

First Demonstration of Optical Frequency Shot-Noise Suppression in Relativistic Electron-Beams and implications to FEL Coherence Enhancement

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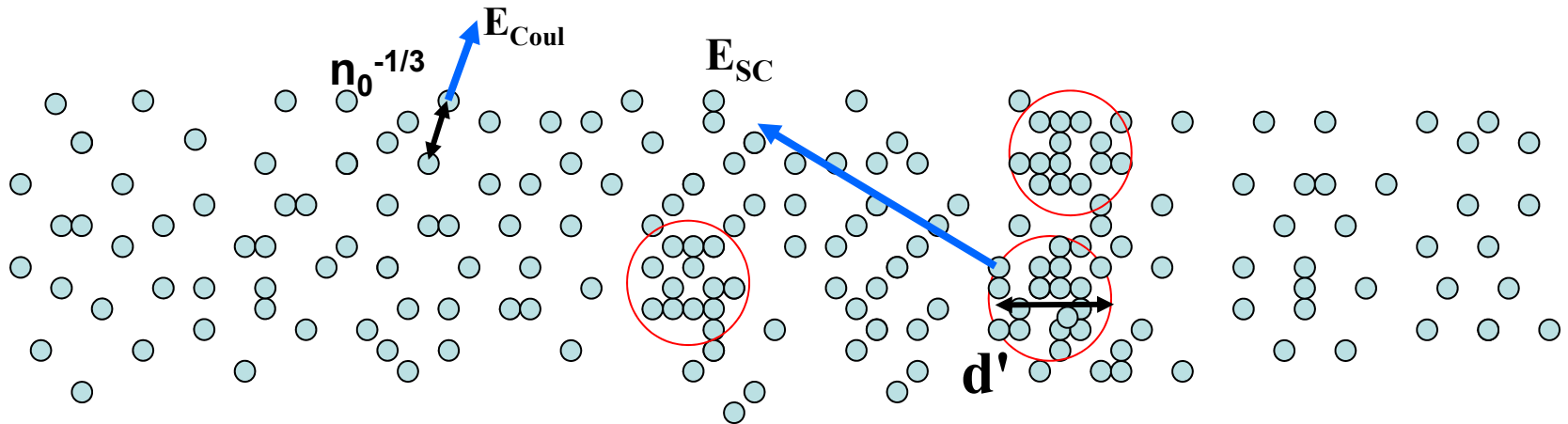


Physics of Collective Micro-Dynamics in a Charged Particle Beam: Homogenization trend

In the beam rest frame:

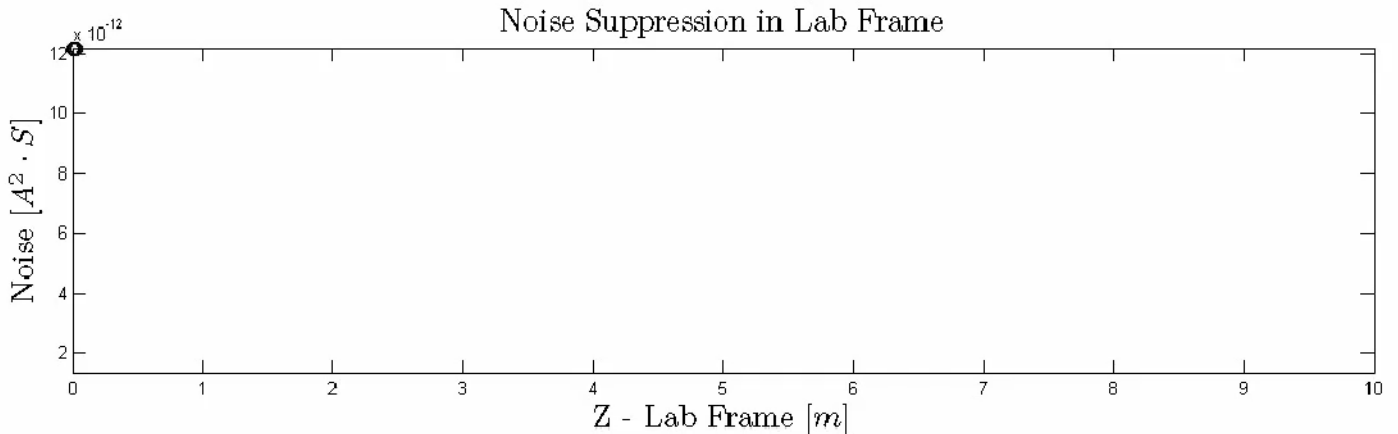
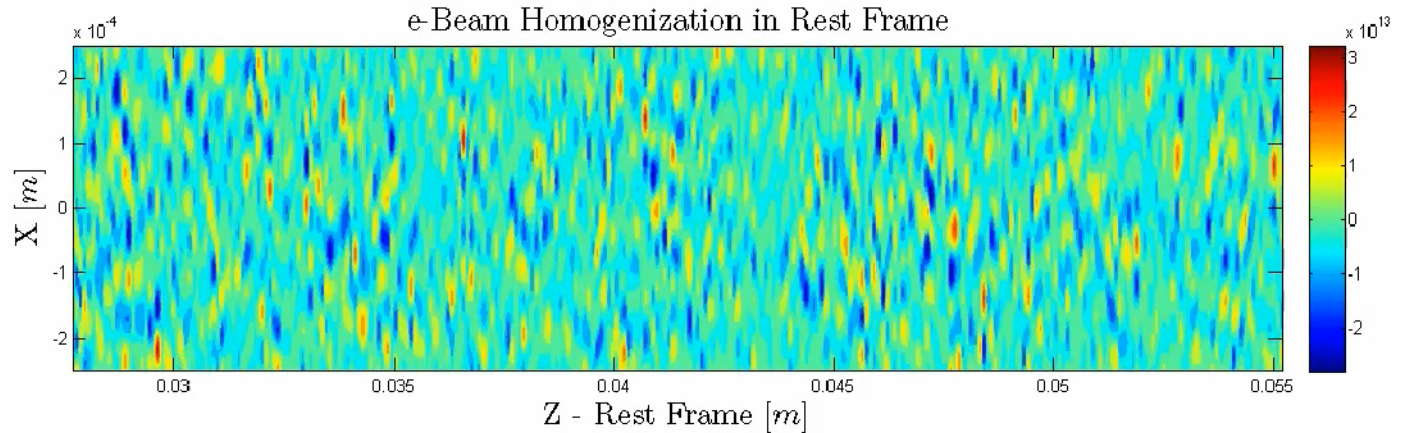
When $\epsilon_{sc} > \epsilon_{coul}$?

Answer: $d' > n_0^{-1/3}$



HOMOGENIZATION ($\lambda=5-10 \mu\text{m}$)

Density
In beam
frame



Nause, A. Dyunin, E. Gover, A. Optical frequency Shot- Noise suppression in electron beams: 3-D analysis. *J. of Appl. Phys.* **107**, 103101 (2010).

ANALYTICAL FLUID-PLASMA LINEAR MODEL

[H. Haus and F. N. H. Robinson, Proc. IRE 43, 981 (1955)]

[A. Gover, E. Dyunin, PRL 102, 154801 (2009)]

Plasma Oscillation in a Uniform e-Beam Drift Section

$$\begin{aligned}\tilde{i}(L_d, \omega) &= \left[\tilde{i}(0, \omega) \cos \phi_p - i \tilde{V}(0, \omega) (\sin \phi_p / W_d) \right] e^{i\phi_b(L_d)} \\ \tilde{V}(L_d, \omega) &= \left[-i \tilde{i}(0, \omega) W_d \sin \phi_p + \tilde{V}(0, \omega) \cos \phi_p \right] e^{i\phi_b(L_d)}\end{aligned}$$

Kinetic voltage:
(axial velocity modulation)

$$\tilde{V}(z, \omega) \propto \tilde{\gamma}(z, \omega) \propto \tilde{v}_z(\omega)$$

Optical phase:

$$\phi_b = \frac{\omega}{v_z} L_d$$

Plasma phase:

$$\phi_p = \theta_{pr} L_d \quad \theta_{pr} = r_p \frac{\omega_{pL}}{v_0} \quad , \quad \omega_{pL} = \left(\frac{e^2 n_0}{m \epsilon_0 \gamma^3} \right)^{1/2}$$

CURRENT SHOT-NOISE SUPPRESSION

$$gain = \frac{\overline{|\tilde{i}(L_d, \omega)|^2}}{\overline{|\tilde{i}(0, \omega)|^2}} = \cos^2 \phi_p + N^2 \sin^2 \phi_p$$

$$N^2 = \frac{\overline{|\tilde{V}(0, \omega)|^2}}{W_d^2 \overline{|\tilde{i}(0, \omega)|^2}} = \left(\frac{\lambda_D}{\lambda} \right)^2$$

$$\left(k_D = \frac{2\pi}{\lambda_D} = \frac{\omega_{pL}}{\delta v_z} \right)$$



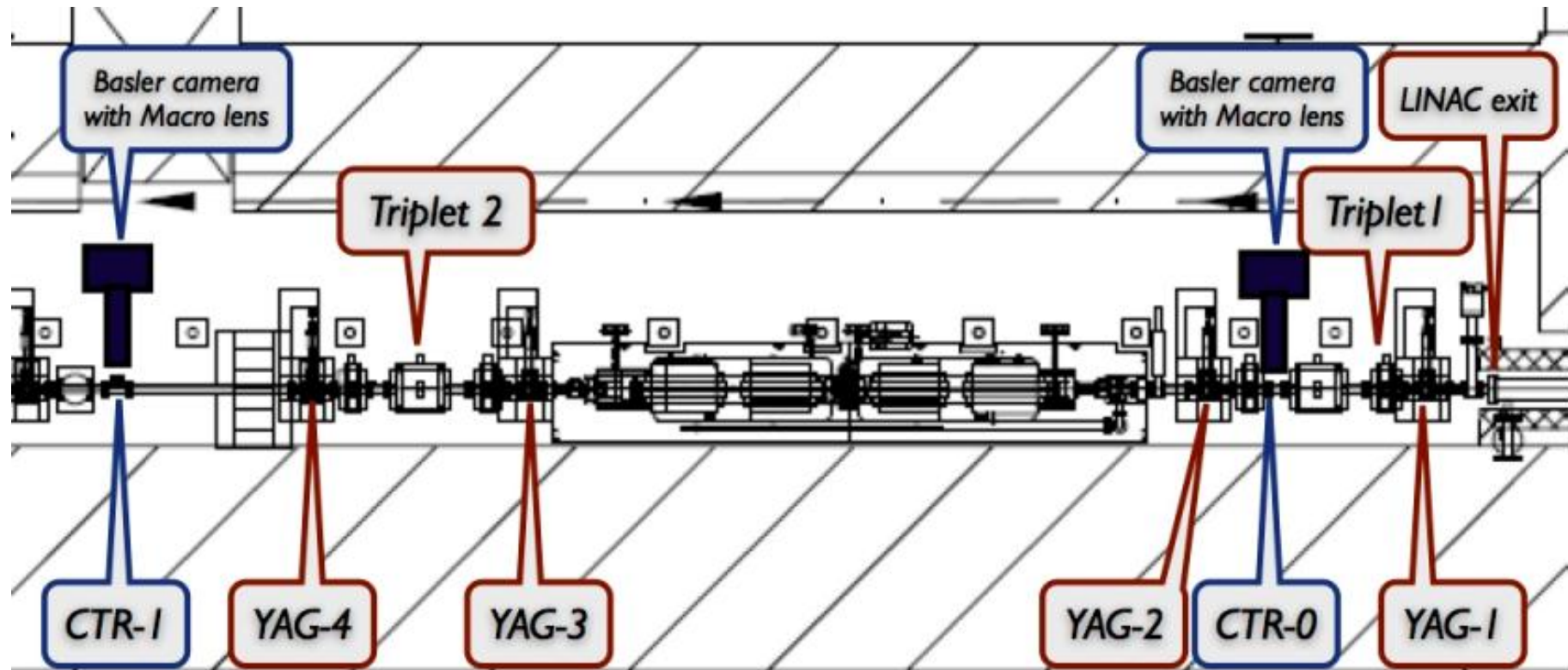
$$gain(\phi_p = \pi/2) = N^2$$

$\ll 1$ For current noise dominated beam.

For LCLS ($\delta E = 3 \text{ keV}$): $N^2 = 2.5 \times 10^{-5}$

NOISE SUPPRESSION EXPERIMENT IN ATF
ARIEL NAUSE - OCTOBER 2011

Experiment Layout



Operating Parameters

Pulse length: 5 ps

Beam energy: 50 – 70 MeV

Beam current: 40-100 A

Emittance: ~ 3 mm-mrad

Initial beam size: 400-500 μm

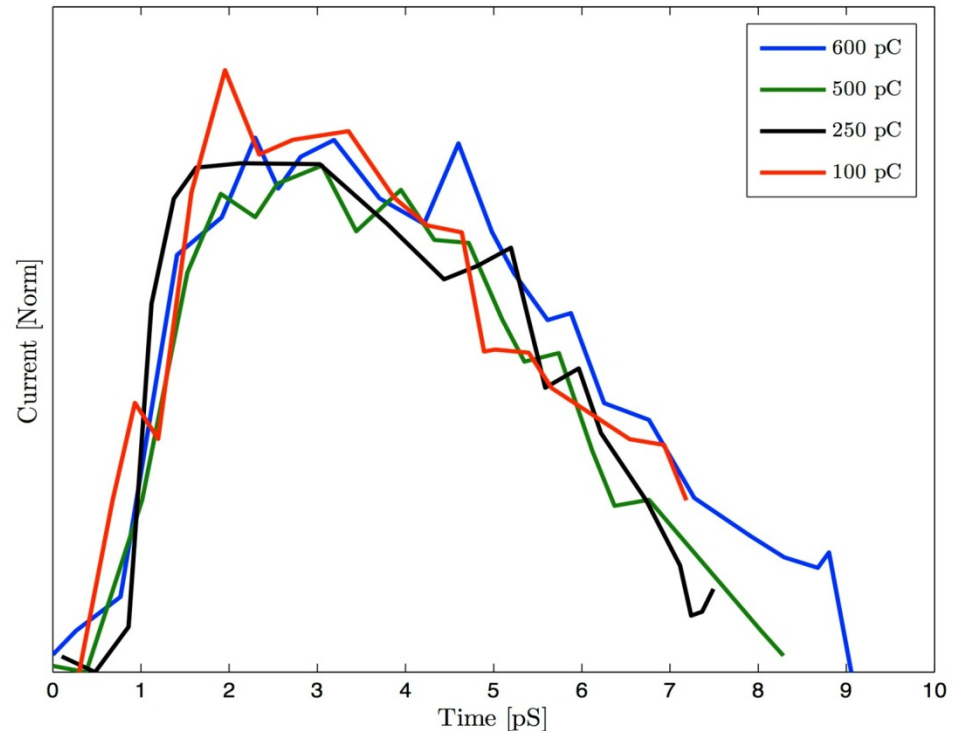
Convergence: ~ 2 mrad

Acceleration phase: on crest

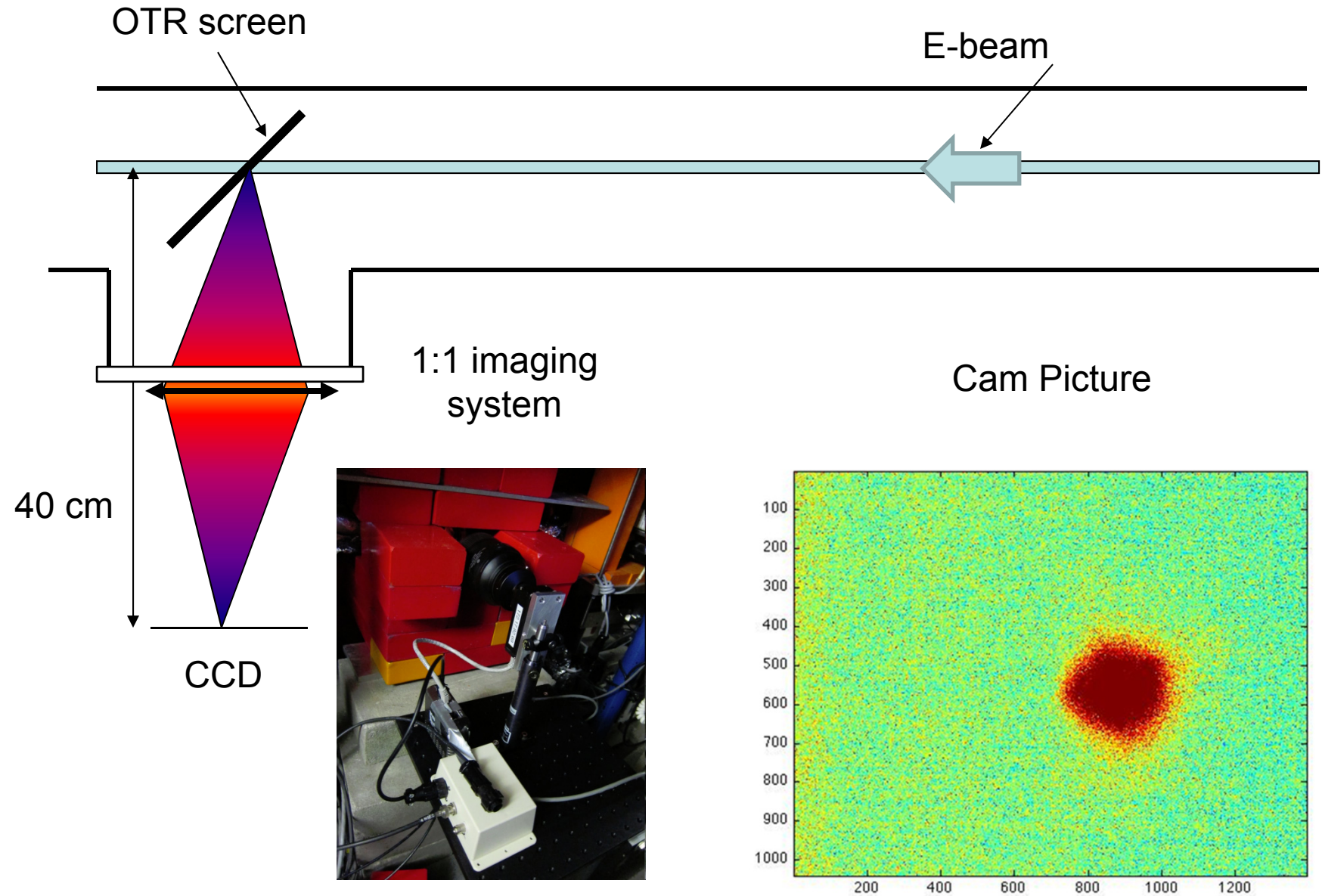
Copper OTR screen

Basler CCD camera equipped with a Nikkor macro lens (100 mm)

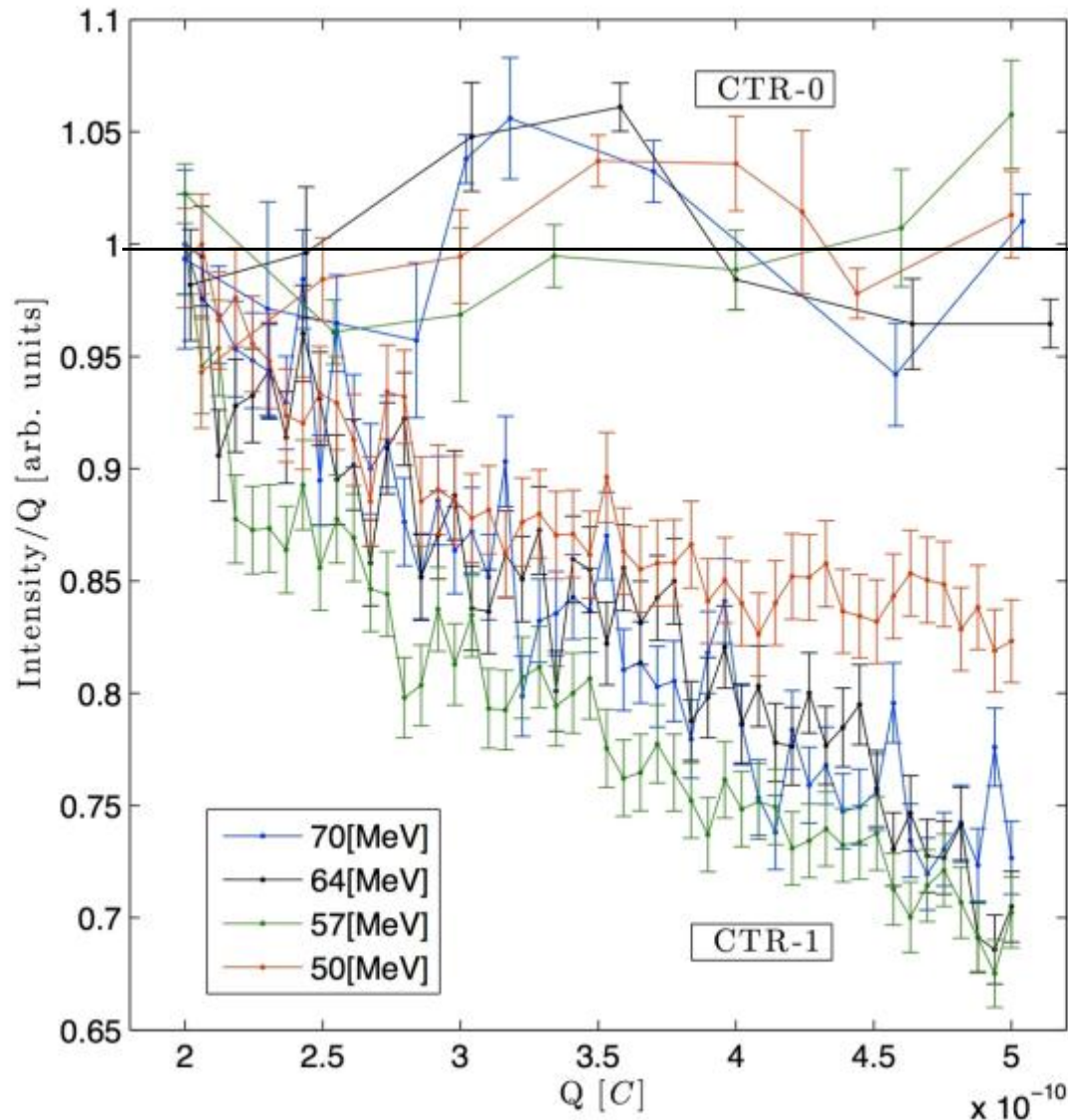
Cam sensitivity: 0.4 – 1 μm



OTR Measurement



Measured OTR Signal per unit charge



Before drift:
linear
dependence
on Q

After drift:
sub-linear
dependence
on Q

GENERAL LINEAR COLLECTIVE MICRODYNAMICS EQUATIONS FOR A BEAM WITH VARYING PARAMETERS

$$\begin{aligned}\frac{d}{d\phi_p} \check{i}(z, \omega) &= -\frac{i}{W(z)} \check{V}(z, \omega) \\ \frac{d}{d\phi_p} \check{V}(z, \omega) &= -iW(z) \check{i}(z, \omega)\end{aligned}$$

$$W(z) = \frac{r_p \sqrt{\mu_0 / \epsilon_0}}{k A_e(z) \theta_p(z)}$$

Plasma-wave “transmission line”
impedance

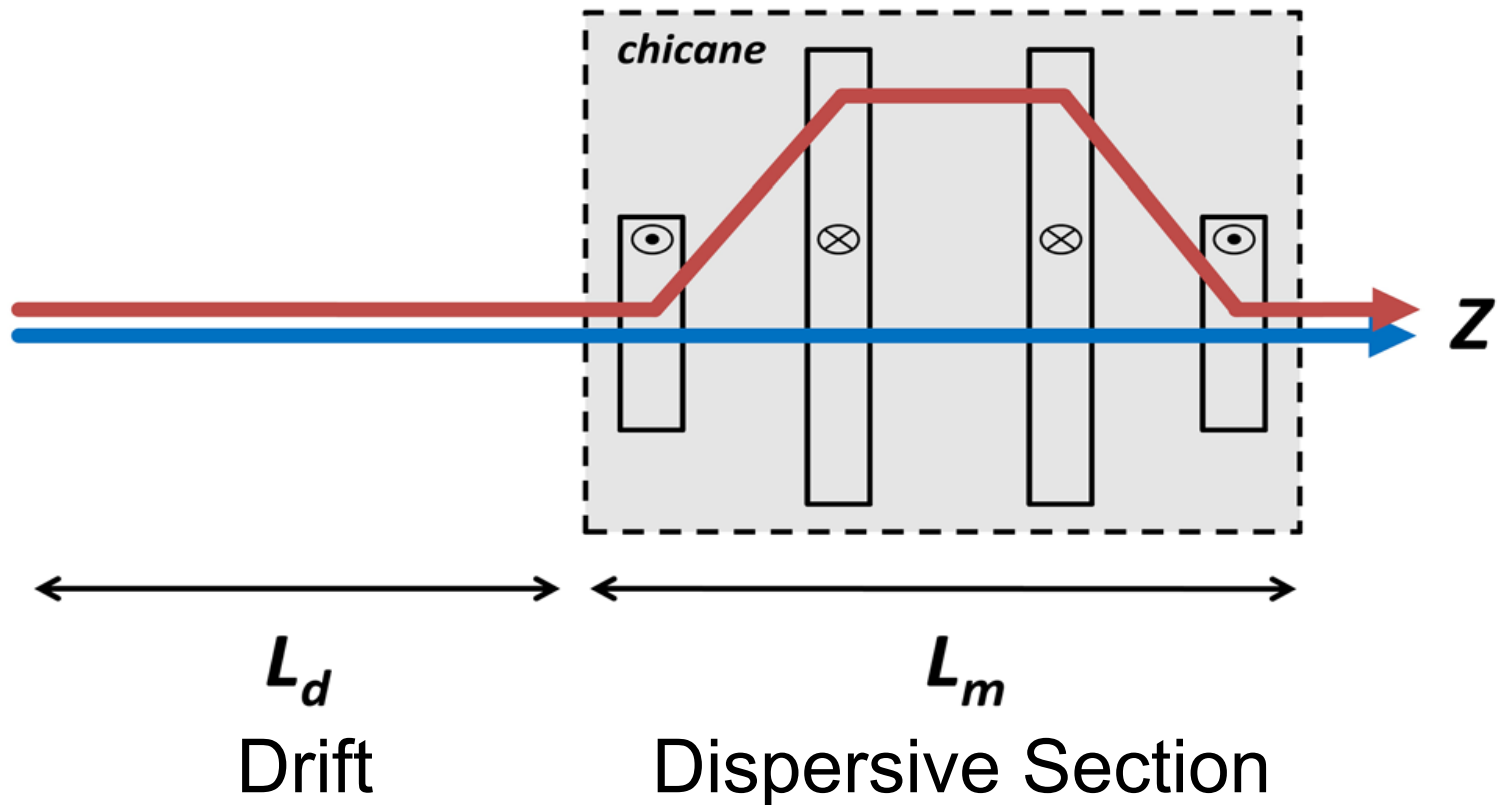
$$W = -i \frac{Z_{LSC}}{r_p \theta_p}$$

Z_{LSC} is the “conventional” beam impedance
per unit length

$$\phi_p(z) = \int_0^z \theta_{pr}(z') dz'$$

$$\theta_p^2(z) = \frac{e Z_0 I_0}{m c^2 A_e(z) \gamma_0 \gamma_{0z}^2(z) \beta_{0z}^3(z)}$$

DRIFT/DISPERSION TRANSPORT



D. Ratner Z. Huang G. Stupakov, Phys. Rev. ST-AB, **14**, 060710 (2011)
A.Gover, E.Dyunin, T.Duchovni, A.Nause, *Phys. of Plasmas*, **18**, 123102 (2011).
Experiment (35% suppression):
D. Ratner, G. Stupakov, Phys. Rev. Lett. 109, 034801 (2012)

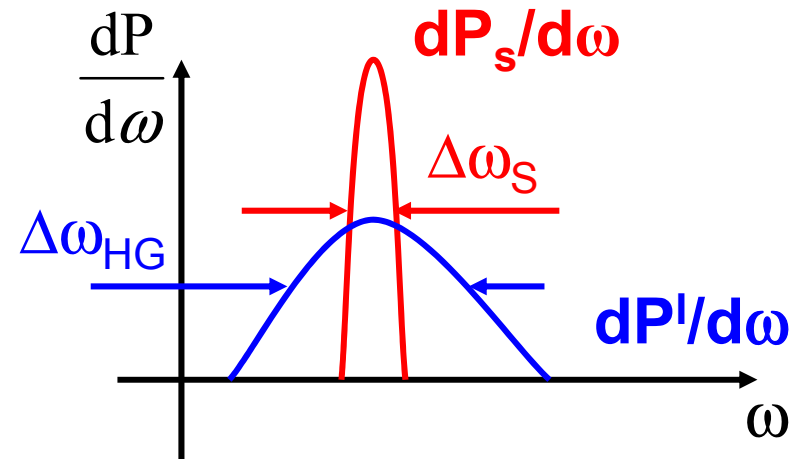
SASE SUPPRESSION

Effective Input noise (NEP)

$$\left(\frac{dP_{in}^{noise}}{d\omega}\right)_{eff} = \left(\frac{dP(L_w)}{d\omega}\right)_{incoh} / G(\omega)$$

Coherence condition

$$[P_s(0)]_{coh} \gg \left(\frac{dP_{in}^{noise}}{d\omega}\right)_{eff} \Delta\omega$$



To dominate Current Shot-Noise:

$$[P_s(0)]_{coh} \gg \frac{eI_b Z_0}{16\pi A_{em}} \left(\frac{a_w}{\gamma\beta_z \Gamma}\right)^2 \Delta\omega$$

(seed radiation injection)

$$|\tilde{i}_s(0)|^2 \gg eI_b \Delta\omega$$

(pre-bunching)

Conditions for Radiation NOISE suppression

A. Gover, E. Dyunin, "Coherence Limits of Free Electron Lasers"

IEEE J. Quant. Electron. **46**, 1511 (2010)

The diagram illustrates the components of the SASE suppression equation. Two green boxes at the top represent the 'Residual Velocity Noise Contribution' and the 'Residual Current Noise Contribution'. Green arrows point from these boxes to the terms N^2 and $\left(\frac{S}{2}\right)^2$ respectively, within a red-bordered equation box. The equation shows the ratio of the SASE power spectrum at a specific phase to its value at zero phase spread, equal to a sum of noise terms multiplied by the same ratio at zero phase spread.

$$\left. \frac{dP_{SASE}}{d\omega} \right|_{\phi_{pd} = \frac{\pi}{2} - \frac{\sqrt{3}}{2}S} = \left(N^2 + \left(\frac{S}{2} \right)^2 \right) \left. \frac{dP_{SASE}}{d\omega} \right|_{L_d=0}$$

SASE suppression factors:

$$S^2 = \left(\frac{\gamma_0}{\gamma_z} \frac{\theta_{pw}}{\Gamma} \right)^2 \ll 1 \quad N^2 = \left(\frac{\lambda_D}{\lambda} \right)^2 \ll 1$$

$S^2 \gg N^2$: SASE suppression limited by current shot-noise

$S^2 \ll N^2$: SASE suppression limited by velocity spread

Short wavelengths limits

For significant suppression

(and negligible Landau damping):

Ballistic condition

(same as Landau for $L_d = \pi/2\theta_p$):

$$N = \frac{\lambda_D}{\lambda} = k \frac{\Delta\beta_z}{\theta_p} \ll 1$$

$$\Delta\phi_p = kL_d\Delta\beta_z \ll 1$$


$$\left\{ \begin{array}{ll} \frac{k}{\theta_p} \frac{\Delta\gamma}{\gamma^3} \ll 1 & \text{✓} \\ \frac{k}{\theta_p} \left(\frac{\varepsilon_n}{\gamma\sigma_x} \right)^2 \ll 1 & \text{?} \end{array} \right.$$

Granularity condition:

$$n_0 A_e \lambda = \frac{I_0}{ec} \lambda \sim 10^4 \gg 1 \quad \text{✓}$$

Short wavelengths limit?

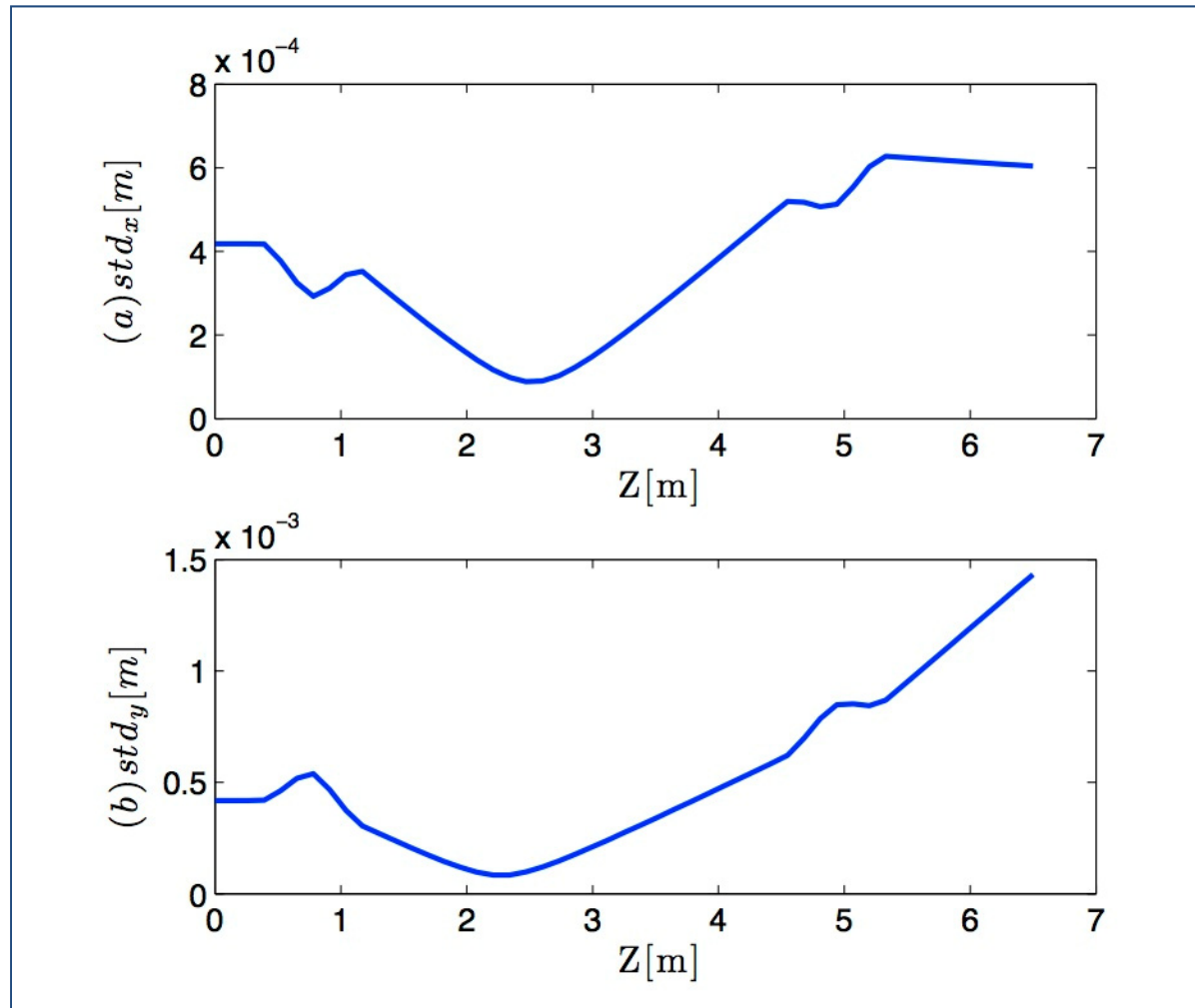
- First demonstrate in UV
- Improve beam parameters (emittance, energy spread) to satisfy $N \ll 1$
- Find optimal transport parameters.

CONCLUSION

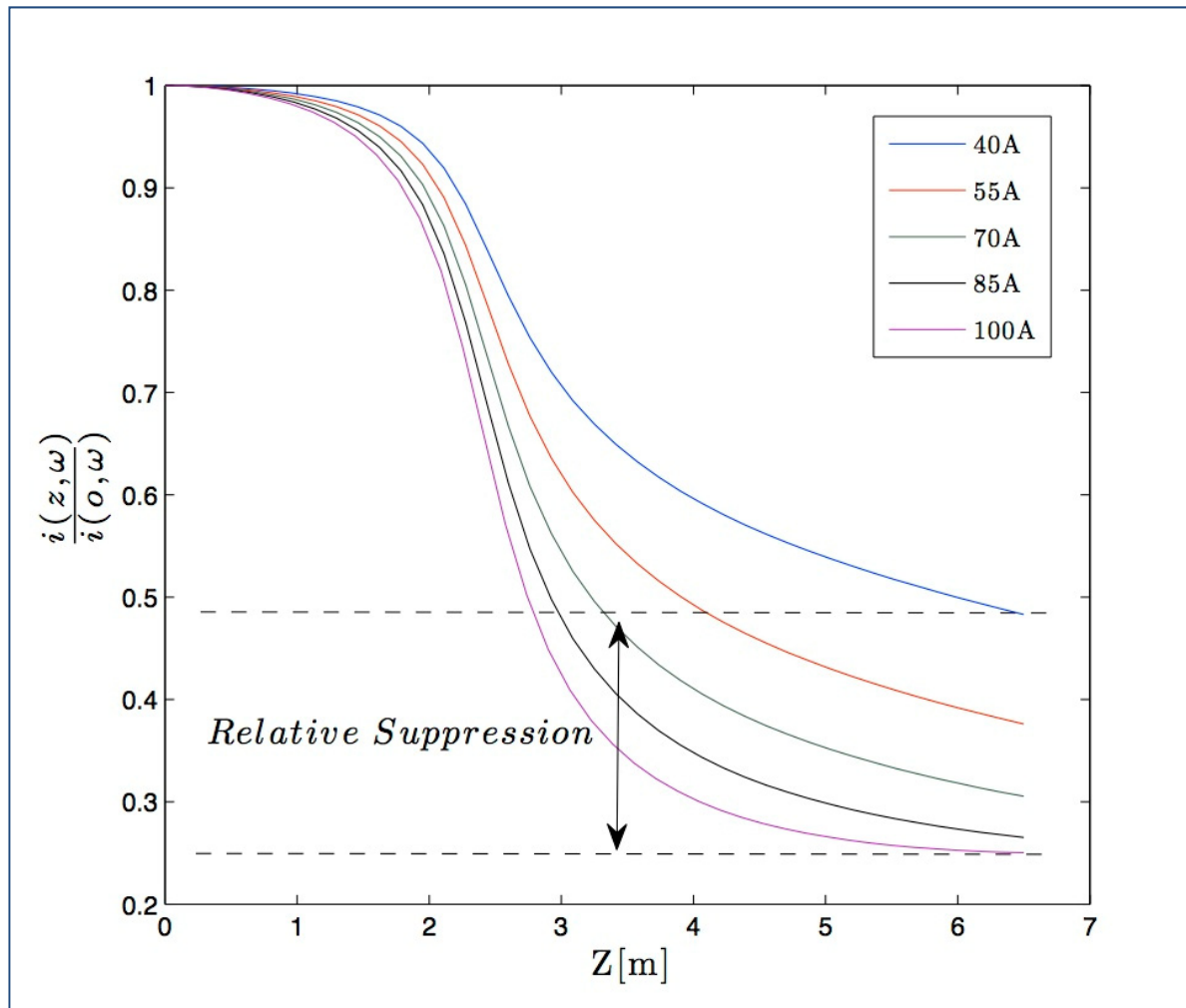
- It is possible to adjust the e-beam current shot-noise level by controlling the longitudinal plasma oscillation dynamics.
- We have demonstrated for the first time such noise suppression at optical frequencies.
- This can be used to enhance FEL coherence and relax seeding power requirement. Further studies will determine short wavelength limit needs
- After elimination of shot noise, IR/XUV FEL coherence is ultimately limited by the quantum input noise $dP / d\omega = \hbar \omega$.

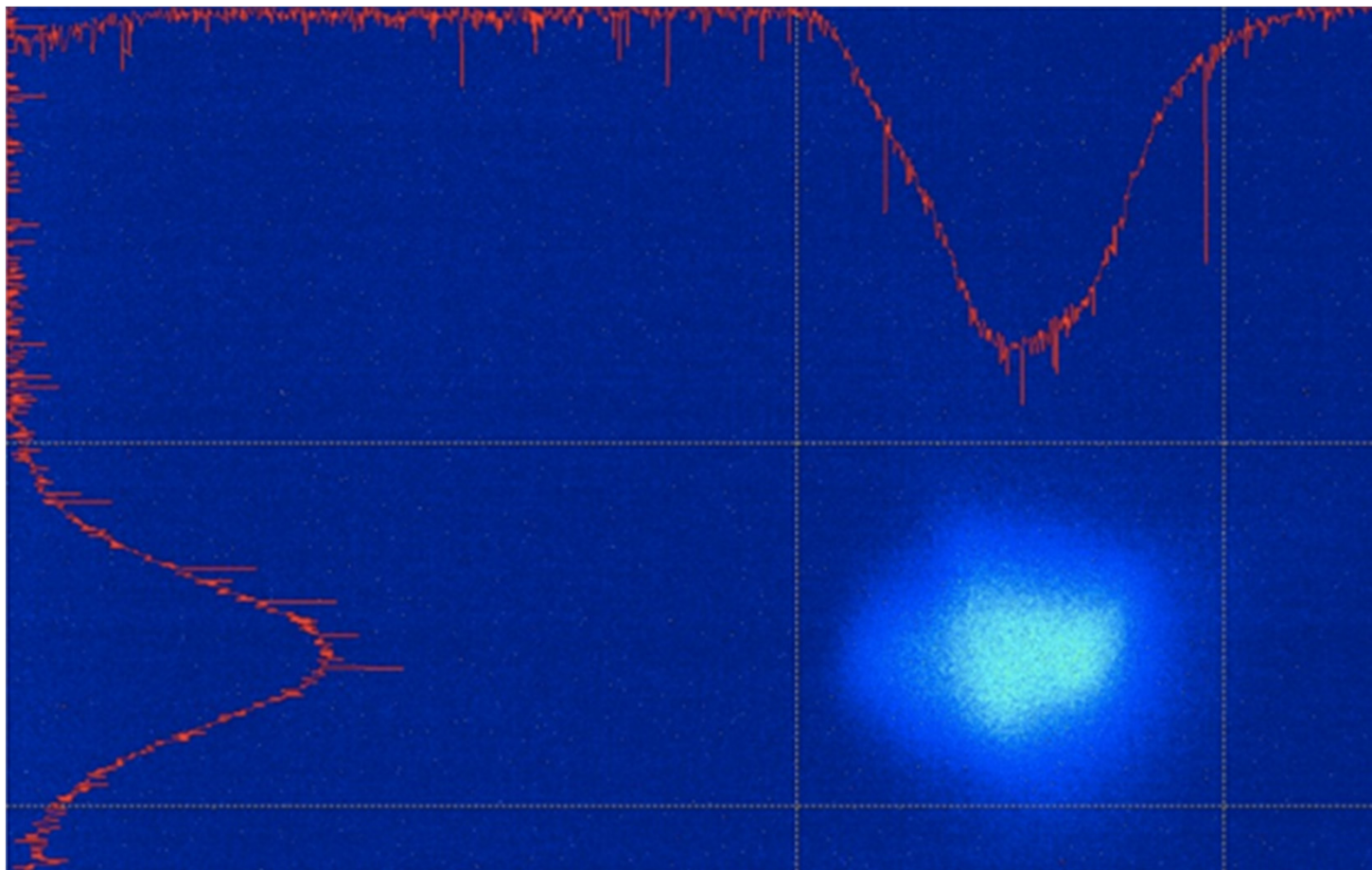
Reserve

Beam Profile Along Trajectory (GPT)



COMPUTATION OF NOISE SUPPRESSION WITH BEAM ANGULAR SPREAD





Fundamental “Schawlow-Townes” Coherence Limits (NEP)

e-Beam current noise
+ energy shot
+ radiation noise:

$$\left(\frac{dP_{in}}{d\omega} \right)_{\min} = A \bullet eI_b + B \bullet \cancel{\delta E_c} + \frac{\cancel{\hbar \omega}}{1 + \cancel{e^{-\hbar \omega / kT}}}$$

Minimum (energy spread limited) e-beam noise:

$$\left(\frac{dP_{in}}{d\omega} \right)_{\min} = \frac{\delta E_c}{\pi} + \frac{\hbar \omega}{1 + e^{-\hbar \omega / kT}}$$

Microwave/THz regime:

$$\left(\frac{dP_{in}}{d\omega} \right)_{\min} = \frac{\delta E_c}{\pi} + \cancel{k_B T} \quad \left(\approx \frac{\delta E_c}{\pi} = \frac{k_B T_c}{\pi} > k_B T \right)$$

(Cathode temperature limited)

Optical/X-UV regime:

$$\left(\frac{dP_{in}}{d\omega} \right)_{\min} = \frac{\cancel{\delta E_c}}{\pi} + \hbar \omega \quad (\approx \hbar \omega)$$

(Quantum limit)

3-D Numerical Simulations

**A. Nause, E. Dyunin, A. Gover,
JAP 107, 103101 (2010).**

100,000 particles over 2 pS duration to increase resolution,
70 MeV energy, 200pC charge

E-BEAM NOISE AND RADIATION SUPPRESSION THEORY

Microwave tube noise suppression

H. Haus and F. N. H. Robinson, Proc. IRE 43, 981 (1955).



Optical noise suppression in a drifting relativistic beam:

Gover, Phys. Rev. Lett. 102, 154801 (2009),

Nause, JAP, 107, 103101 (2010)

Optical noise suppression with a dispersive section:

Rathner, PhysRevSTAB 14 060710 (2011)

Gover, Phys. Plasmas 18, 123102 (2011)

SASE noise suppression:

Gover, JQE46, 1511 (2010)

Short wavelength limit:

R. Bonifacio, Optics Communications 138 (1997) 99-100

K-J Kim, Shanghai FEL conférence 2011

3-D Homogenization Trend

A simple physical argument:

Inter-particle Coulomb force: $\epsilon_{\text{Coul}} = e^2 / 4\pi\epsilon_0 n_0'^{-1/3}$

Space-charge force: $\epsilon_{sc} = e^2 \Delta N' / 2\pi\epsilon_0 d'$

Poisson statistics: $\Delta N' = N'^{1/2}$ $N' = (\pi d'^3 n_0' / 6)^{1/2}$

When $\epsilon_{sc} > \epsilon_{\text{Coul}}$?

$$\frac{\epsilon_{sc}}{\epsilon_{\text{Coul}}} = \left(\frac{2\pi}{3} \frac{d'}{n_0'^{-1/3}} \right)^{1/2} > 1$$

Answer: $d' > n_0'^{-1/3}$

Note: Process leads to velocity spread growth

At the Cathode (no-correlation point)

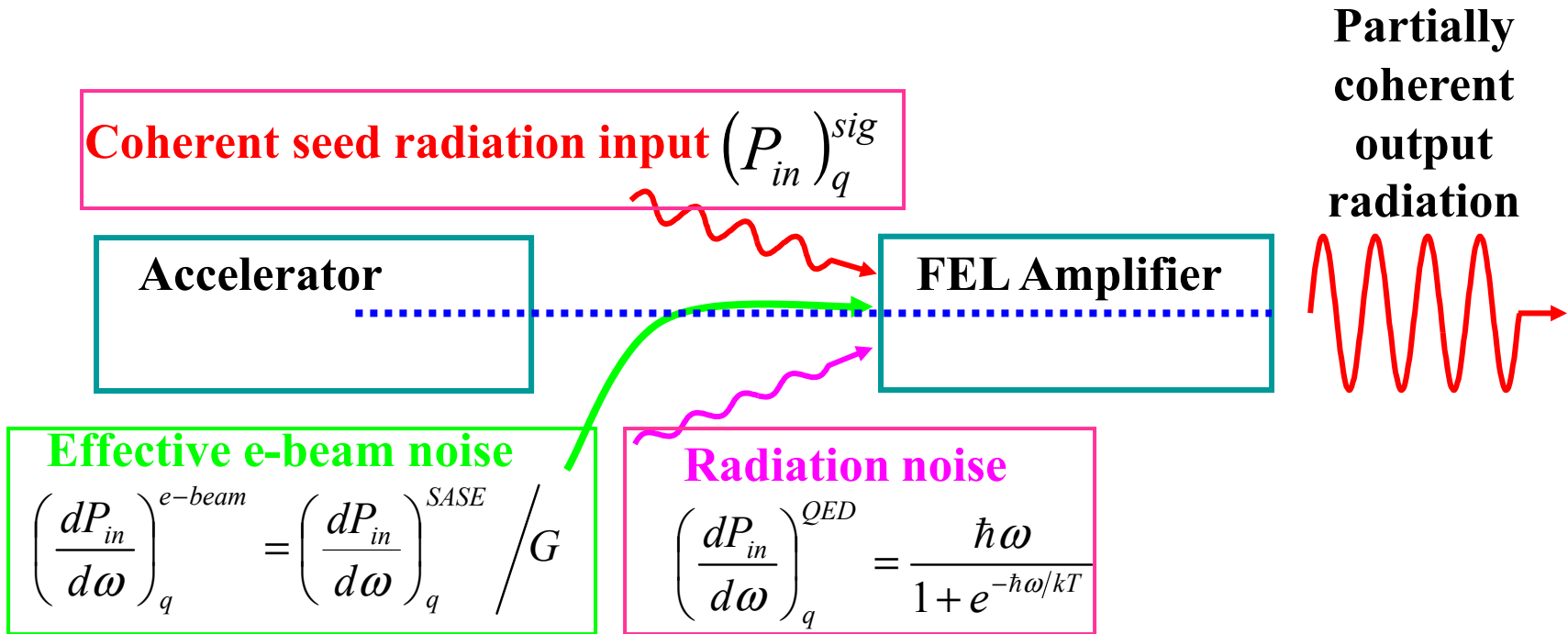
Current and velocity noise – uncorrelated:

$$\overline{|\check{i}(\omega)|^2} = \frac{1}{T} \left\langle |\check{i}(\omega)|^2 \right\rangle_{N_T} = eI_b$$

$$\overline{|\check{v}(\omega)|^2} = \frac{1}{T} \left\langle |\check{v}(\omega)|^2 \right\rangle_{N_T} = \frac{(\delta E_c)^2}{eI_b}$$

$$\left(\overline{|\check{I}(\omega)|^2} \right)^{1/2} \left(\overline{|\check{V}(\omega)|^2} \right)^{1/2} = \delta E_c$$

NOISE INPUTS INTO FEL AMPLIFIER (NOISE EQUIVALENT POWER - NEP)



Fundamental Coherence Limits

A Conservative of motion in a non-dissipative e-beam transport section:

$$\overline{|\check{i}_c|^2} \overline{|\check{V}_c|^2} = (\delta E_c)^2$$

Minimum input noise:

$$\left(\frac{dP_{in}}{d\omega} \right)_{\min} = \frac{\delta E_c}{\pi} + \frac{1}{1 + e^{-\hbar\omega/kT}}$$

Microwave/THz regime:

$$\left(\frac{dP_{in}}{d\omega} \right)_{\min} = \frac{\delta E_c}{\pi} + k_B T \quad \left(\approx \frac{\delta E_c}{\pi} \geq \frac{k_B T_c}{\pi} \right)$$

(Cathode temperature limited)

Optical regime:

$$\left(\frac{dP_{in}}{d\omega} \right)_{\min} = \frac{\delta E_c}{\pi} + \hbar\omega \quad (\approx \hbar\omega)$$

(Quantum limit)