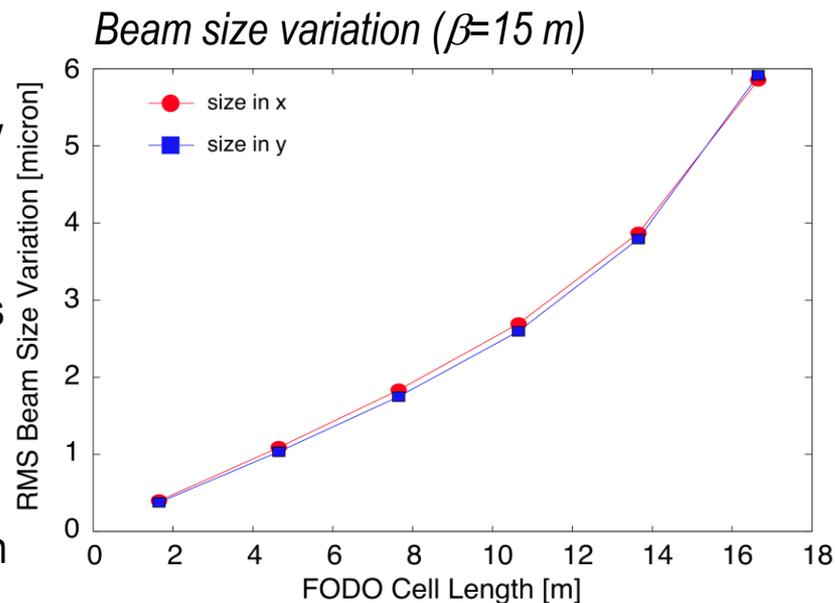
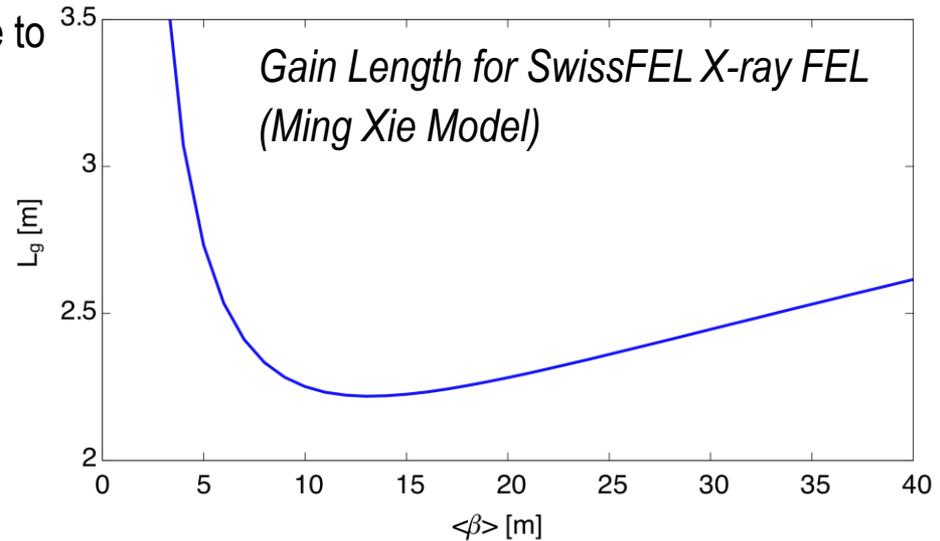
**Growth Rates and Coherence Properties of FODO-lattice
based X-ray Free Electron Lasers**

- **Focusing and FODO Lattices**
- **FEL Eigenmodes**
- **Coherence**
- **FEL Performance of hard X-ray FEL (1 Ångstrom)**
- **FEL Performance of soft X-ray FEL (1 nm)**
- **Optimization of Focusing Strength**

- External focusing increases the FEL performance due to the higher electron density and larger FEL parameter.
- Stronger focusing also transfers kinetic energy into betatron oscillation, slowing down the electrons proportionally to its betatron amplitude.

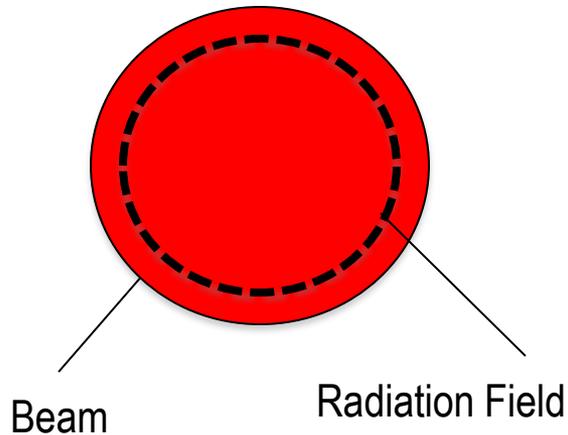
→ Optimizing the beta function

- FODO lattice to provide focusing
- Achievable beta-function is limited by the FODO cell length (instable limit)
- Longer wavelength or low emittances pushes optimum beta-value below limit
- Strong variations in beam sizes when operating at focusing limit

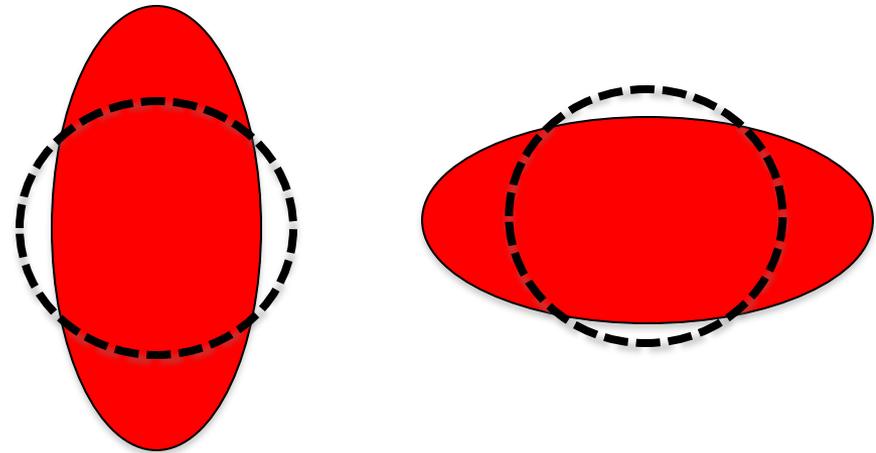


For low diffraction at hard X-ray FELs the optical mode stays inside the electron beam

Round, rigid electron beam



Beam Size Variation in FODO Lattice



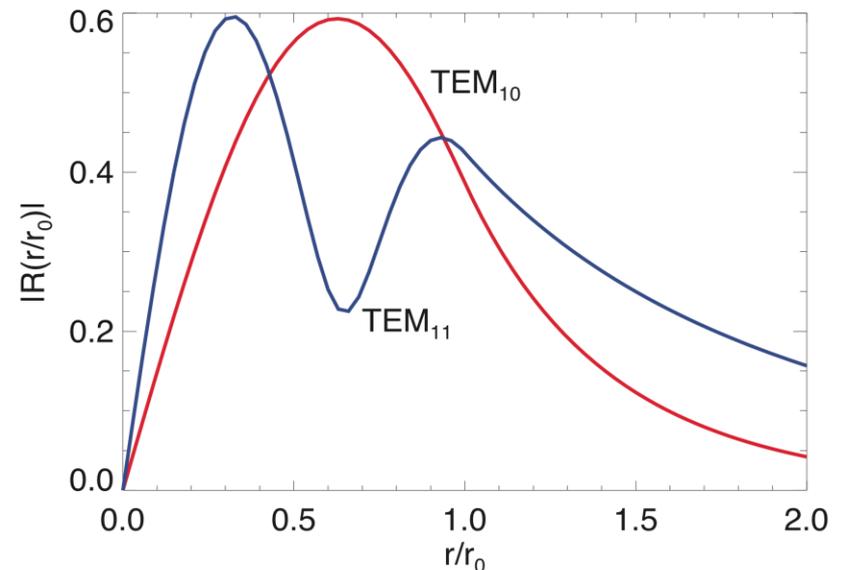
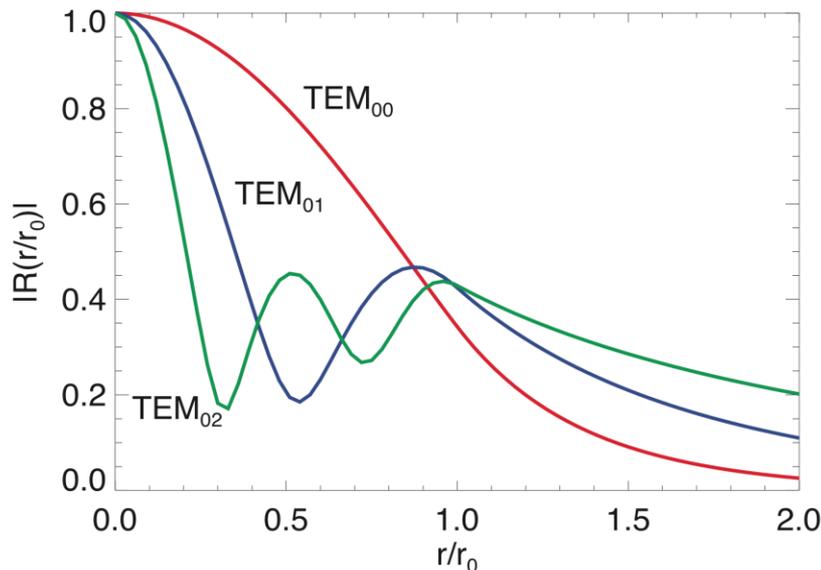
- Electrons are pushed out of the radiation field, reducing the amount of electrons emitting into the FEL mode
- Effect should be more pronounced for strong variation, e.g. FODO cells at the stable transport limit

- The FEL wave front with a growth rate Λ has to fulfill the field equation of the form:

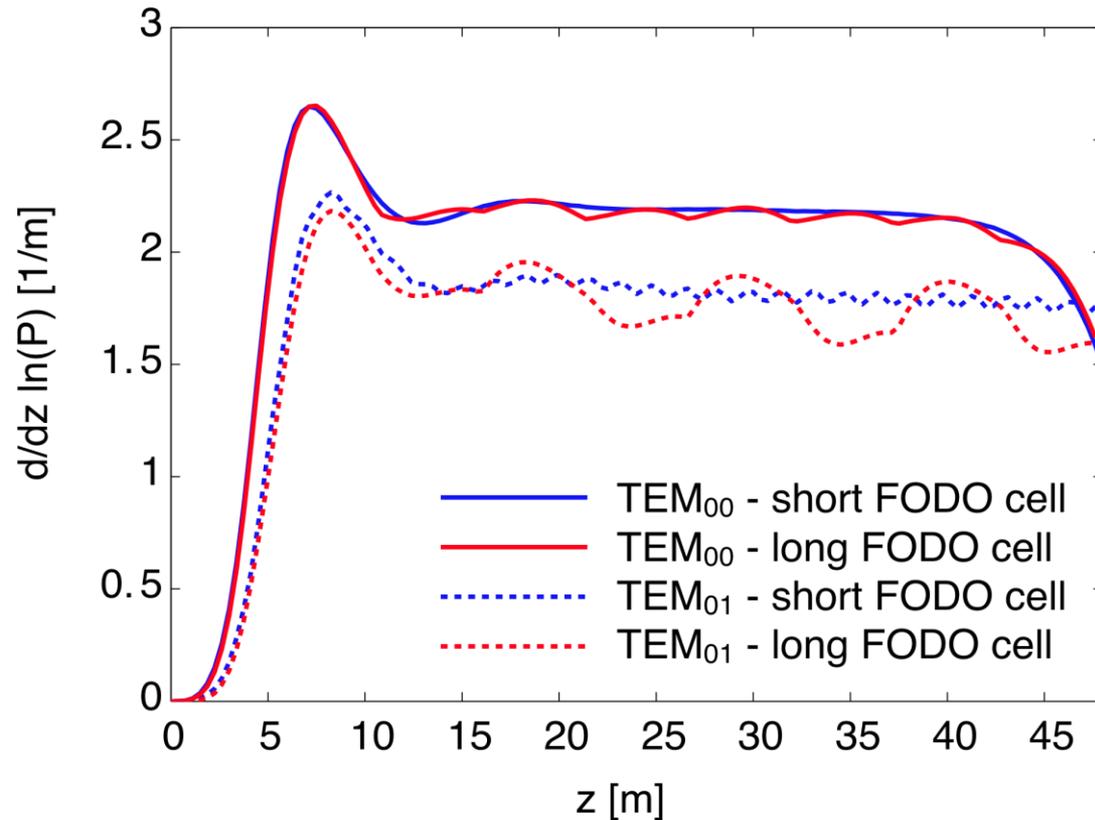
$$\left[\nabla_{\perp}^2 + V(r, \Lambda) \right] A = 0$$

- The transverse electron distribution acts like a 2D potential in Quantum Mechanics.
- Note: $V(r, \Lambda) = 0$ is the 1D dispersion equation with the given electron density at the radius r .
- Like in QM a valid solution for A has to be finite when $|A|^2$ is integrated over the transverse plane.
- Discrete solutions define growth rates Λ_{mn} for each FEL eigenmode.

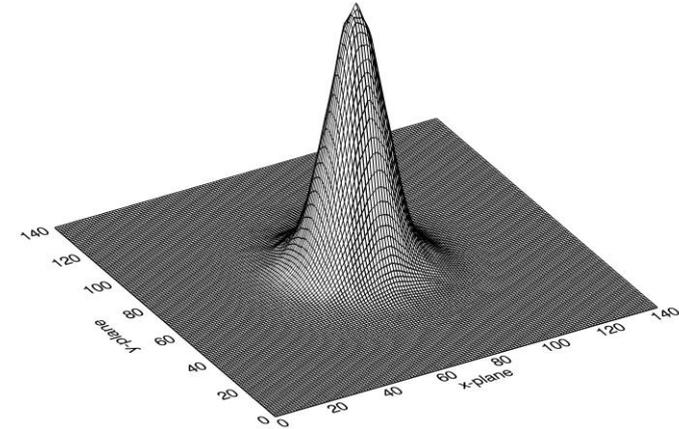
Example: $A(r, \phi) = R_{mn}(r) e^{im\phi}$ for step profile with radius r_0



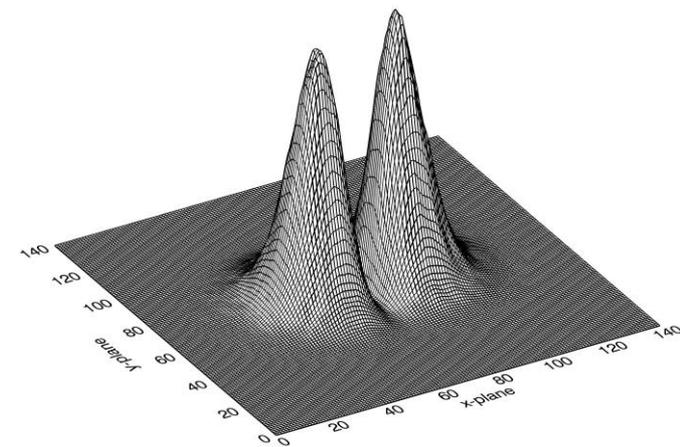
- Seeding with different Gauss-Hermite Modes
- Only TEM_{00} and TEM_{01} couple well to FEL eigenmodes
- TEM_{01} mode more sensitive to beam size variations
- Higher modes „leak“ into lower modes. Enhanced by stronger beam size variations



Fundamental Gauss Mode (TEM_{00})

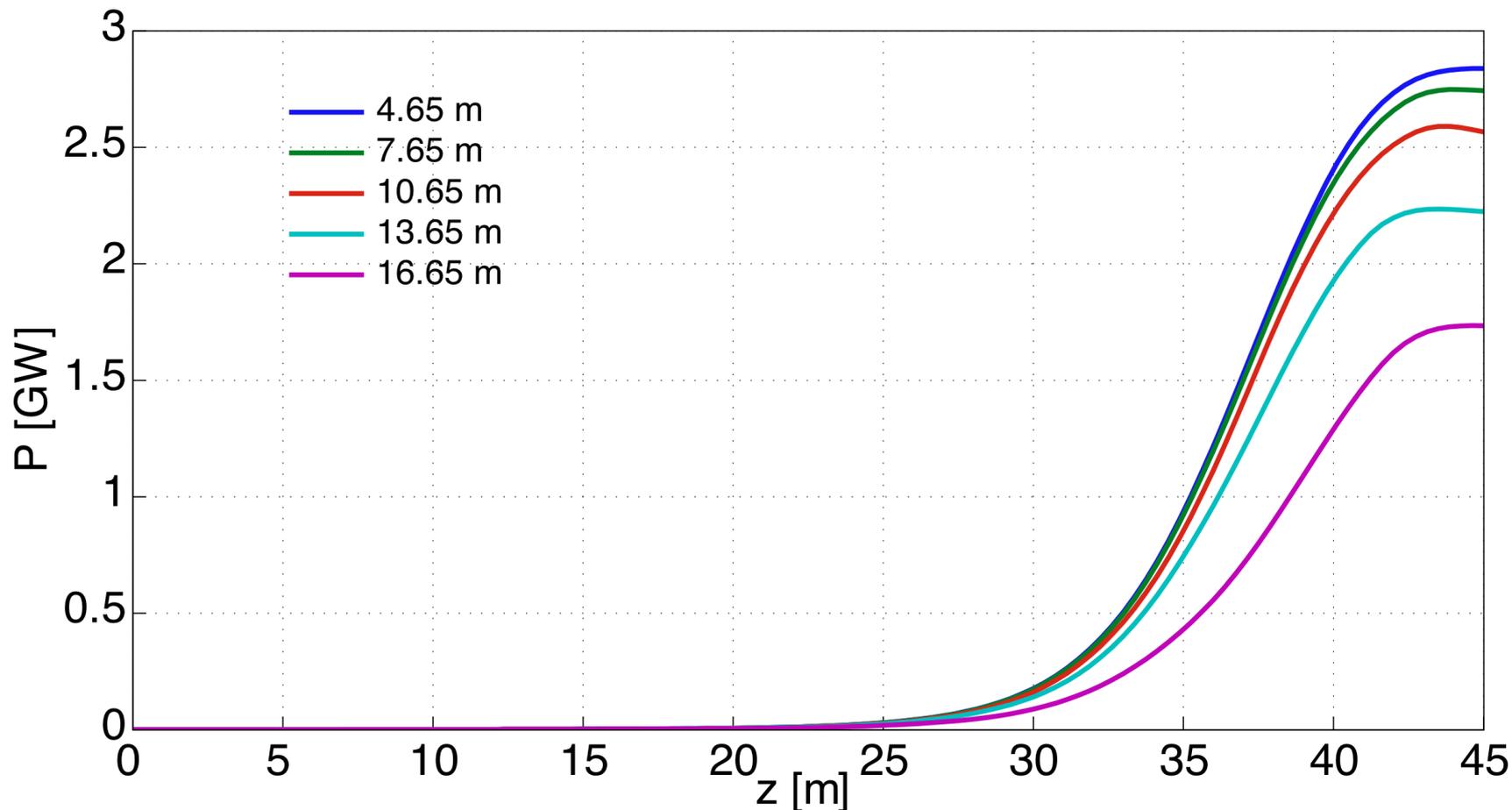


Gauss-Hermite Mode (TEM_{01})



FEL Performance for different FODO Cell Lengths

- Simulation with SwissFEL parameters at 1 Å for various cell length of the FODO lattice.
- Noticeable drop in saturation power for cell length of 10 m and longer, slight increase in saturation length



- Coherence is the correlation between two points in space and time, averaged over many shots:

$$\Gamma_{12}(\tau) = \langle E(\mathbf{r}_1, t + \tau) \cdot E^*(\mathbf{r}_2, t) \rangle$$

- Considering only transverse coherence (t=0) the complex coherence function

$$\mu_{12} = \frac{\Gamma_{12}(0)}{\sqrt{\Gamma_{11}(0)\Gamma_{22}(0)}}$$

Is a measure coherence between two points ($\mu_{12}=1 \rightarrow$ full coherence, $\mu_{12}=0 \rightarrow$ incoherent).

- A quantity describing the entire beam is the intensity weighted average of the coherence function:

$$\zeta = \frac{\int |\mu_{12}|^2 \Gamma_{11} \Gamma_{22} d\mathbf{r}_1 d\mathbf{r}_2}{\left[\int \Gamma_{11} d\mathbf{r}_1 \right]^2}$$

Very impractical to evaluate ζ , computational and memory wise

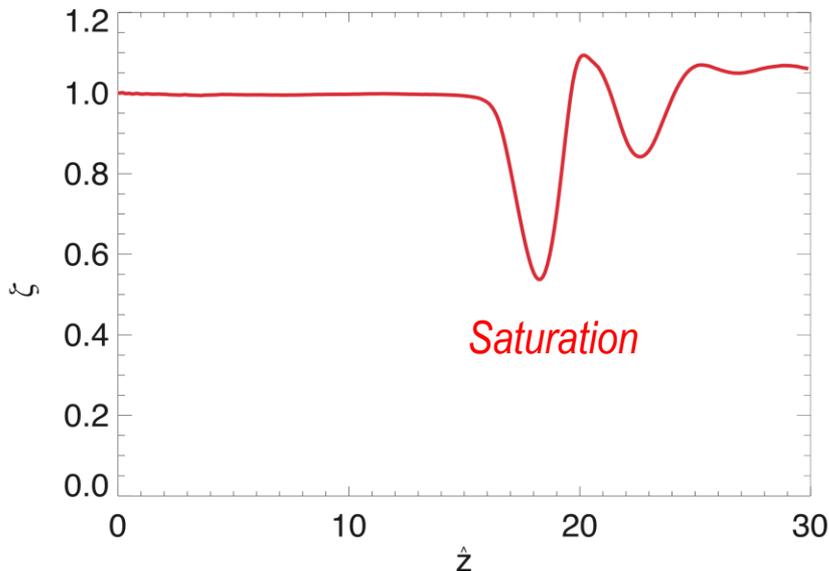
The SASE process is a stationary process and the average over many shots can be replaced by an average over a single long pulse.

Due to the seed by a “white noise” (shot noise) signal the coherence factor is identical to the fluctuation of the instantaneous power **within the linear regime** (see *SALDIN et al*)

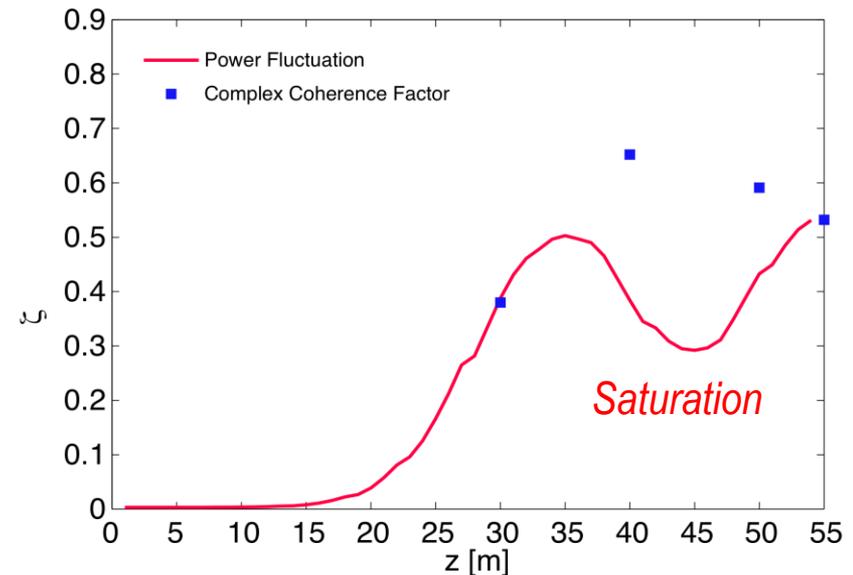
$$\zeta = \frac{\langle (P - \langle P \rangle)^2 \rangle}{\langle P \rangle^2} = \frac{1}{M_L \cdot M_T}$$

M_L : Number of longitudinal modes
 M_T : Number of transverse modes

1D Simulation ($M_T=1$)

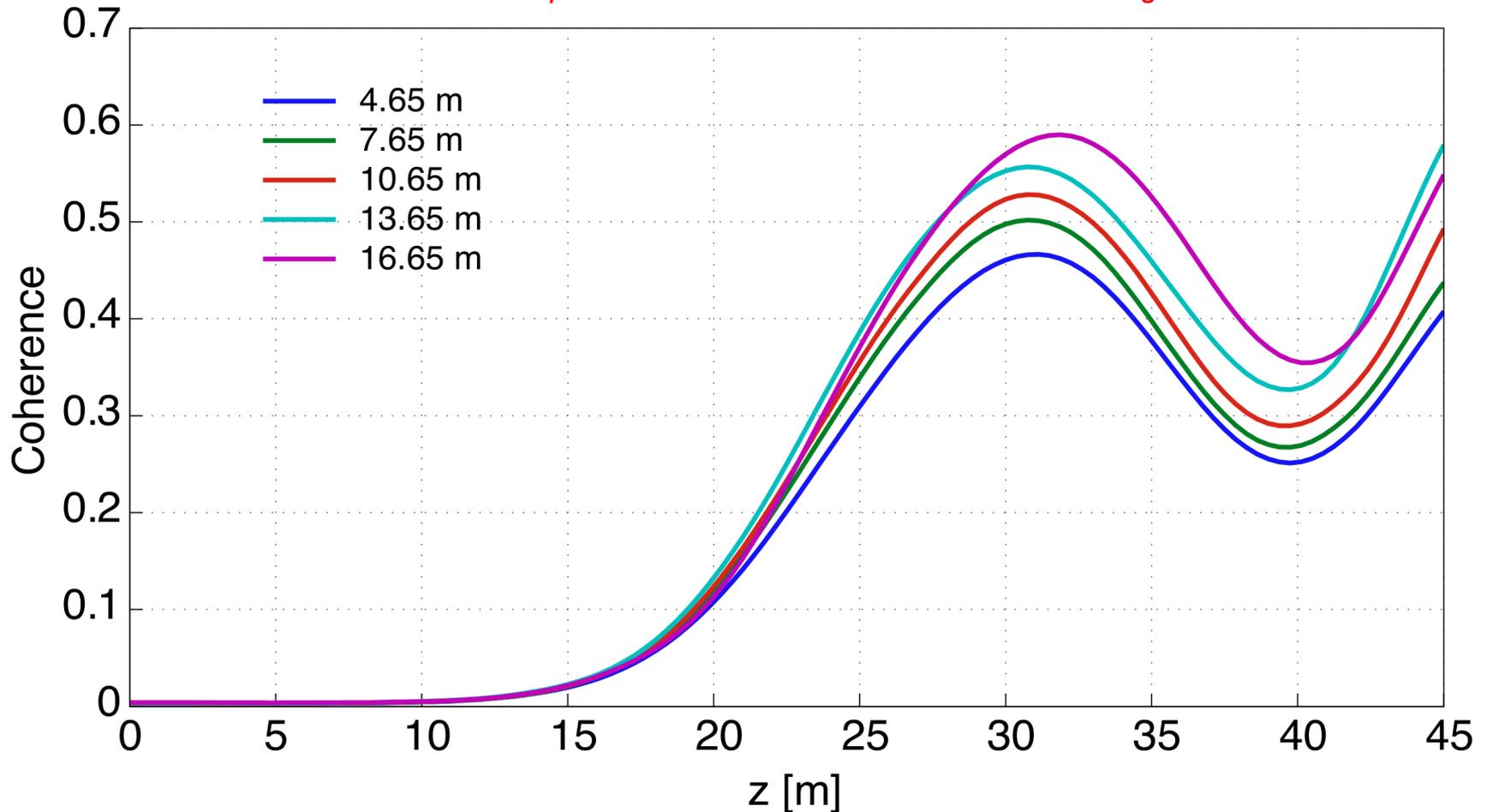


3D Simulation (Genesis 1.3)

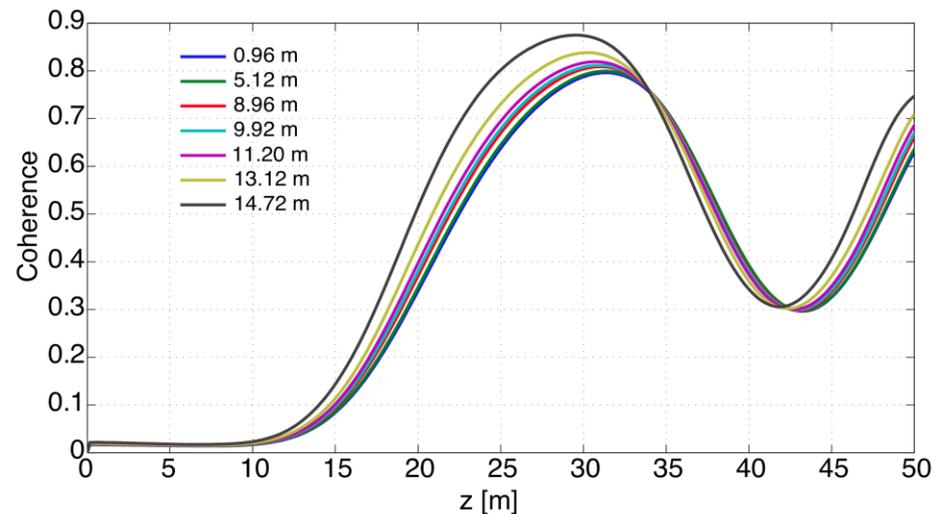
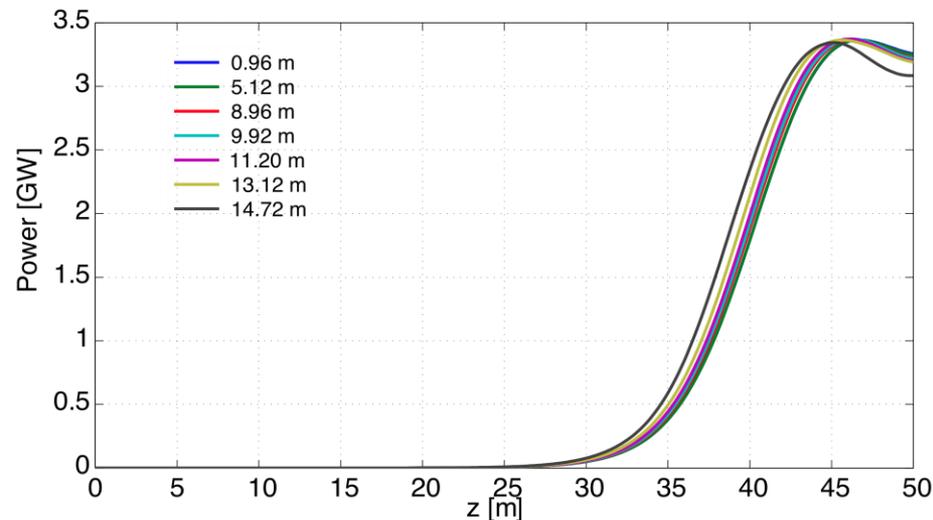


- Stronger growth in coherence for large beam size variations
- Higher modes more affected by FODO lattice than fundamental mode
- Dominance of fundamental mode grows faster → improved coherence

Difference of up to 30% between different FODO cell lengths



- Modeled after the SwissFEL soft X-ray beamline.
- In difference to hard X-ray case, the radiation mode (37 microns) is larger than electron beam size (32 microns).
- Very weak dependence of FEL performance on FODO cell length. Slight improvement arises from quadratic term of beta-function, which somehow reduces the average beam size.
- Still a noticeable dependence on coherence. A larger beam size variation improves upon coherence.



With given FODO lattice (9.5 m cell length) the beta function is optimized for best FEL performance.

Ming Xie Prediction for best gain length

$$\beta = 12 \text{ m}$$

Best gain length (Simulation)

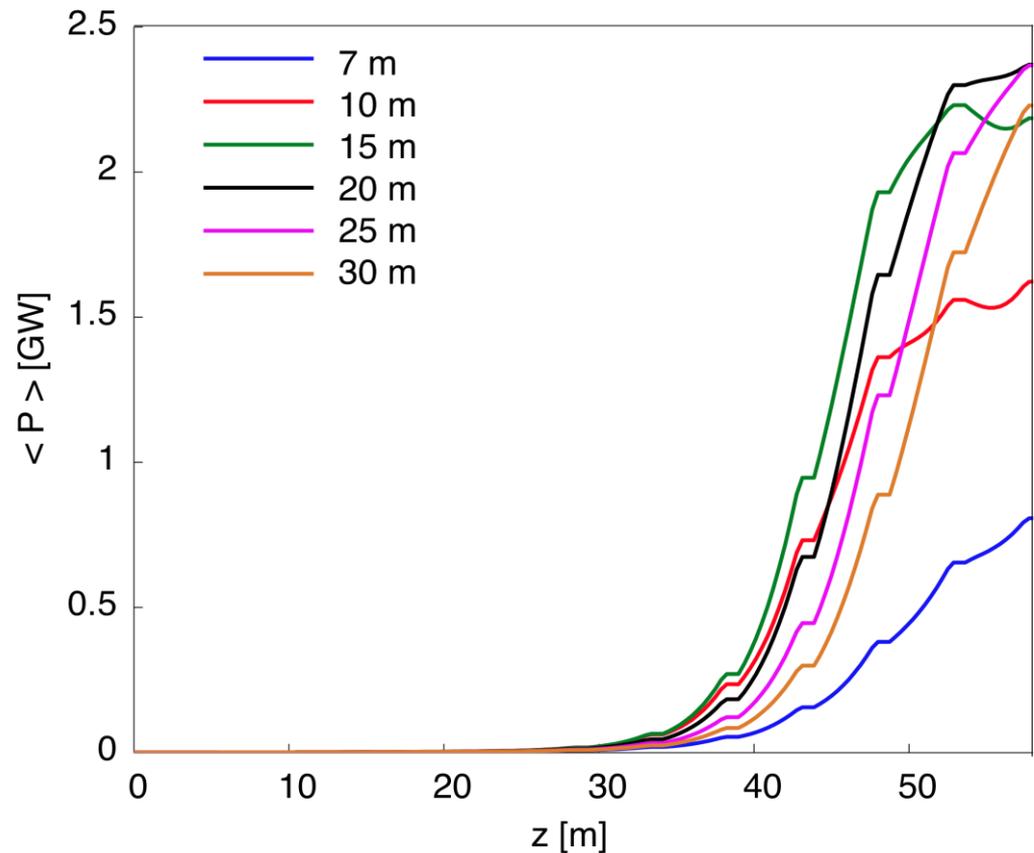
$$\beta = 15 \text{ m}$$

Best saturation power

$$\beta = 20\text{-}25 \text{ m}$$

Best coherence

$$\beta = 10 \text{ m}$$



FODO lattices are essential to optimize X-ray FELs

Strong variation of electron beam sizes:

- **Less saturation power and longer saturation length**
- **Improved coherence for SASE FEL**

Effect is not seen for longer wavelength (1 nm and longer)

Relaxed focusing for given FODO lattice can increase saturation power.