



Progress towards X-ray FEL0 and Concept for Sub-Å Stabilization for Nuclear Resonance Metrology



Kwang-Je Kim

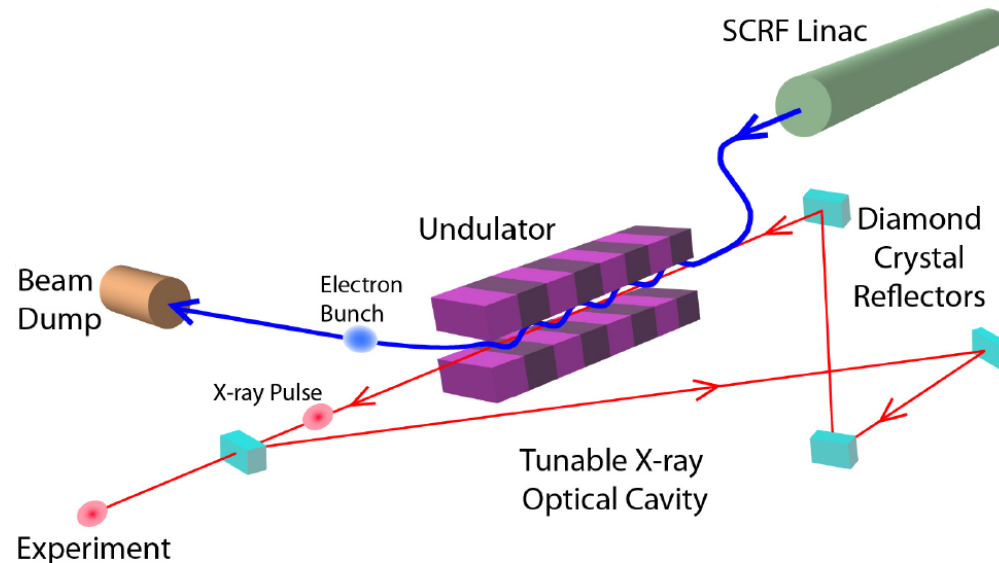
ANL & U of C

FEL2012

August 26-31, 2012

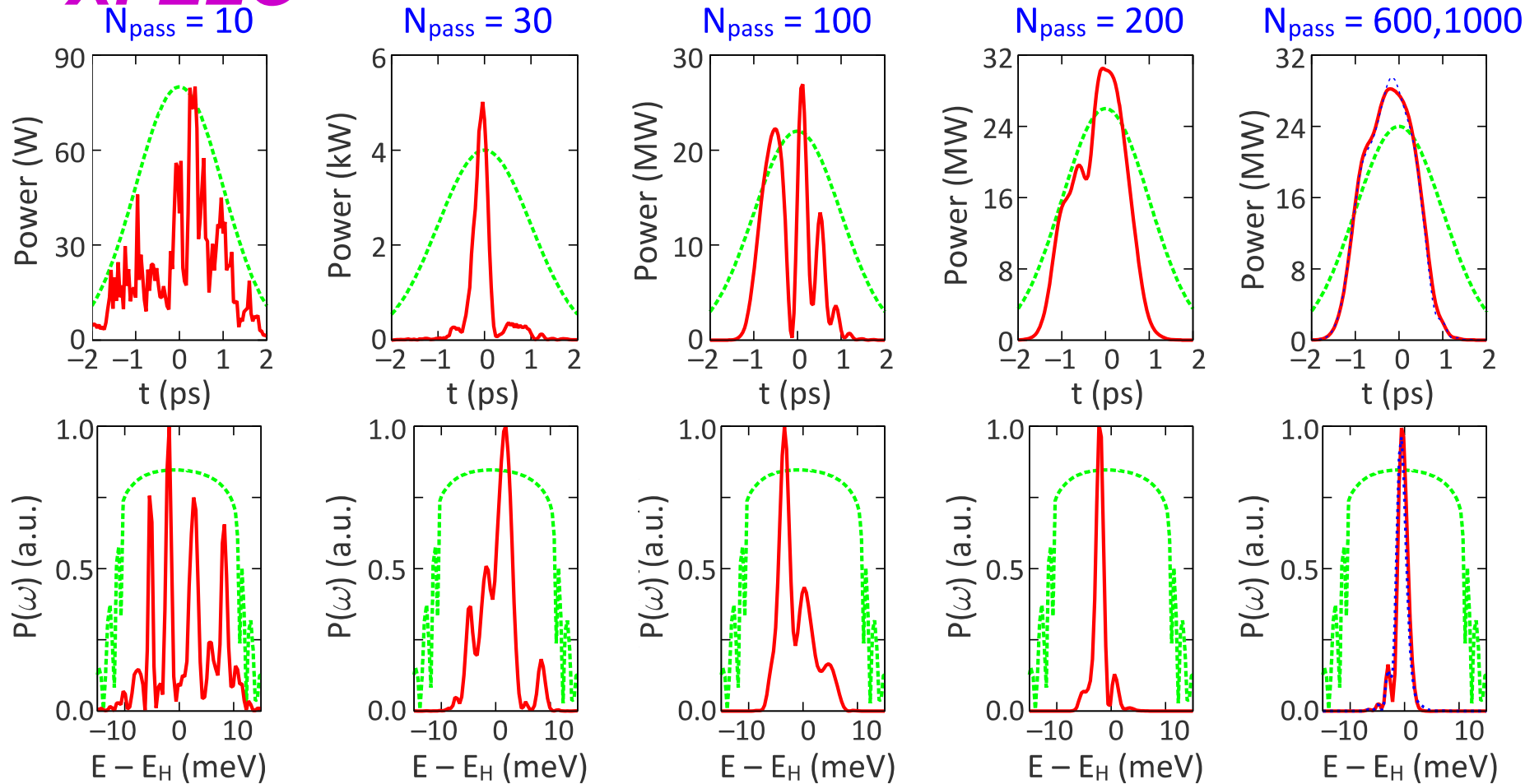
Nara, Japan

X-Ray FEL Oscillator (XFEL-O)



- **X-ray FEL oscillator is feasible by using Bragg mirrors**
 - R. Collela and A. Luccio, 1984; KJK, Y. Shvyd'ko, and S. Reiche, 2008)
- **Tuning is possible with a four mirror configuration**
 - R. M.J.Cotterill, (1968)KJK & Y. Shvyd'ko (2009)
- **Crystal choice becomes independent from Bragg planes→ choose diamond for its best mechanical and thermal properties**

Temporal and spectral evolution of XFELs



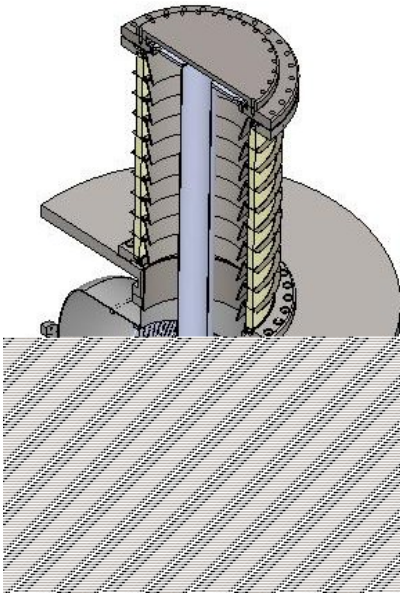
R.R. Lindberg, K.-J. Kim, Yu. Shvyd'ko, and W.M. Fawley, *Phys. Rev. ST-AB*. **14**, 010701 (2011)

Example Parameters

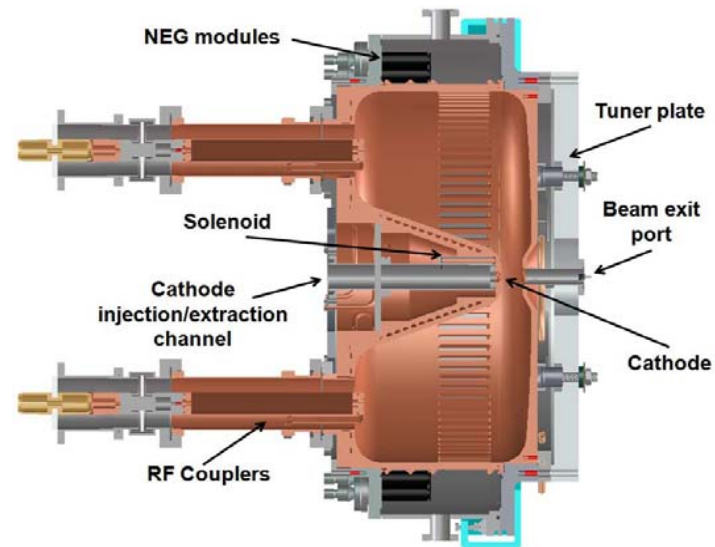
- Electron bunches: 7 GeV, 25 pC, 0.5 ps (rms)
- Undulator: $K=1.38$, $\lambda_U=1.65$ cm, $L_U=30$ m
- X-ray cavity roundtrip length: 90 m
- Diamond (337) crystal
- 30% gain=15% RT loss+5% out-coupling+10% net gain
- Output x-ray pulse characteristics—complementary to high-gain XFEL:
 - $\hbar\omega=14.4$ keV, > 5% tuning range
 - 6.5×10^8 photons/pulse,
 - Full coherence and high stability
 - BW : $\Delta\hbar\omega \sim 3$ meV (FWHM) \rightarrow 700 fs(FWHM)
 - Rep rate 3 MHz (RT length=90m)
- See R. R. Lindberg and KJK (WEPD27), J. Zemella, et al.,(WEPD29), R. Hajima, et al., (WEPD30)

Accelerator technology: Injector

- Injectors for ERL projects (KEK/JAEA, Cornell) at the soft x-ray FEL NGLS at NGLS will meet the XFEL requirements



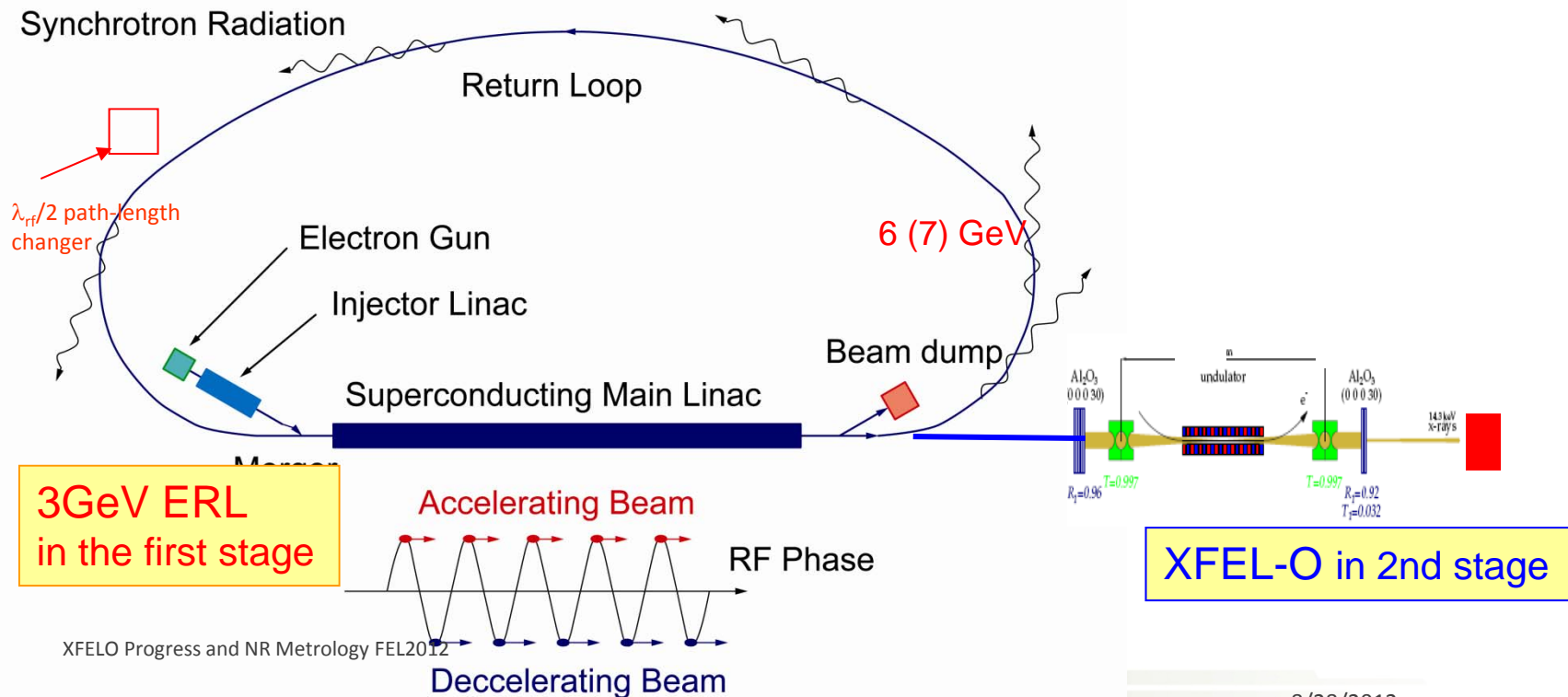
N. Nishimori, et al, WEPD51



C. Papadopoulos, et al., MOPD31

Main accelerator

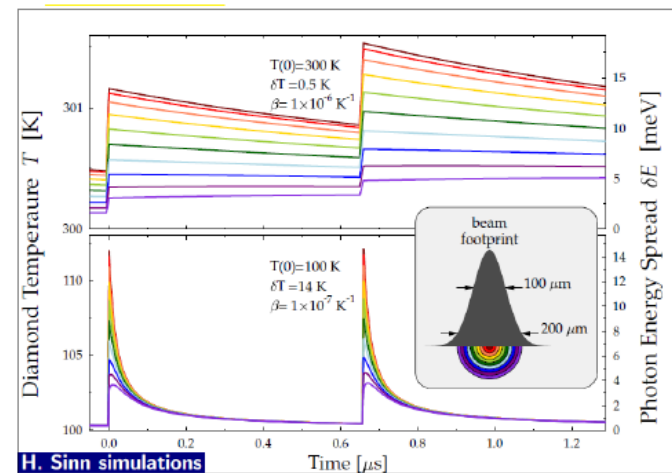
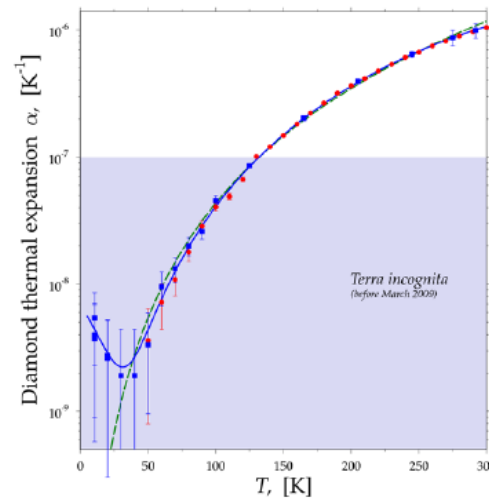
- CW SCRF accelerators for >5 GeV is feasible with or without recirculation passes: ERLs at KEK/JAEA, Cornell
- A pulsed mode XFEL can be operated at European XFEL (J. Zemella, et al., WEPD29)
- Most parameters of an ultimate storage ring retrofitted to large HEP rings (PEP, Tevatron, PETRA<>>) are compatible except for the longitudinal emittance



XFEL Progress and NR Metrology FEL2012

Synthetic diamond for an XFEL O appears feasible

- High reflectivity >99% has been demonstrated
 - Yu. V. Shvyd'ko, S. Stoupin, V. Blank, and S. Terentyev, Nat. Photonics 5, 539 (2011)
- Intra- and inter pulse heating/expansion can be avoided by cooling the diamond crystals to $T < 100\text{K}$
 - S. Stoupin and Yu. V. Shvyd'ko, Phys. Rev. Lett. 104, 085901
- The charge imbalance due to photo-emission may be avoided in high-purity crystal (BNL diamond amplifier gun)



***XFEL**O applications*

Science Opportunities with an XFEL
Workshop, APS, May 5th, 2010



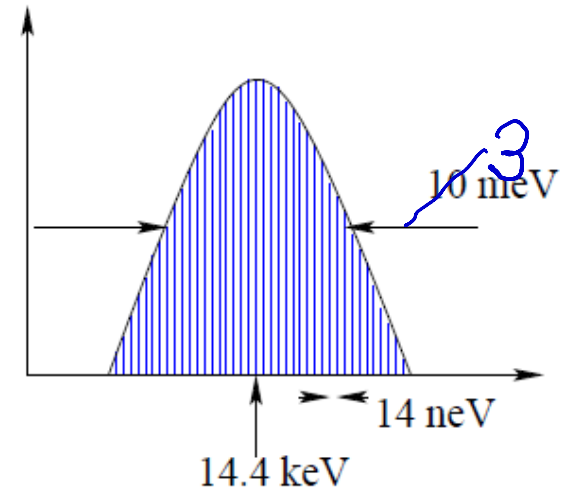
- **High resolution spectroscopy**
 - **Inelastic x-ray scattering**
- **Mössbauer spectroscopy**
 - **10^3 /pulse, 10^9 /sec Moessbauer γ s (14.4 keV, 5 neV BW)**
- **X-ray photoemission spectroscopy**
 - **Bulk-sensitive Fermi surface study with HX-TR-AR PES**
- **X-ray imaging with near atomic resolution (~ 1 nm)**
 - **Smaller focal spot with the absence of chromatic aberration**
- ***Capture coherent atomic motion (< 1 ps resolution) in real space***

Further opportunities for an XFEL-O: Extreme temporal coherence length from sub-mm \rightarrow km $\rightarrow 10^8$ m (spontaneous emission limit)!

- The XFEL-O output pulses, being a copy of the same circulating intra-cavity pulse, are phase-coherent
- If the pulse spacing T is controlled accurately, $\Delta T \ll \lambda/c$, the spectrum of XFEL-O output has a comb structure
- The stabilized XFEL-O may establish x-ray-based length standard and have applications in some fundamental physics

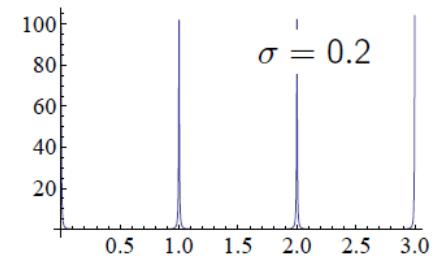
Comb structure of a stabilized XFEL

- Comb spacing: $\hbar\omega_T = \hbar 2\pi/T = 14 \text{ neV}$ is well-matched to ^{57}Fe resonance: $\hbar\omega = 14.4 \text{ keV}$, $\Gamma = 4.8 \text{ neV}$
- There are $\sim 2 \times 10^5$ lines (modes) in the FEL BW:
 $\hbar\Delta\omega = 3 \text{ meV}$
- The integer n is centered around $n_R = \omega_R/\omega_T = L/\lambda_R \rightarrow$ the comb line moves one unit if the cavity length L is changed by λ_R
- $6.5 \times 10^8 / 2 \times 10^5 = 3000 \text{ photons/mode}$

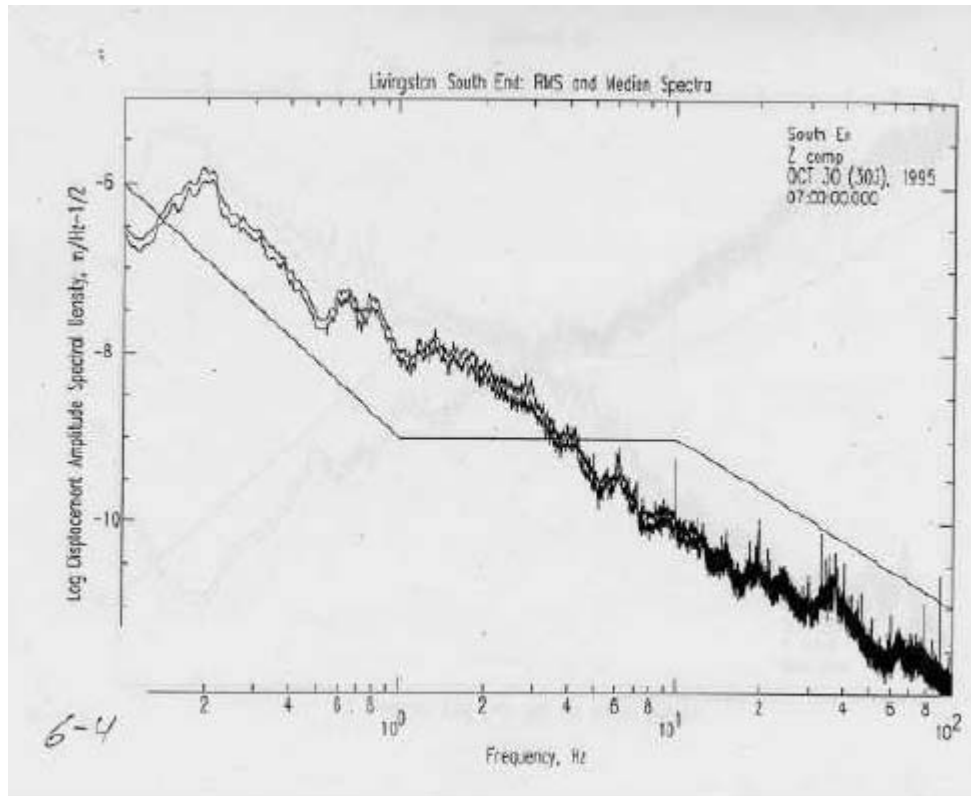


Line broadening due to phase noise

- $E(t) = \sum_n e^{i\psi_n} E_0(t - nT)$ $\psi_n = \varphi_n + \sum_{k=0}^n \phi_k$
- **Uncorrelated:** φ_n & **cumulative:** ϕ_k
- **Gaussian random variables:** $\langle \varphi_k^2 \rangle = \sigma_\varphi^2$, $\langle \phi_k^2 \rangle = \sigma^2$
- $\tilde{E}(\omega) = \frac{1}{\sqrt{2\pi}} \int dt e^{i\omega t} E(t) = \sum_{n=1}^N e^{i\psi_n} e^{in\omega T} \tilde{E}_0(\omega)$
- **Without phase error, this is a spectral comb with envelope $\tilde{E}_0(\omega)$ with peaks at $\omega = n\omega_T$; $\omega_T = 2\pi/T$**
- $\langle |\tilde{E}(\omega)|^2 \rangle = NF(\sigma_\varphi^2, \sigma^2, \omega) |\tilde{E}_0(\omega)|^2$
- $F(\sigma_\varphi^2, \sigma^2, \omega) = \left(1 - e^{-\sigma_\varphi^2} + e^{-\sigma_\varphi^2} \frac{\sinh(\sigma^2/4) \cosh(\sigma^2/4)}{\sinh^2(\sigma^2/4) + \sin^2(\pi\omega/\omega_T)} \right)$
- $\sigma \ll 1, \sigma_\varphi \ll 1, \omega \approx n\omega_T \Rightarrow F = \frac{\sigma^2}{\sigma^2 + (2\pi(\omega - n\omega_T))^2}$
- **Spontaneous emission: $\sigma \sim 10^{-4}$**



Ground motion(seismic)



<http://www.ligo.caltech.edu/docs/G/G010325-00.pdf>

- Empirical formula for spectral power:

$$SP = 10^{-18} \left(\frac{10\text{Hz}}{f} \right)^4 \left[\frac{m^2}{\text{Hz}} \right]$$

- Require

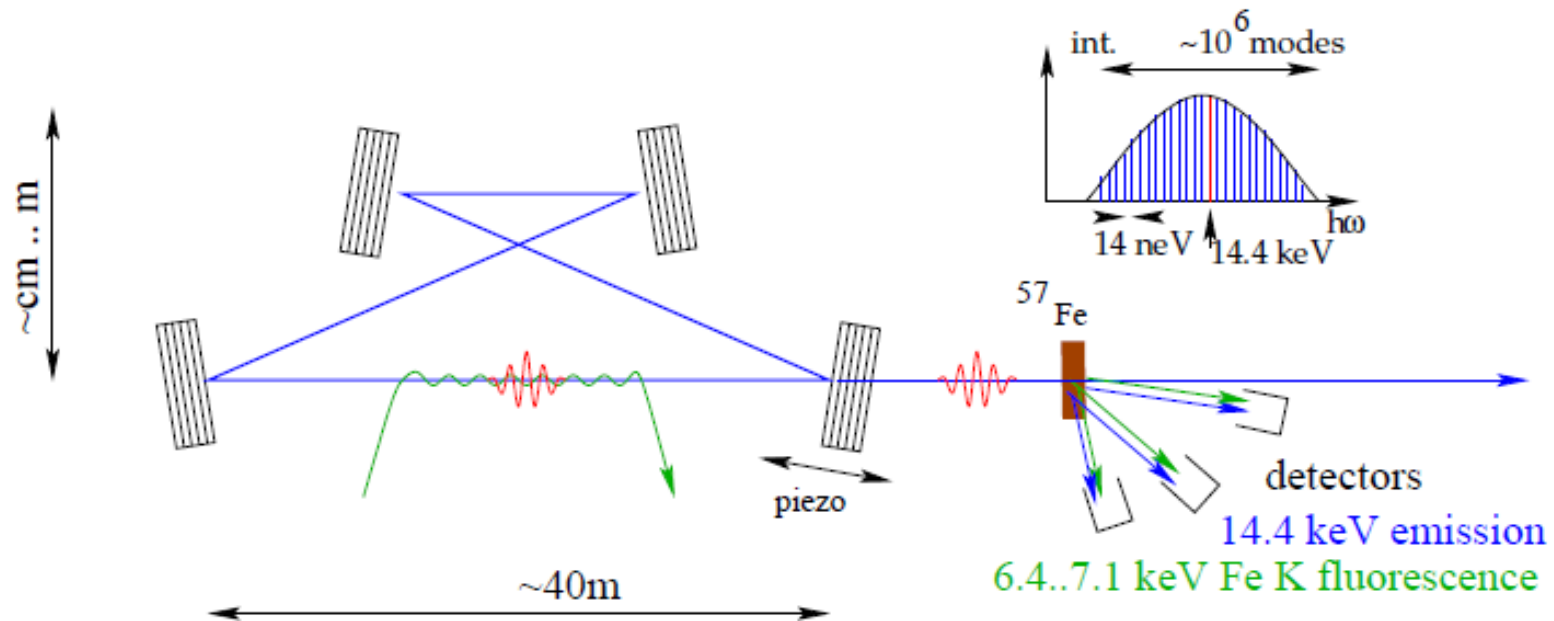
$$\sqrt{\int_{f_0}^{\infty} df SP} \leq \frac{\lambda}{100}$$

→ $f_0 > 2 \text{ kHz}$ for $\lambda \sim 1 \text{ \AA}$

- Feedback for $f_0 < 2 \text{ kHz}$

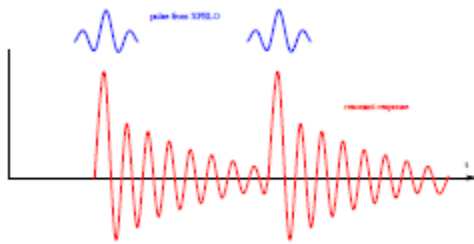
Cavity stabilization by tying one of the comb lines to a nuclear resonance line (B.W. Adams and K.-J. Kim, WEPD31)

Find the signal maximum by changing the cavity length $\sim 1 \text{ \AA}$ and lock to it

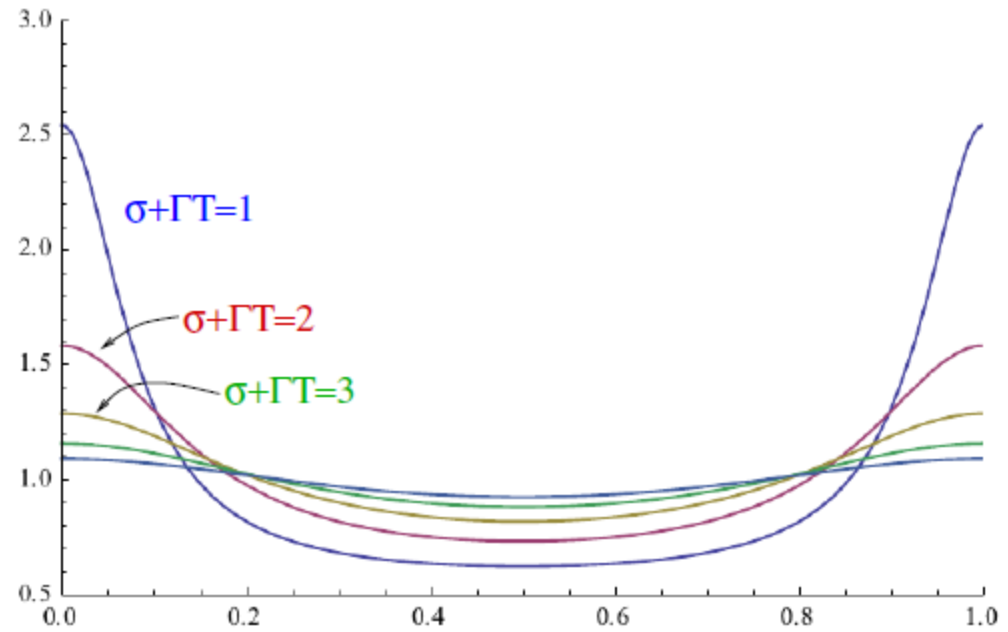


Excitation of nuclear resonance at ω_R of width Γ

- $$n_{abs}(\omega_R, \Gamma) = \kappa \int d\omega \frac{(\Gamma/2)^2}{(\omega - \omega_R)^2 + (\Gamma/2)^2} \langle |\tilde{E}(\omega)|^2 \rangle$$
$$= Nn_0 F(\sigma_\phi^2, \sigma^2 + \Gamma T, \omega_R)$$
- For $\Gamma T = 2.14$ (corresponding to ^{57}Fe and $T=300$ ns) the contrast is still sufficient



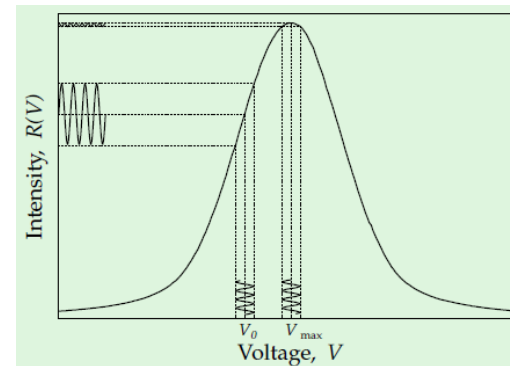
Contrast reduction due to ring-down between pulses



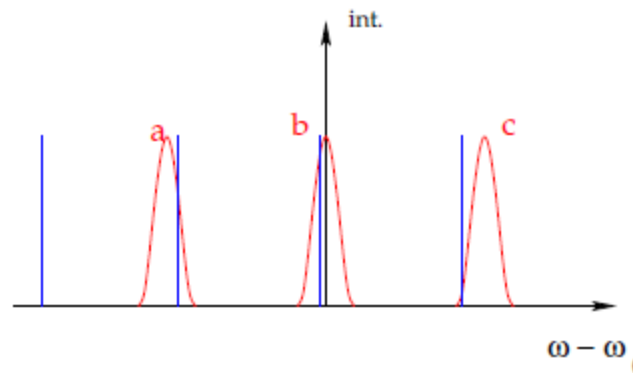
Feedback to prevent line drift at time scale $t > \tau = 1/f_0$

Methods to detect the tuning signal

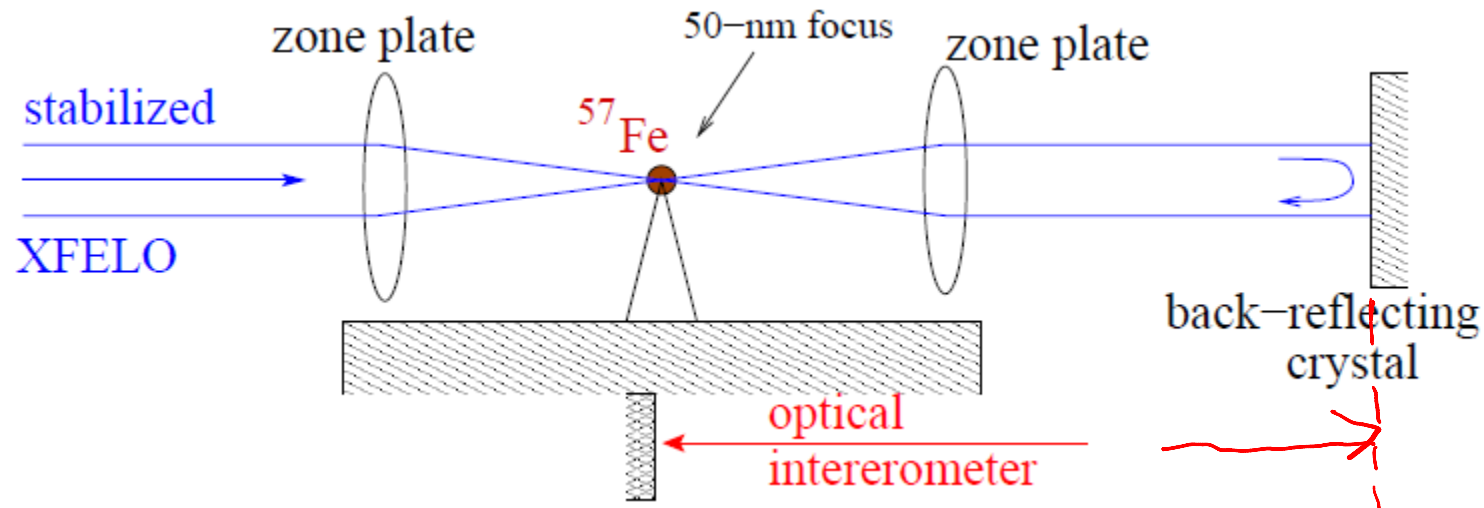
- Null-detection FB used at LIGO—dithering the crystal position



- Employ three samples slightly off-tuned to each other by relative motion



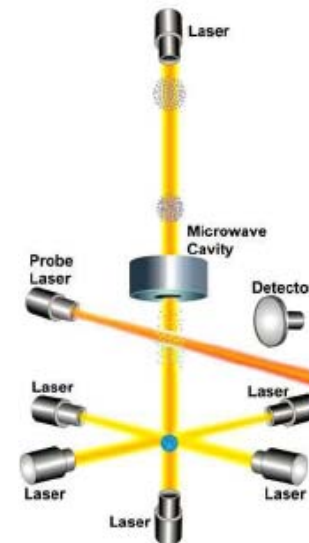
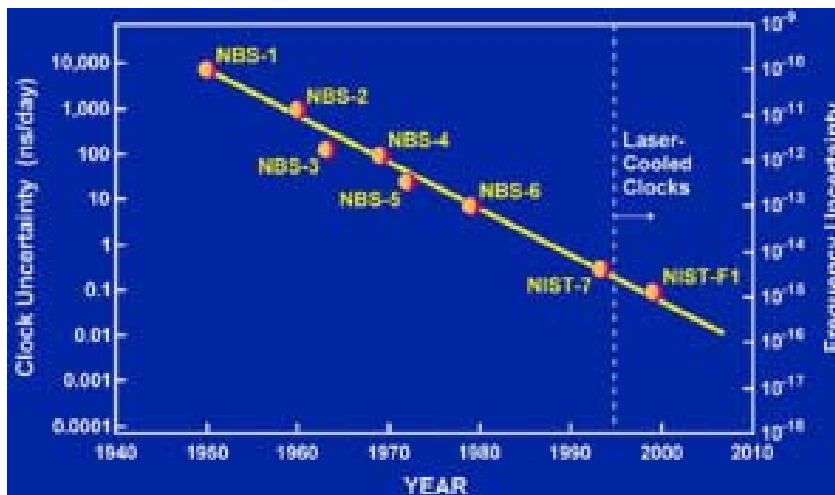
Precision measurement of NR wavelengths



- Stabilization does not determine absolute value of NR wavelength
- Form standing wave pattern using the output of the stabilized XFEL-O using ^{57}Fe resonance
- Count the interference peaks with a single ^{57}Fe atom fixed to a travelling assembly travelling a distance Δz ($< cT/2$) precisely determined by optical length standard. Assume $\Delta z = 10$ m.
- Not necessary to probe all of the 10^{12} peaks! Sparse sampling with pre-existing knowledge of λ

X-ray time standard

- Measure ^{57}Fe λ_R to 10^{-11} accuracy (currently known to 10^{-7}), even to 10^{-13} by interpolation.
- Other NR isotopes for higher accuracy, eg., ^{45}Sc (10^{-22})
- Replace atomic-clock standards with NR standard
- Nuclei are much better isolated from environmental perturbations than the electronic transitions used in atomic clocks



most accurate: Al^+ ion in trap: $8.6 \cdot 10^{-18}$ (arXiv:0911.4527)

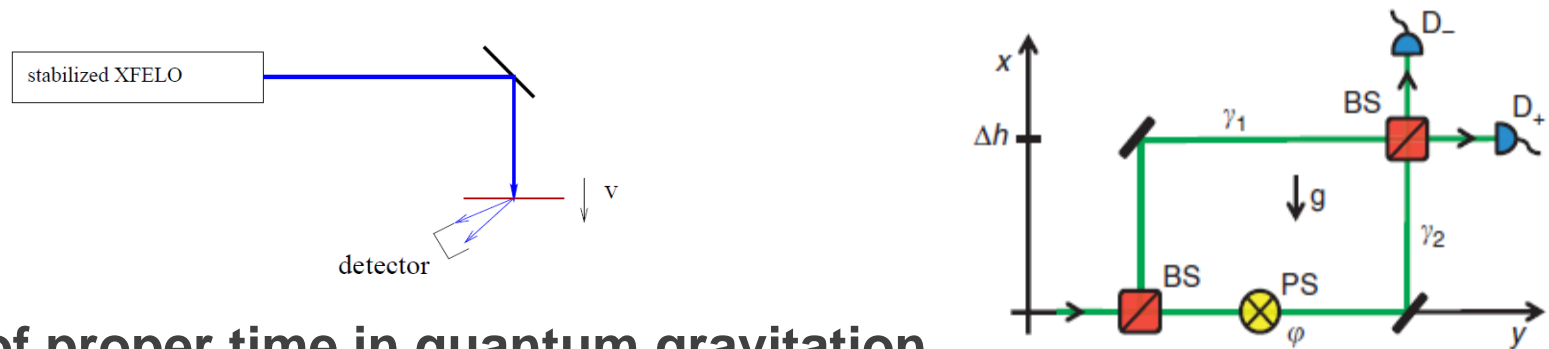
<http://www.nist.gov/pml/div688/grp50/primary-frequency-standards.cfm>

Characteristics of NR-stabilized XFEL-O enabling applications for technology and fundamental sciences

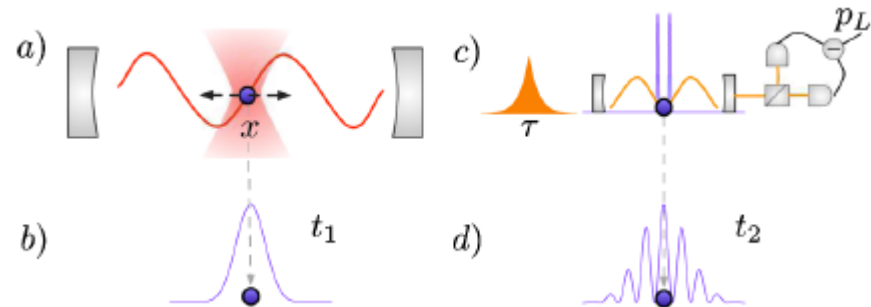
- **Higher wavelength accuracy**
 - as small as 10^{-22}
- **Small dimension of moving clock**
 - one NR atom ~ 1 pm
 - A cm-size atomic clock is feasible but with less accuracy $\sim 10^{-11}$
- **X-rays can be focused to a smaller spot-size**

Some possible applications

- Precision measurement of B- and g-field → “find oil !”



- The role of proper time in quantum gravitation
 - M. Zych, F. Costa, I. Pikovski, C. Brukner, Nature Comm. (2011)
- Macro-particle aspects in collapse of quantum wavefunction (the fate of Schrödinger cat)—improve the resolution of experiment proposed by Kaltenbaek, et al., PRL 107, 020405 (2011);



Conclusions

- An XFEL produces fully coherent, meV BW radiation, complimentary to high-gain XFEL
- Accelerator technology for an XFEL is available
- Diamond optics appears feasible/promising
- An XFEL will drastically improve experimental techniques developed for the 3rd generation sources
- Extending the XFEL to produce x-ray comb stabilized by locking to NR
- A stabilized XFEL may have application in nanotechnology and fundamental sciences