

Longitudinal Shaping of Relativistic Electron Bunches with Applications to the Plasma Wakefield Accelerator

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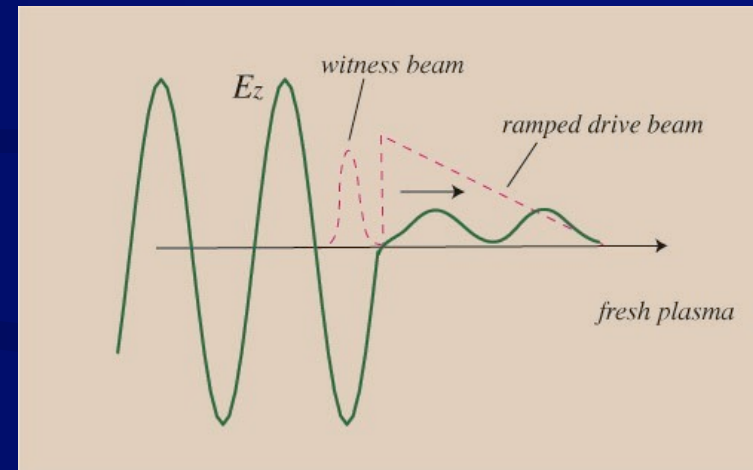
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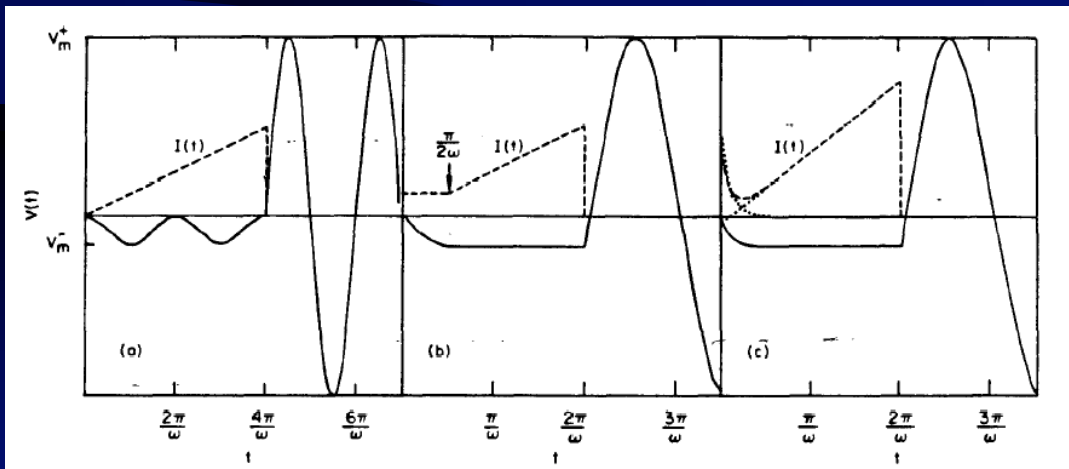
*2009 Particle Accelerator Conference
May 7, 2009*

Optimal Drive Beam Profile for Blowout Regime of PWFA

- PWFA: plasma wakefield accelerator
- electron beam-driven plasma waves
- acc. fields on order of multi-GeV/m
- acceleration of drive tail or witness bunch



Chen, P., Su, J., and Dawson, J. SLAC PUB 3662 (1985).



triangle

doorstep

ideal

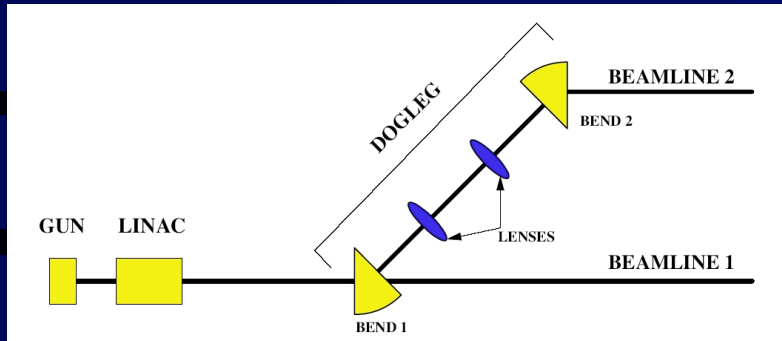
Transformer Ratio:

$$E_+ = \text{acc. field}; E_- = \text{decc. field}$$

$$R = E_+ / E_- = k_p L_z$$

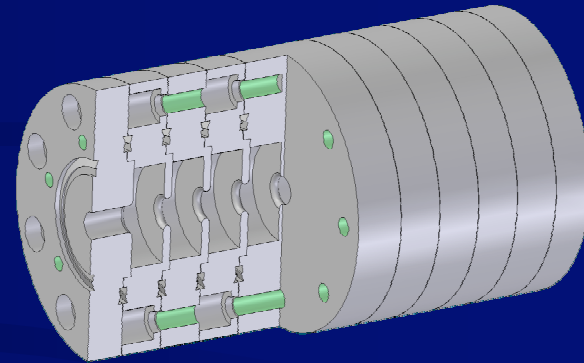
$$R > 2 \text{ if } L_z > 2k_p$$

Focus of this Talk



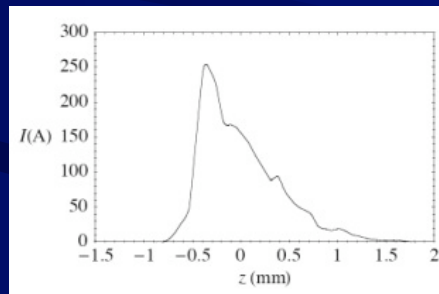
dogleg compressor

+



temporal diagnostic

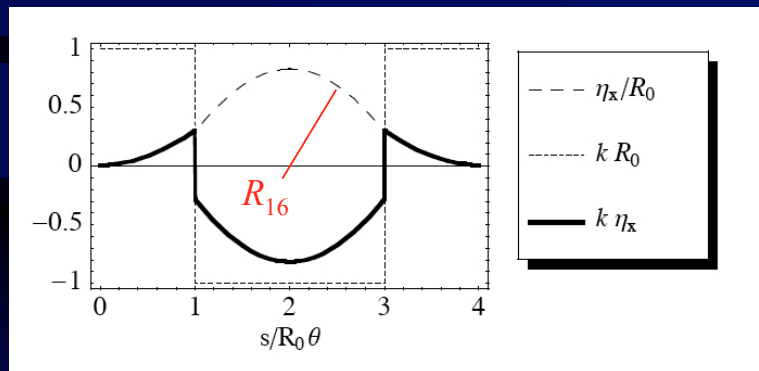
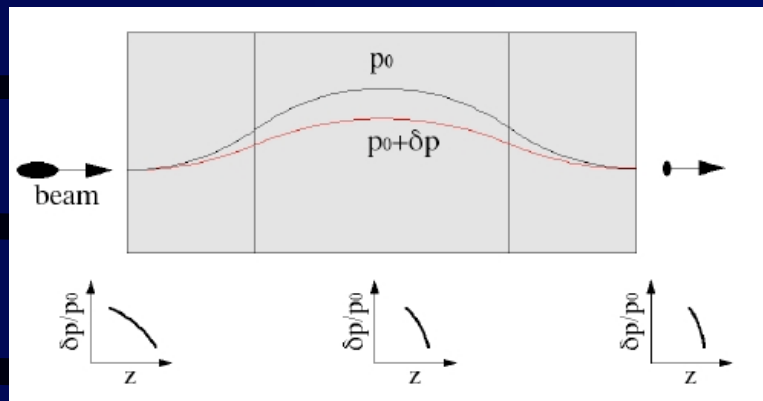
=



- Generation of electron beam with ramped current profile
- Temporal diagnostic with sub-ps resolution
 - transverse deflecting mode cavity
- Experimental verification of ramping mechanism

How Does a Dogleg Compress the Beam?

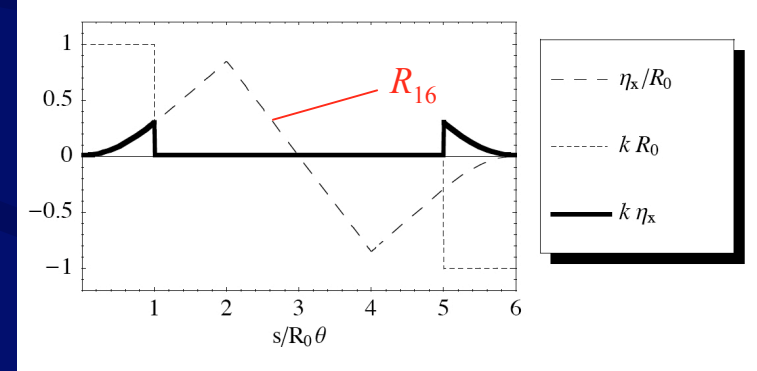
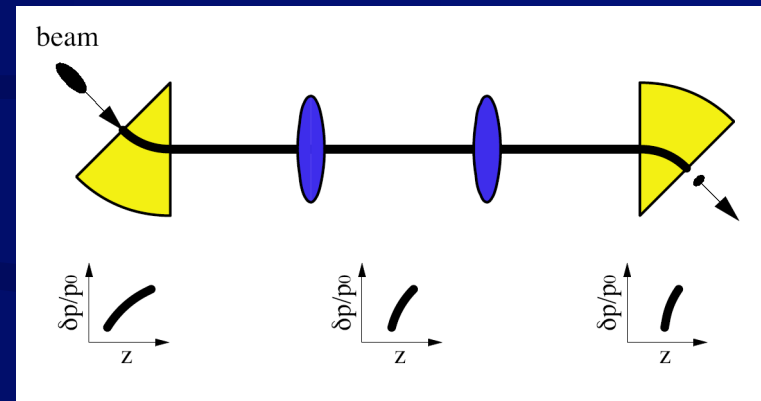
chicane



higher-energy particles travel a shorter path

$$R_{56} = \frac{\partial z}{\partial \delta} > 0 \quad \text{"positive longitudinal dispersion"}$$

dogleg

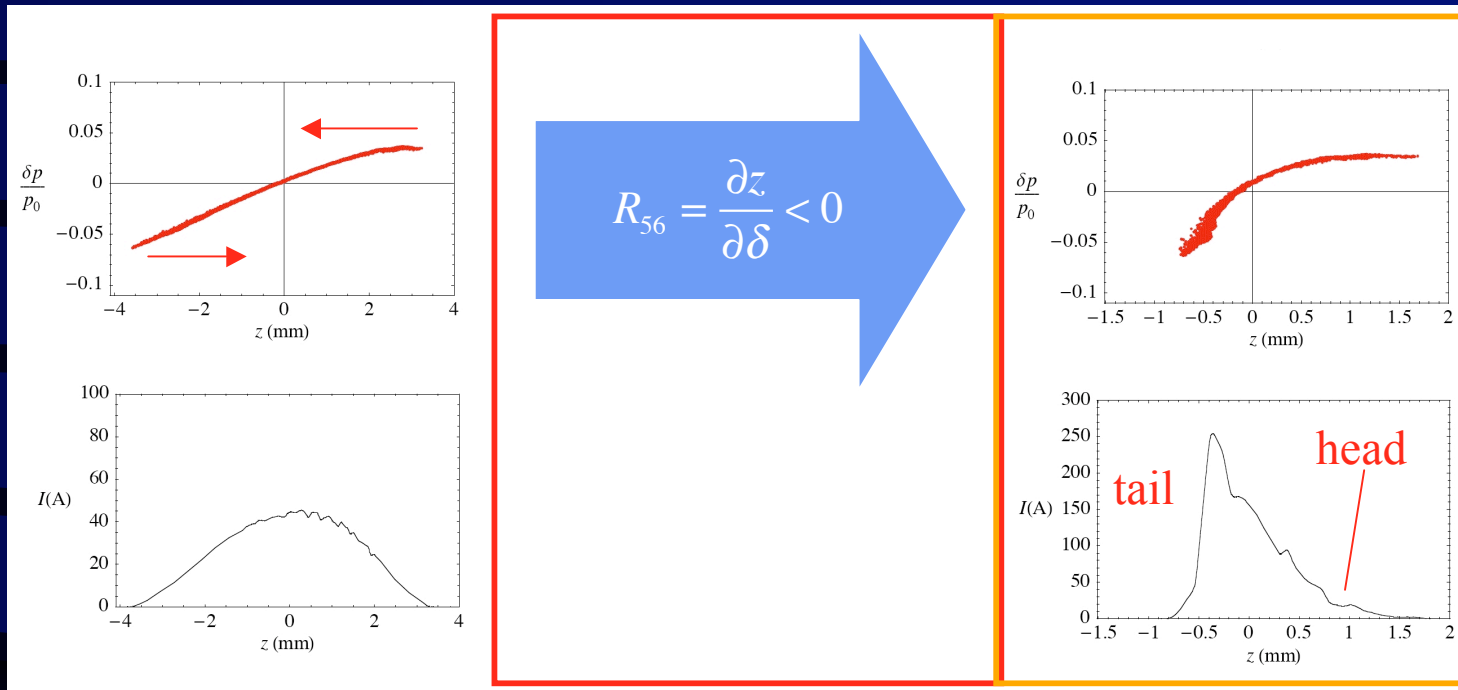


higher-energy particles travel a longer path

$$R_{56} = \frac{\partial z}{\partial \delta} < 0 \quad \text{"negative longitudinal dispersion"}$$

Ramped Beam Mechanism

Artificial mathematical manipulation of a chirped particle distribution



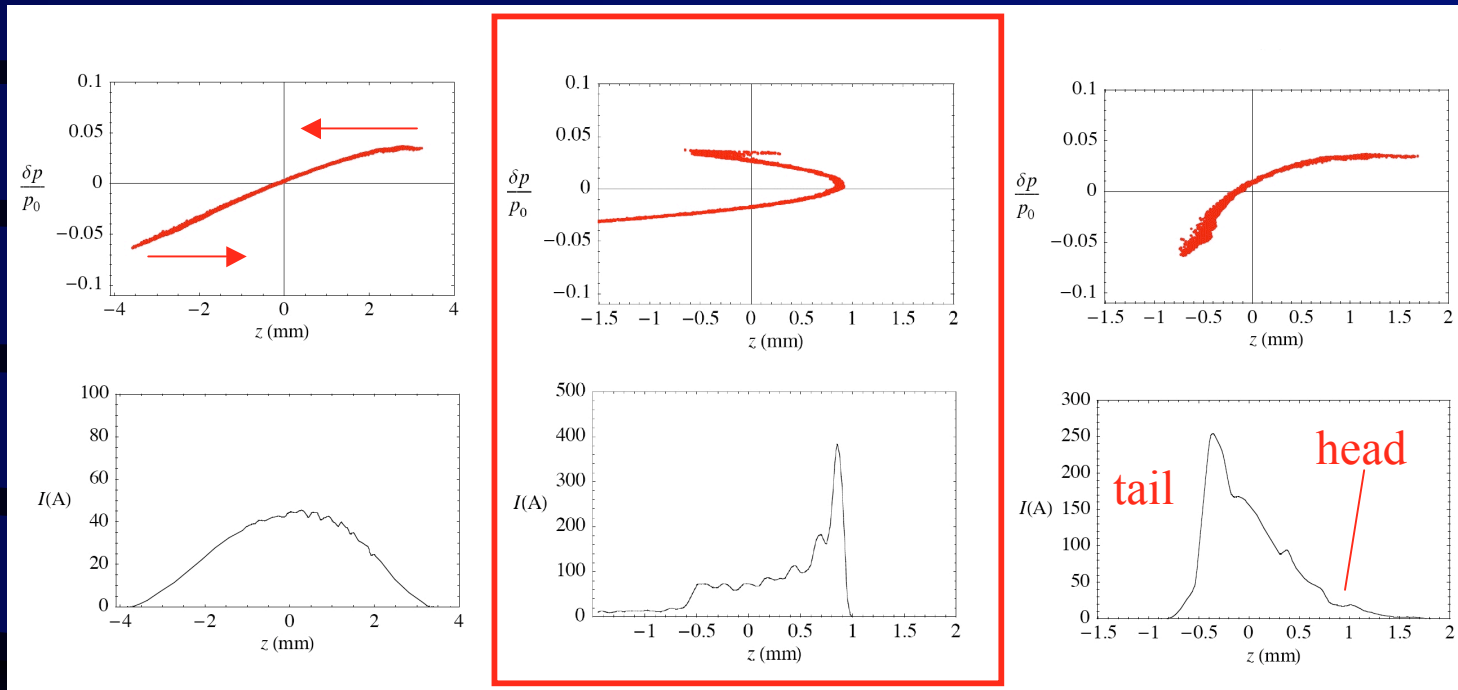
linear transformation: $z_f = z_0 + R_{56}\delta$

T_{566} arises from chromatic focusing errors in horizontally focusing quads and then grows in the subsequent drift sections (2nd order x-z correlation).

Solution: sextupole corrector magnets near the horizontally focusing quads.

Ramped Beam Mechanism

Artificial mathematical manipulation of a chirped particle distribution



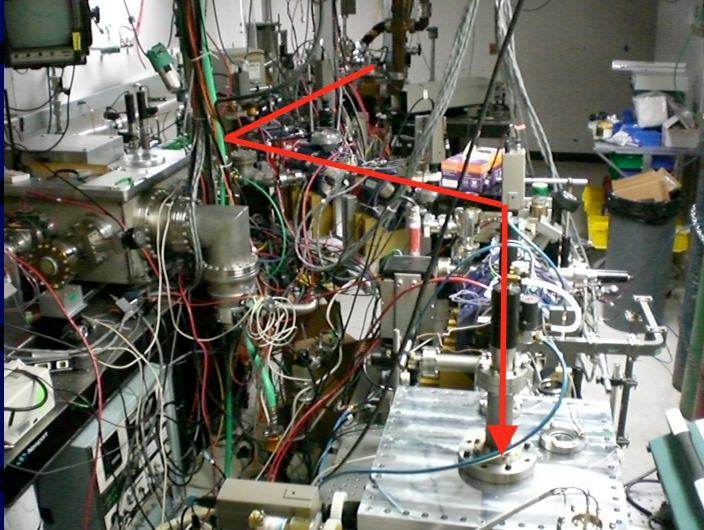
non linear transformation: $z_f = z_0 + R_{56}\delta + T_{566}\delta^2$

T_{566} arises from chromatic focusing errors in horizontally focusing quads and then grows in the subsequent drift sections (2nd order x-z correlation).

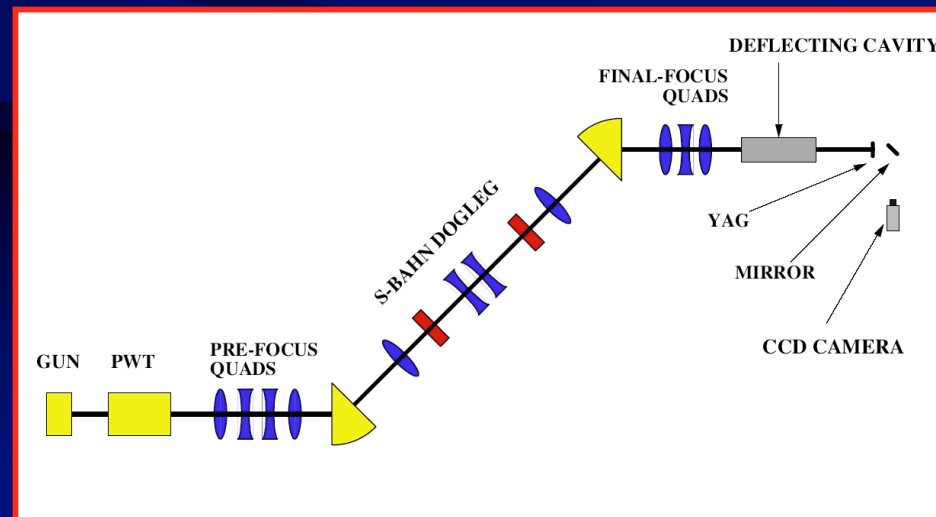
Solution: sextupole corrector magnets near the horizontally focusing quads.

UCLA Neptune Beamline

With Dogleg Compressor

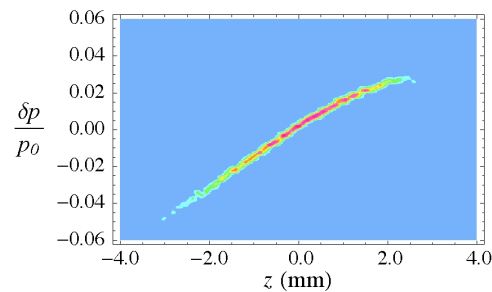


- Is a “dogleg” or dispersionless translating section.
- Half-chicane with focusing elements between the bends.
- Can be operated in a nondispersive mode with symmetric beta function and 2π betatron advance.
- Like a chicane, may be used as a bunch-length compressor.
- Nominal first order temporal dispersion ($R_{56} = -5\text{cm}$) is suitable for beam-shaping.
- Sextupoles required to compensate 2nd order longitudinal dispersion.

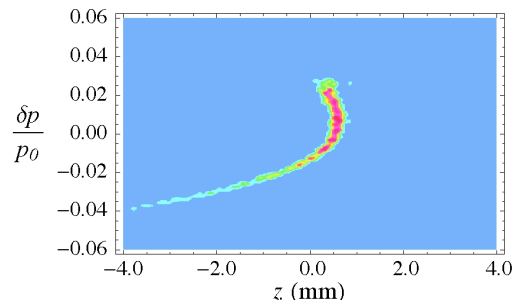


Neptune Dogleg Compressor

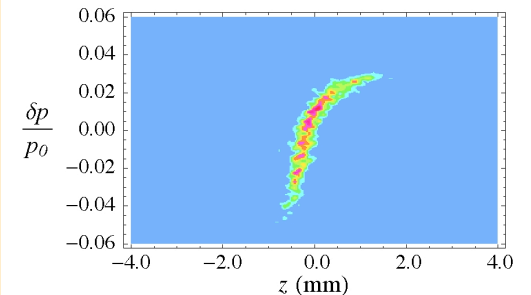
PARMELA Simulation Results: 1000 particles, 300pC



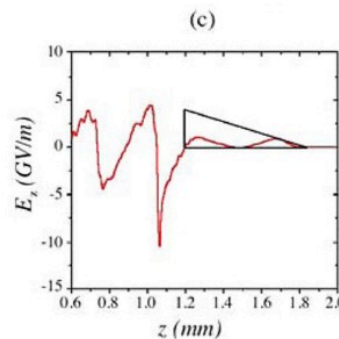
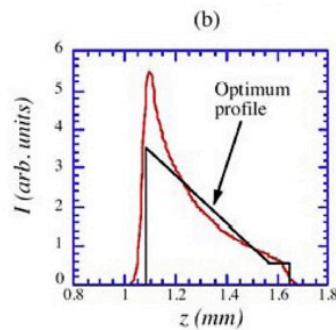
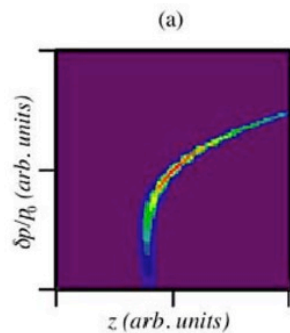
Initial



Final: Sextupoles Off



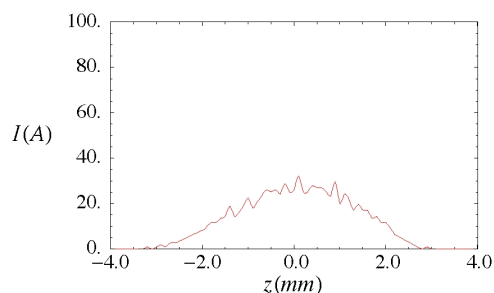
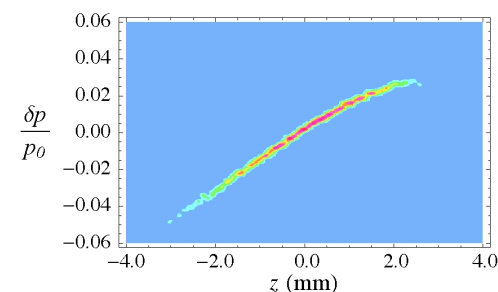
Final: Sextupoles On



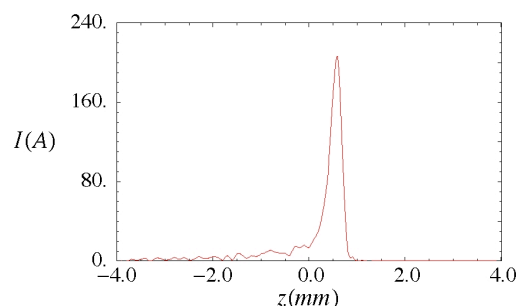
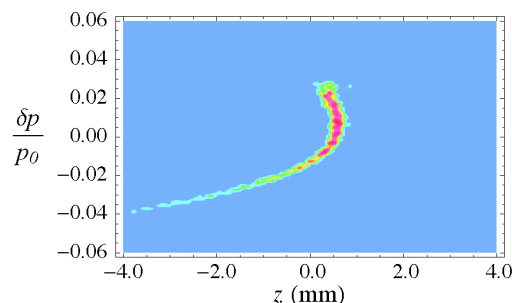
- 2D PIC Simulation
- 5 GeV/m gradients
- 6 nC drive beam w/ $n_0 = 2 \times 10^{16} \text{ cm}^{-3}$

Neptune Dogleg Compressor

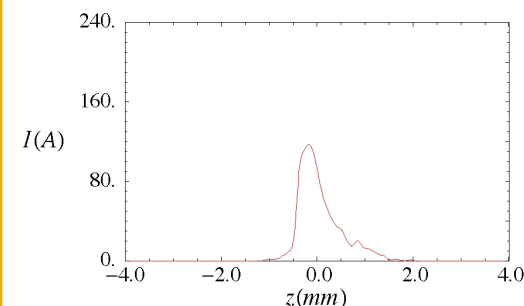
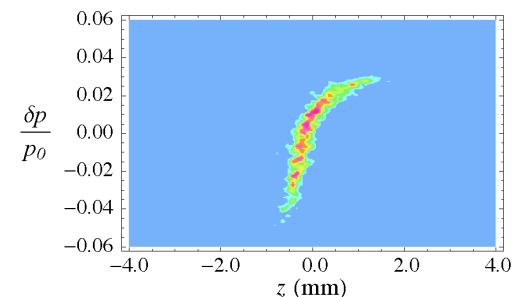
PARMELA Simulation Results: 1000 particles, 300pC



Initial



Final: Sextupoles Off

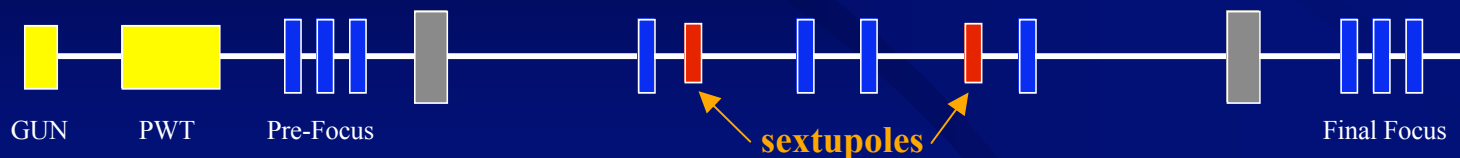


Final: Sextupoles On

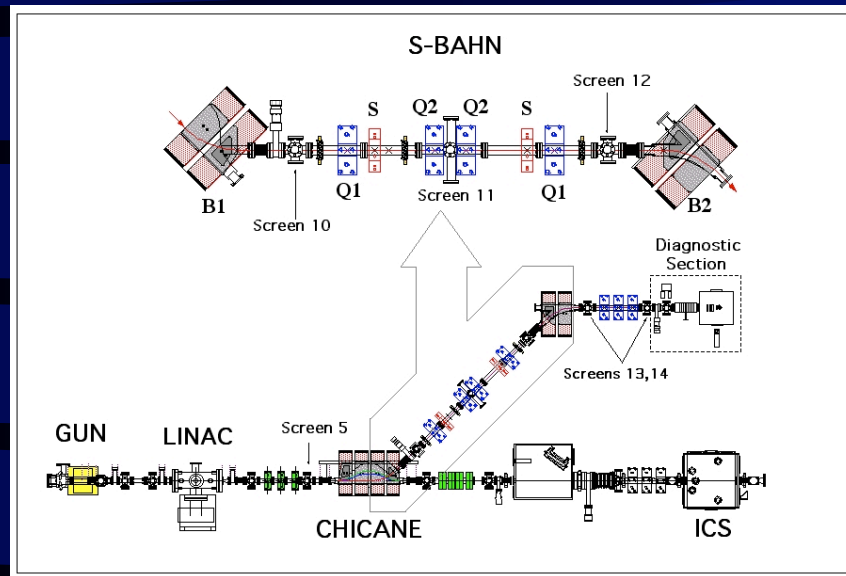
$$\epsilon_{x,N}(\text{initial}) = 4.9 \mu\text{m}$$

$$\Delta\epsilon_{x,N} = 9.9 \mu\text{m} + 12.7 \mu\text{m} = 22.6 \mu\text{m}$$

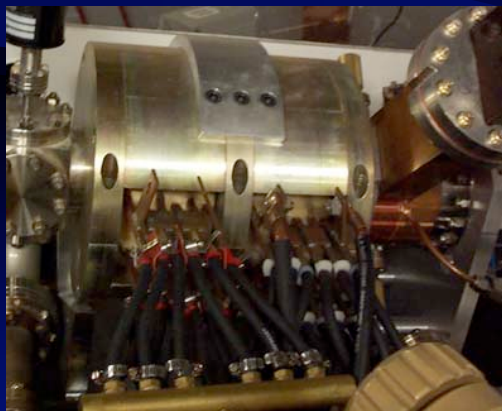
space-charge nonlinear total



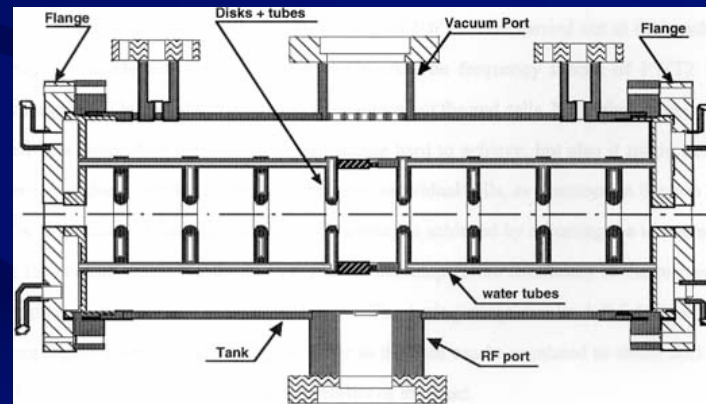
The UCLA Neptune Laboratory



Beam Charge:	100pC ---> 500pC
Beam energy:	up to 15 MeV
Emittance:	$\epsilon_N = 4$ mm mrad
Power Source:	18 MW Klystron
RF Frequency:	2.856 GHz
Cathode laser:	60 μ J at $\lambda = 266$ nm
Laser pulse length:	5-7 ps RMS

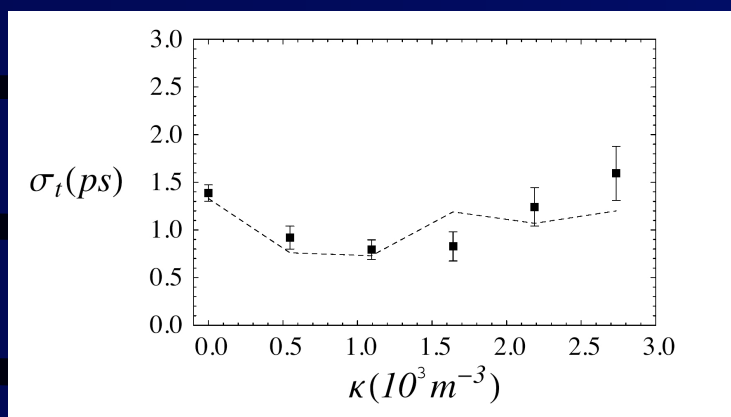


1.6-Cell Photoinjector

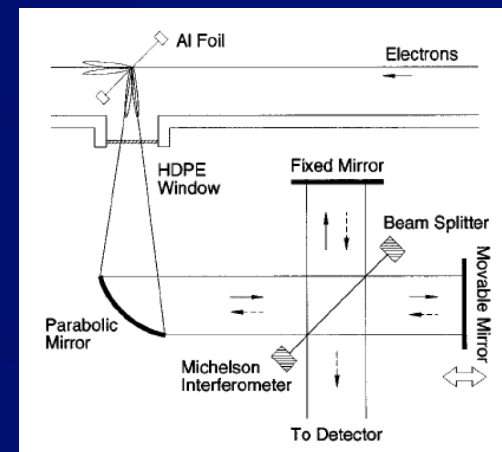


7&2/2 Cell PWT Linac

Coherent Transition Radiation Measurements of Compression



- Martin-Puplett CTR Interferometer
- Bunch length measurement by autocorrelation.
- Sub-picosecond resolution obtainable.



Simulations predict “ramped” beam occurs near point of maximum compression ($\kappa=1094 \text{ m}^{-3}$).

Empirical CTR values obtained using method of [Murokh, et al., Nucl. Instr. Meth. Phys. Res. A, **410**, 452 (1998)].

Empirical analysis assumes a gaussian profile, which is not necessarily the case here.

Theoretical curve obtained from PARMELA + ELEGANT simulation, with autocorrelation algorithm.

PARMELA
gun and linac
5000 macroparticles

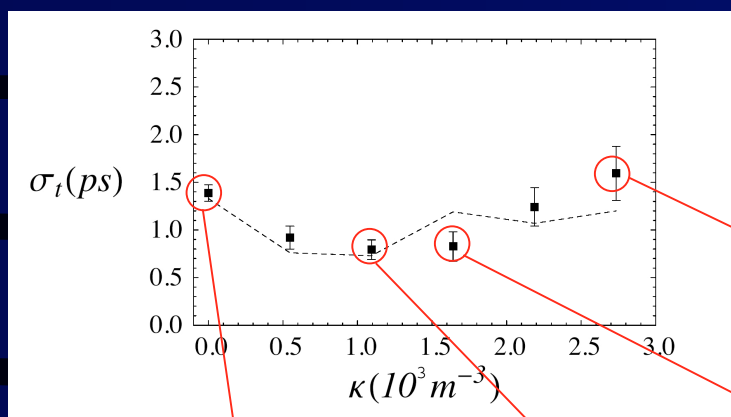


ELEGANT
prefocus and s-bahn
60% collimation

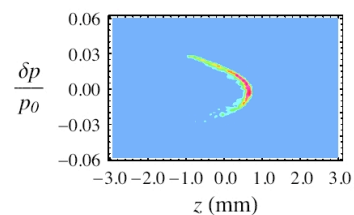
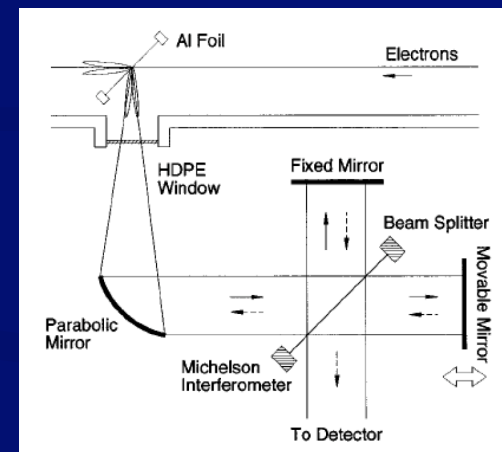


MATHEMATICA
1.interferogram reconstruction
2.triple-gaussian fit procedure

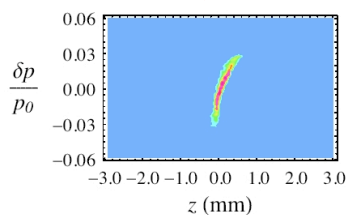
Coherent Transition Radiation Measurements of Compression



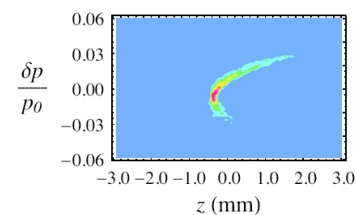
- Martin-Puplett CTR Interferometer
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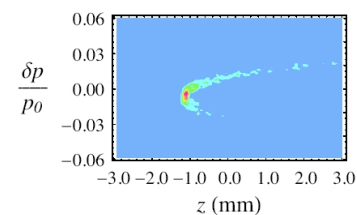
$\kappa = 0 \text{ m}^{-3}$



$\kappa = 1094 \text{ m}^{-3}$



$\kappa = 1641 \text{ m}^{-3}$

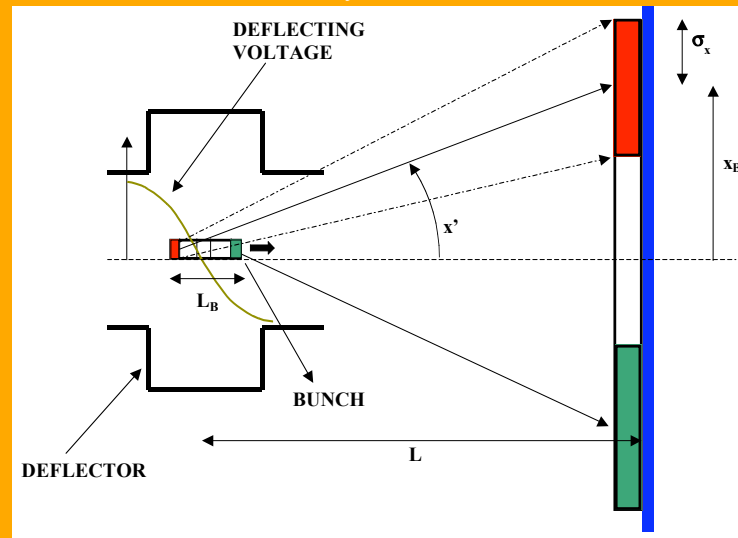


$\kappa = 2735 \text{ m}^{-3}$

A Better Temporal Diagnostic

Deflecting Mode Cavity

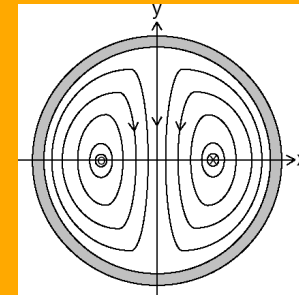
Courtesy of D. Alesini



Lowest dipole mode is TM_{110}
 Zero electric field on-axis (in pillbox approx.)
 Deflection is purely magnetic
 Polarization selection requires asymmetry

$$x' = \frac{\pi f_{RF} L_B \sqrt{2 P_{RF} R_{\perp}}}{cE/e}$$

$$x_B = \frac{\pi f_{RF} L L_B \sqrt{2 P_{RF} R_{\perp}}}{cE/e}$$



J.D. Fuerst, et. al., DESY Report CDR98, 1998

Pillbox Fields

$$E_z = E_0 J_1(\kappa r) e^{i\phi};$$

$$B_r = B_0 \frac{J_1(\kappa r)}{\kappa r} e^{i\phi};$$

$$B_\phi = i B_0 J'_1(\kappa r) e^{i\phi};$$

on axis
 $\kappa r = 0$

$$E_z = 0;$$

$$B_x = \frac{B_0}{2};$$

$$B_y = i \frac{B_0}{2};$$

Deflecting Mode Cavity

Power and Resolution

screen deflection: $\sigma_{x,f} = \sqrt{\sigma_0^2 + \sigma_{def}^2}$ $\sigma_{def} = 2\sigma_z L \frac{\pi V_{\perp} f}{cU/e}$

$$V_{\perp} \gg V_{min} = \frac{\sigma_{x,0} U/e}{L\pi\sigma_t f} \quad \sigma_{t,min} = \frac{\sigma_{x,0} U/e}{L\pi V_{\perp} f}$$

$$V_{\perp,design} = 3V_{min} = 545kV \quad \sigma_{t,min} = 545fs$$



$$\Delta t = \frac{\sigma_0 U/e}{L\pi f R_{\perp}^{1/2} \sqrt{n P_{in}}}$$

$\sigma_{x,f}$ = beam size at screen with deflector on;

$\sigma_0 = 0.3mm$ = beam size at screen with deflector off;

$L = 43cm$ = drift from deflector to screen;

$f = 9.6GHz$ = RF frequency;

$V_{\perp}$ = deflecting voltage;

$R_{\perp} = 820k\Omega$ = transverse shunt impedance per cell;

P_{in} = input RF power;

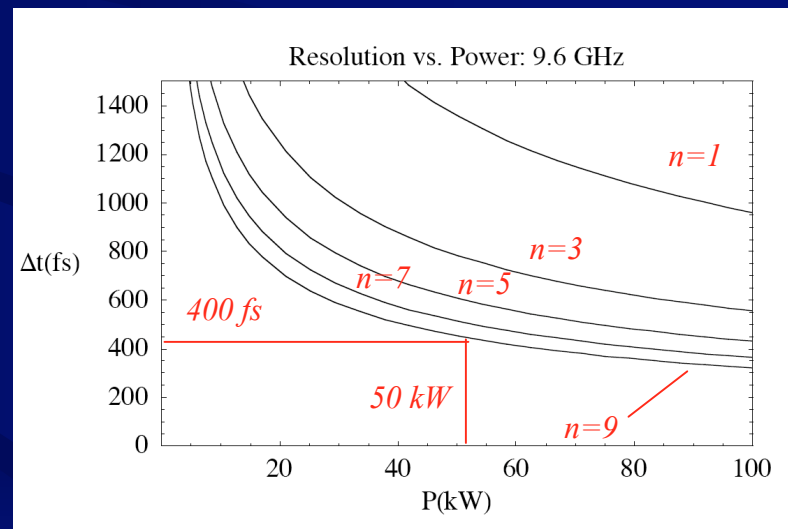
$U = 12MeV$ = electron beam energy;

ϕ_0 = deflector injection phase = 0;

$\sigma_{t,min}$ = minimum resolvable rms bunch length;

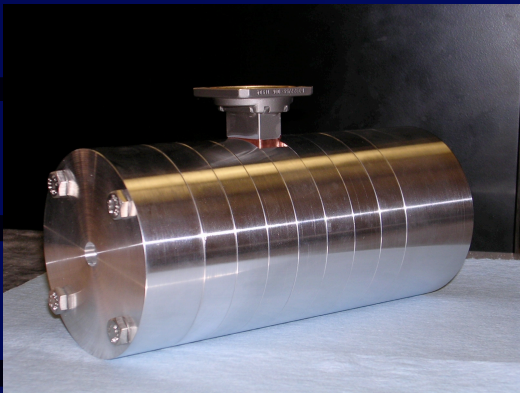
$\Delta x = 30\mu m$ = spatial resolution of screen & optics;

Δt = effective temporal resolution of deflector;



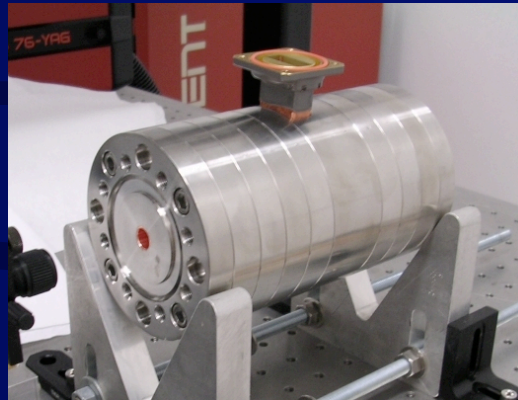
9 cells; 50 kW; 400 fs resolution

Overview of Design Process



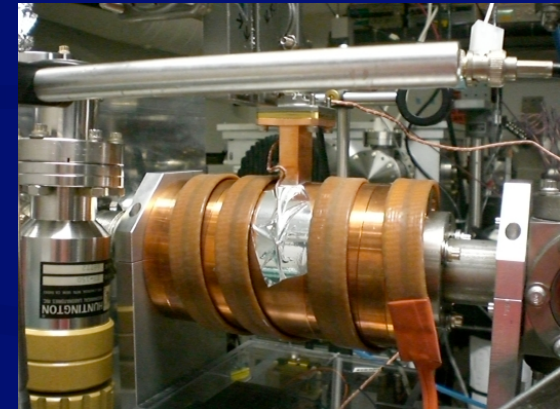
cold test prototype

Aluminum
9-cell
9.3 GHz
cold-test only
clamped
no polarization separation



steel prototype

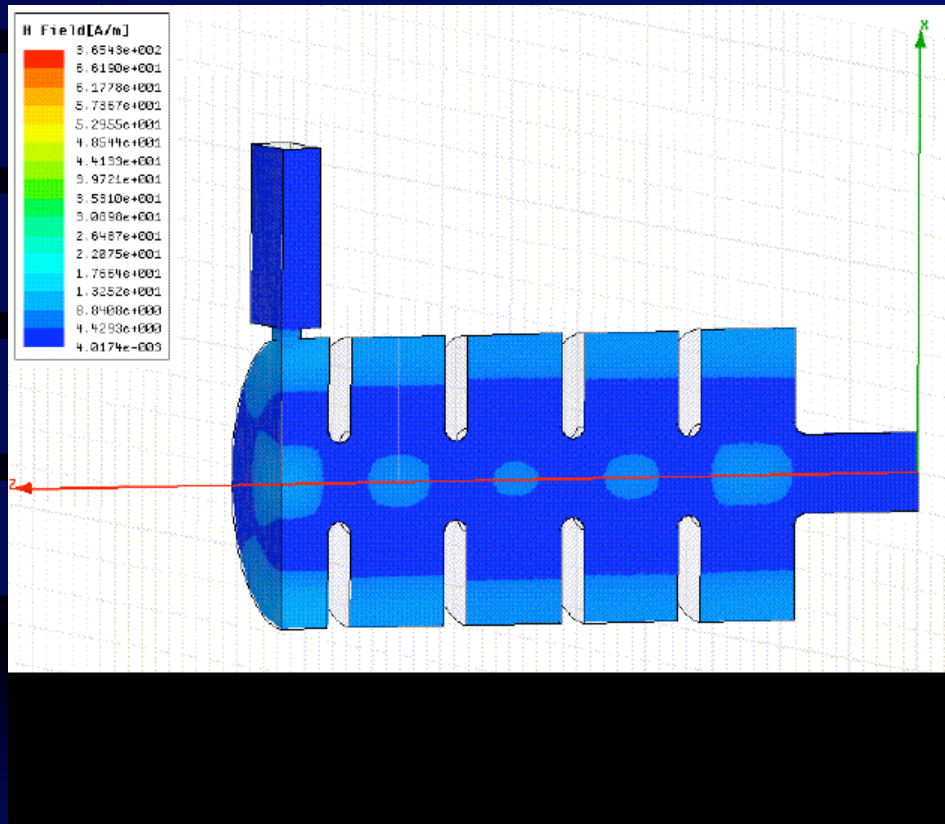
Steel with Cu coating
9-cell
9.5 GHz
cold-test only
conflat flange design
no polarization separation



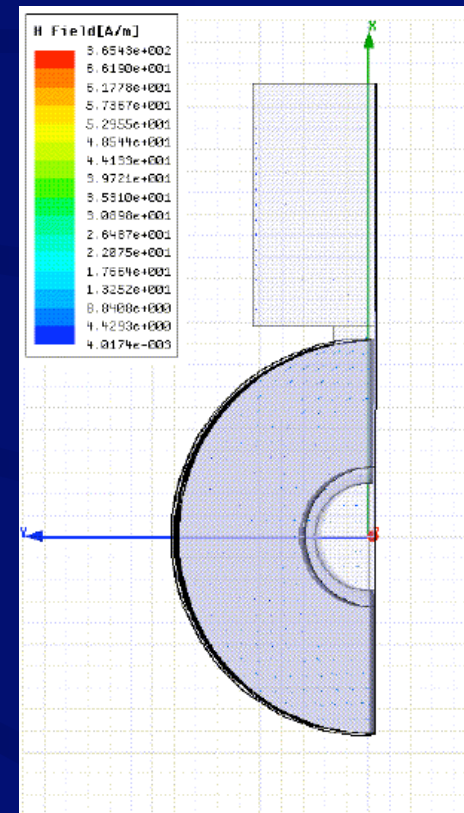
final design

GlidCop Al-15
9-cell
9.59616 GHz
tested up to 50kW peak pwr
conflat flange design
EDM'd polarization holes

Deflecting Cavity Animations

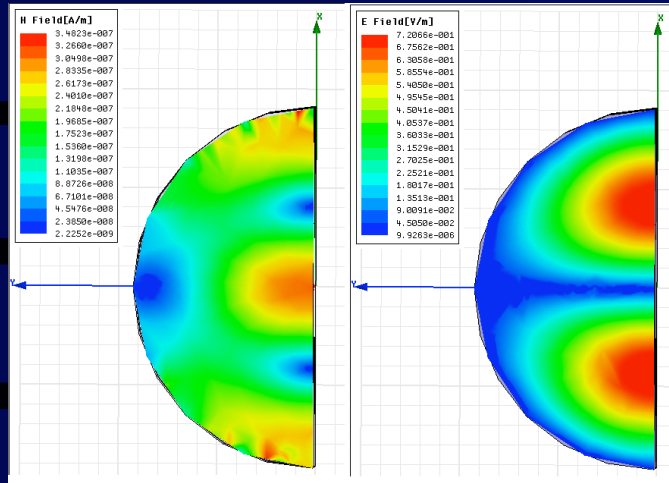


H-field complex magnitude

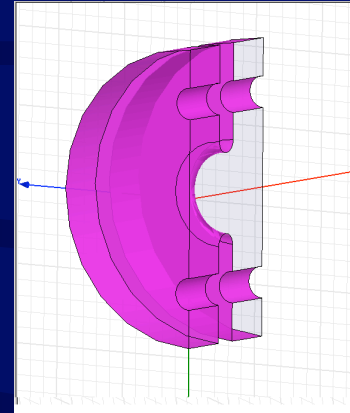


H-field vector plot

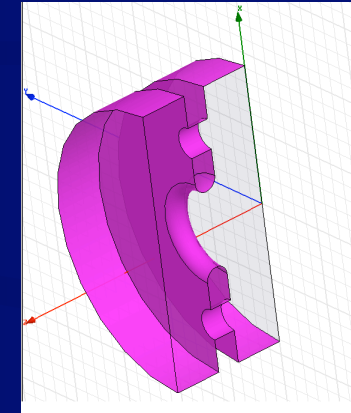
Deflecting Cavity: Polarization Separation



Rods

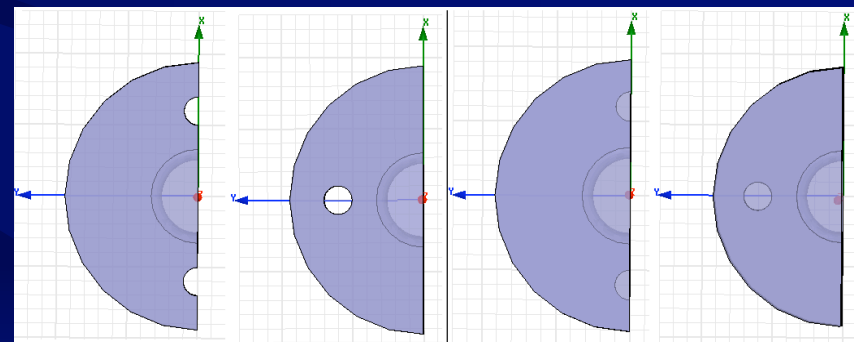


Holes



hole/rod radius = 2 mm

- Rods give larger mode separation but shift the desired mode too much
- Holes give less mode separation but don't perturb the desired mode.
- In final design, holes used with radius reduced to 1mm, giving a mode separation of 1 MHz.



Undesired
+1358 MHz

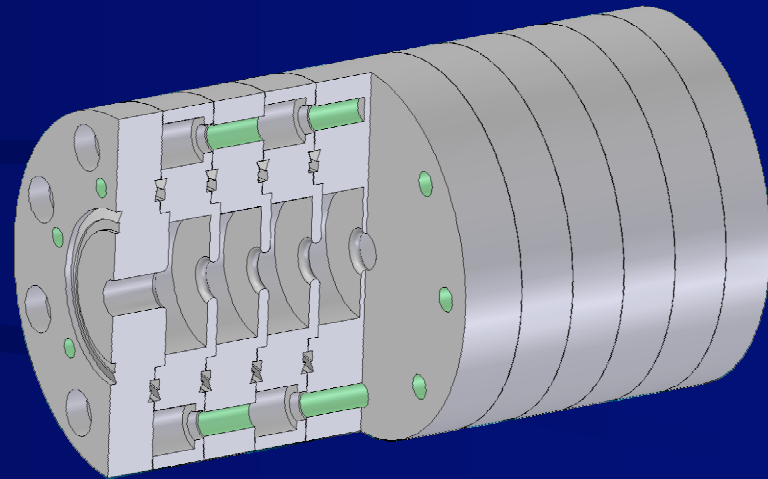
Desired
+53 MHz

Undesired
-7 MHz

Desired
-2 MHz

Final Cavity Design

- 9-cell standing wave structure
- center-fed input RF
- reconditioned VA-24G klystron
- no brazing between cells
- cells are stacked CF vacuum flanges



CAD drawing of stacked cells

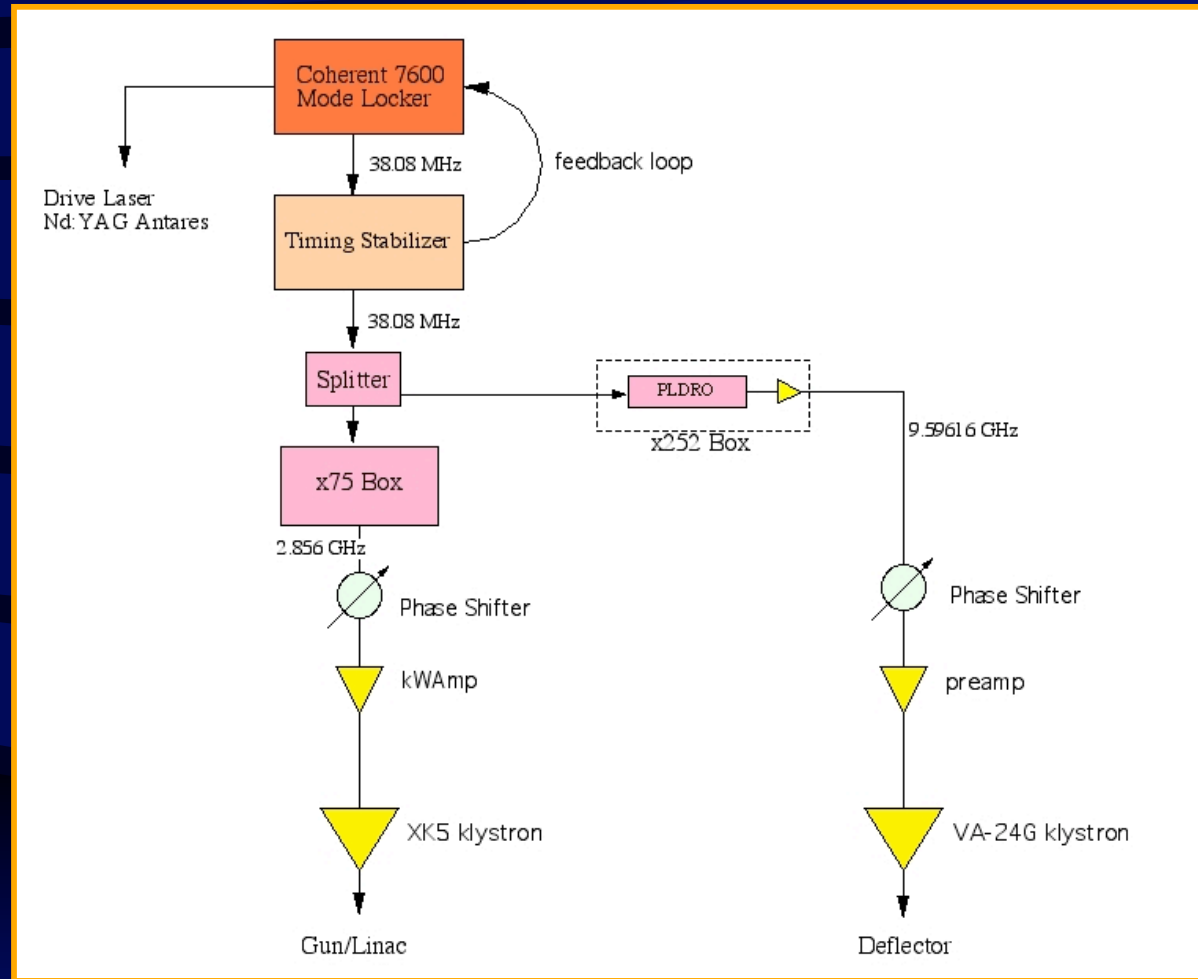


x-band klystron (50 kW peak)



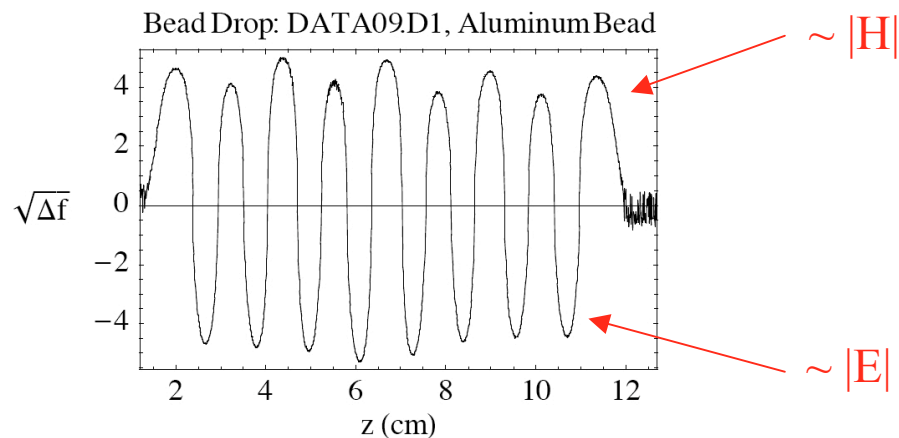
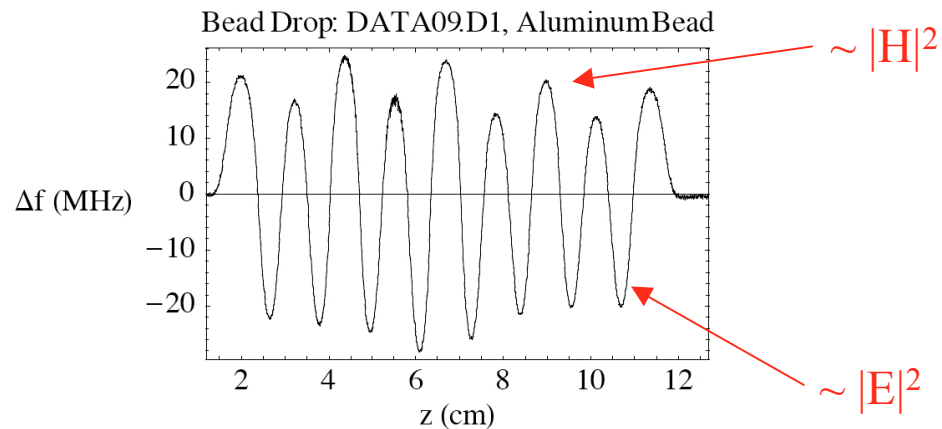
one cell with polarization holes

S-Band/X-Band RF System



- S and X-Band frequencies are multiples of modelocker freq of drive laser
- Ensures phase stability of gun, linac, laser, and deflector

Bead Pull Results



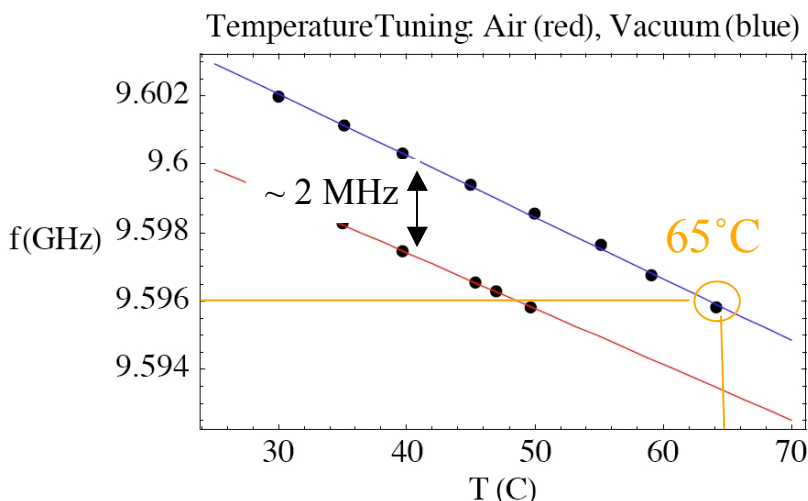
After brazing input coupler

- Bead pull using aluminum bead
- Data proportional to $|E|^2$ and $|H|^2$
- Field flatness $\sim \pm 5\%$
- Data taken at room temp (24 C)

field flatness $\sim 10\%$

$f_0 = 9.60084$ GHz ; $\Delta f = 1.5$ MHz
 $\beta = 0.870$; VSWR = 1.15
 $Q_L = 6359$; $Q_0 = 11889$; $Q_e = 13672$

Temperature Tuning



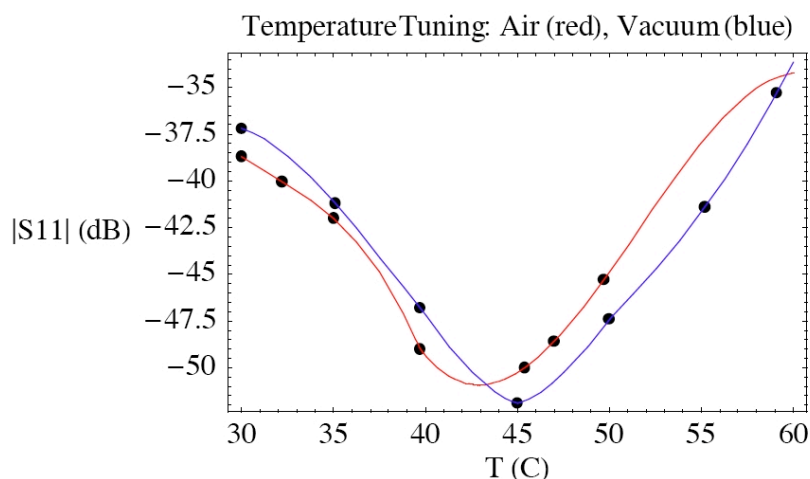
Frequency vs. Temperature

- using heater tape and thermocouple
- PID temperature feedback control
- dots are measured data
- solid lines are linear fits
- $df/dT = -179 \text{ kHz/}^\circ\text{C}$

$$f_0 = 9.596 \text{ GHz} ; \Delta f = 1.5 \text{ MHz}$$

$$\beta = 1.036 ; \text{VSWR} = 1.03637$$

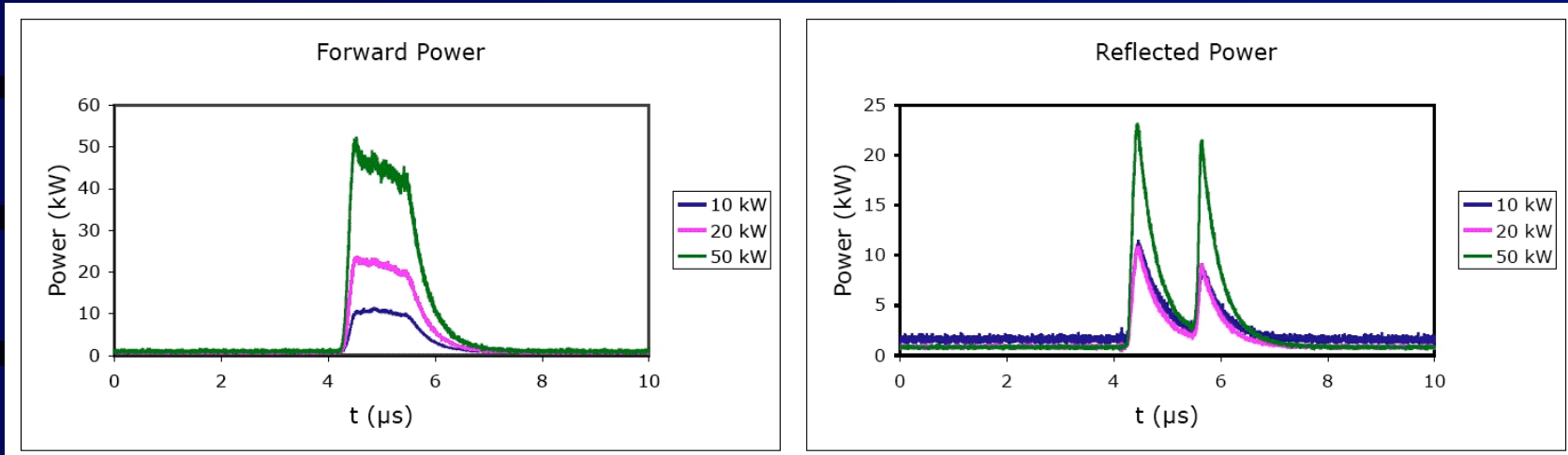
$$Q_L = 6638 ; Q_0 = 13043 ; Q_e = 13517$$



Reflectance vs. Temperature

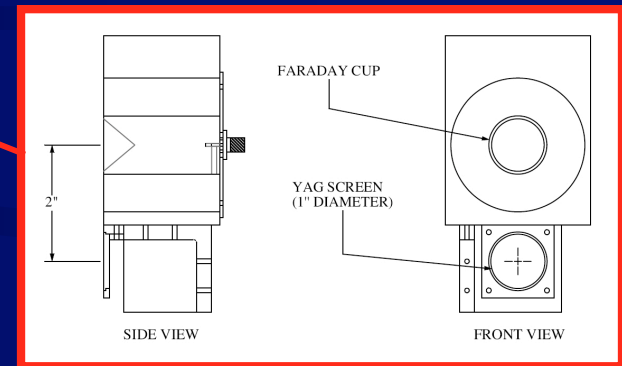
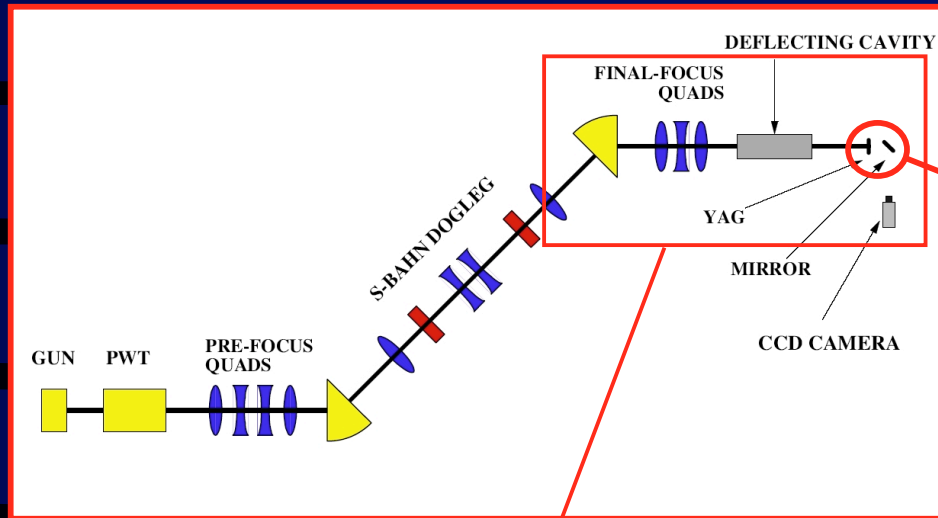
- dots are measured data
- solid lines are interpolations
- at optimal freq in vacuum (9.59616 GHz), cavity is slightly overcoupled (-35 dB @ 62 C)
- therefore operating $\beta = 1.036$ in vacuum

High Power RF Measurements

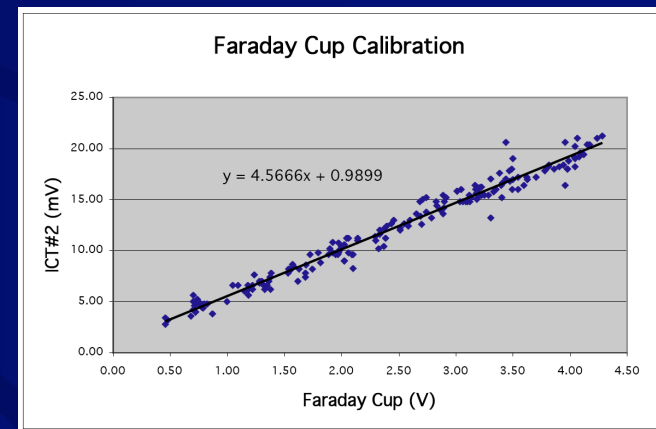
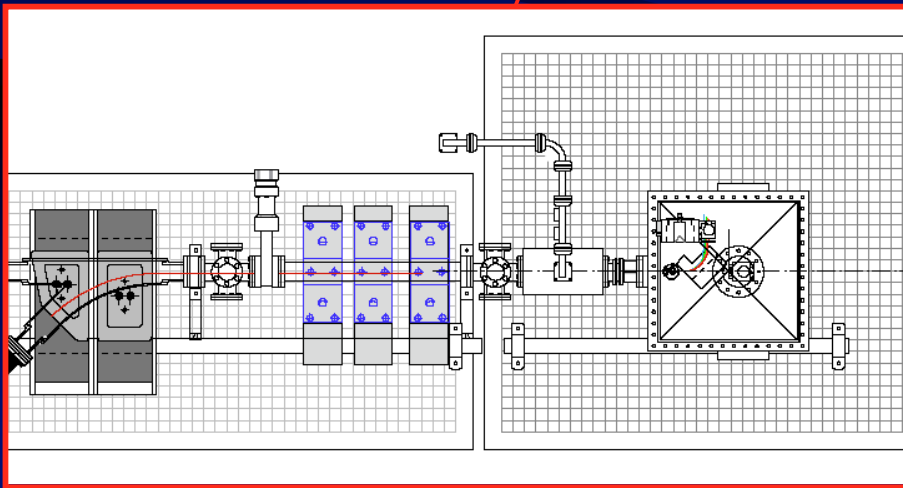


- oscilloscope traces for several attenuation settings
- measured on deflecting cavity waveguide power coupler
- maximum forward power level is 50 kW

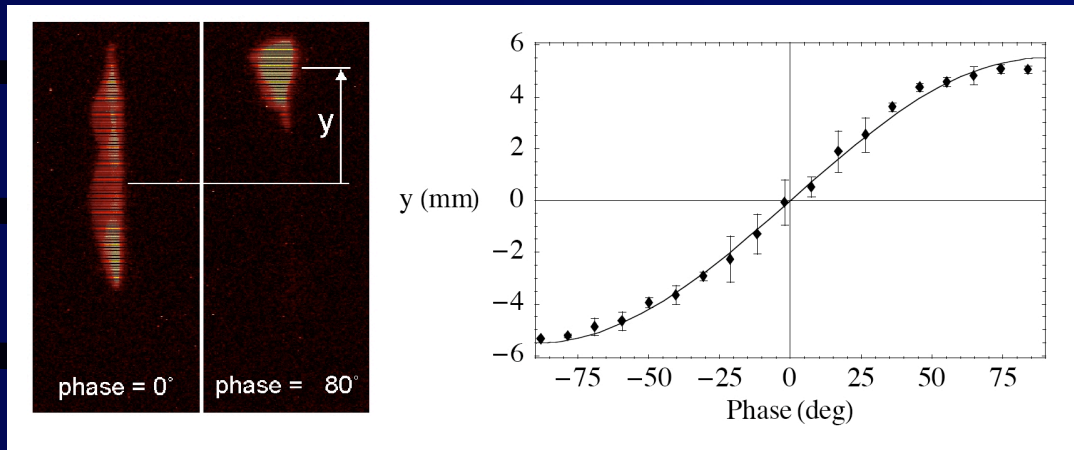
Experimental Setup



- pop-in faraday cup / 1" YAG
- YAG = yttrium aluminum garnet



Deflection vs. RF Phase



- Solid curve = sine function fit
- Amplitude = $eV_0L/p_0c = 5.5$ mm

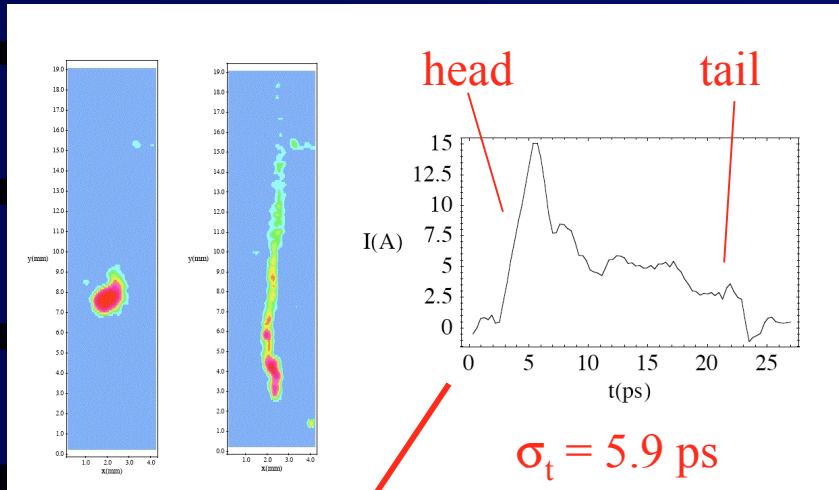
$$y_{cen} = \frac{eV_0L}{p_0c} \sin(\phi_0)$$

Method	Forward Power (kW)	V_0 (kV)
Phase Scan	9.6 kW	232
RF (50 Ω termination)	12.75	267
RF (1 M Ω termination)	12.15	261

- Comparison with RF values
- Calibrated crystal detector
- Assumption: shunt impedance = 5.6 M Ω (sim. value)

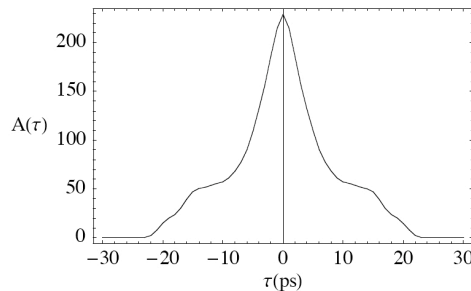
Deflecting Cavity: Uncompressed Beam

off on

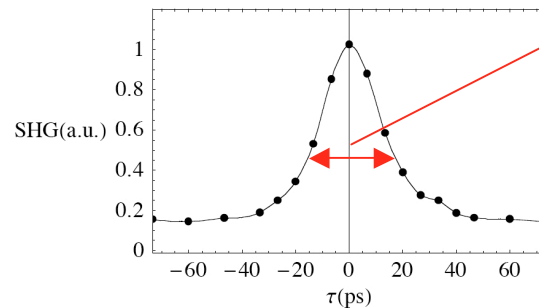


- beam is on-crest in linac (no chirp)
- therefore not compressed in dogleg
- beam appears asymmetrical
- somewhat long pulse
- a lot of structure in the tail
- in some streaks, it is more pronounced
- structure related to nonlinear xtals(?)

beam profile autocorrelation



IR drive laser autocorrelation



FWHM = 28.8 ps (IR)

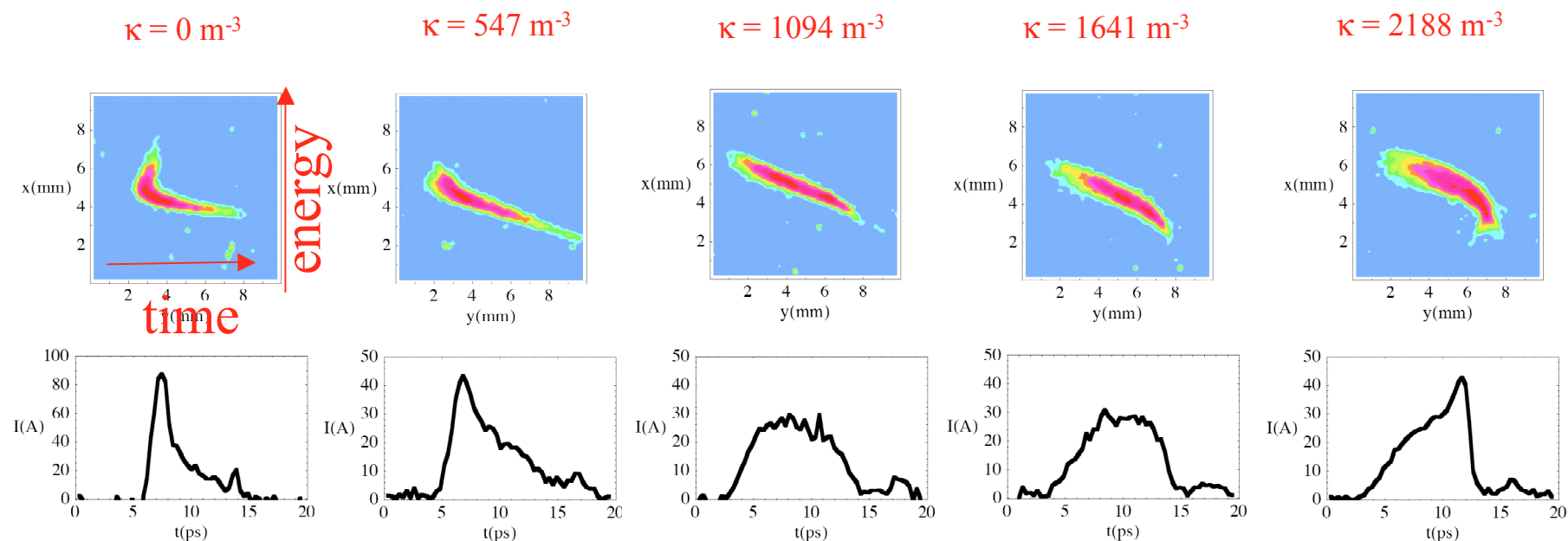


doubling xtals
(factor of 2)

FWHM = 14.4 ps (UV)

$\sigma_{\text{rms}} \sim 7$ ps

Deflecting Cavity: Compressed Beam

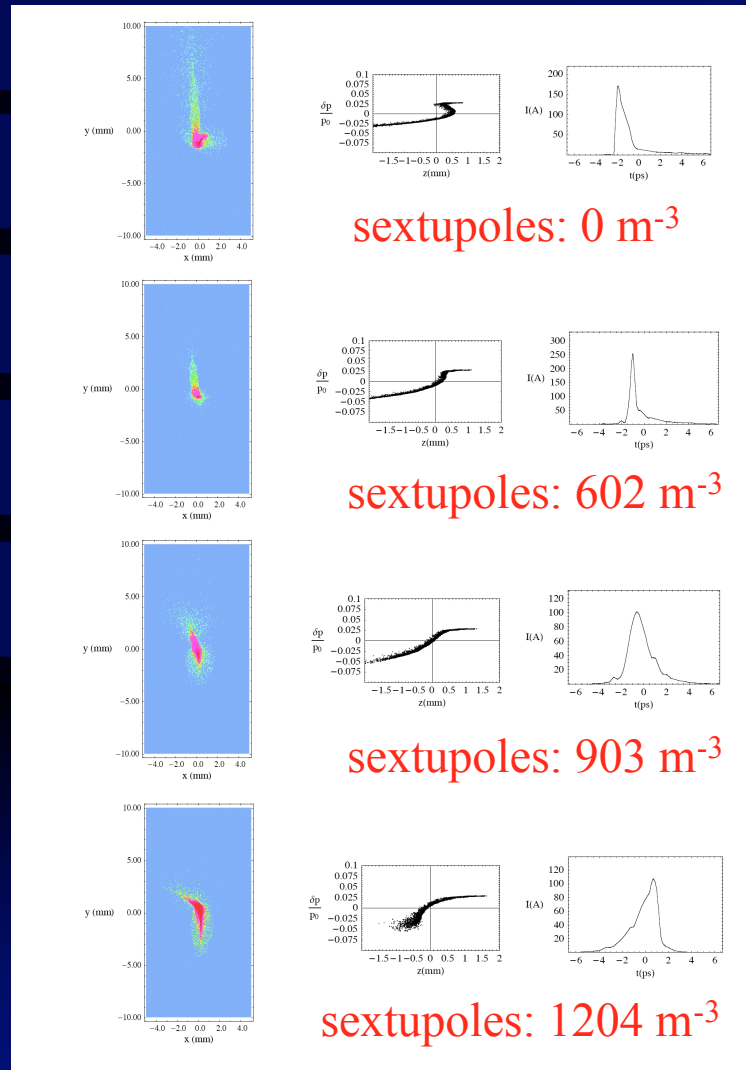


- chirped 20° in linac, 234 pC of charge at 11.8 MeV with $V_0 = 400 \text{ kV}$
- residual horizontal dispersion produces pseudo-phase space reconstruction
- combination of linear and nonlinear effects (R_{16} & T_{166})
- ramping mechanism clearly visible
- due to asymmetry of initial pulse, overcompensation with sextupoles needed

R. J. England, J. B. Rosenzweig, and G. Travish, PRL **100**, 214802 (2008).

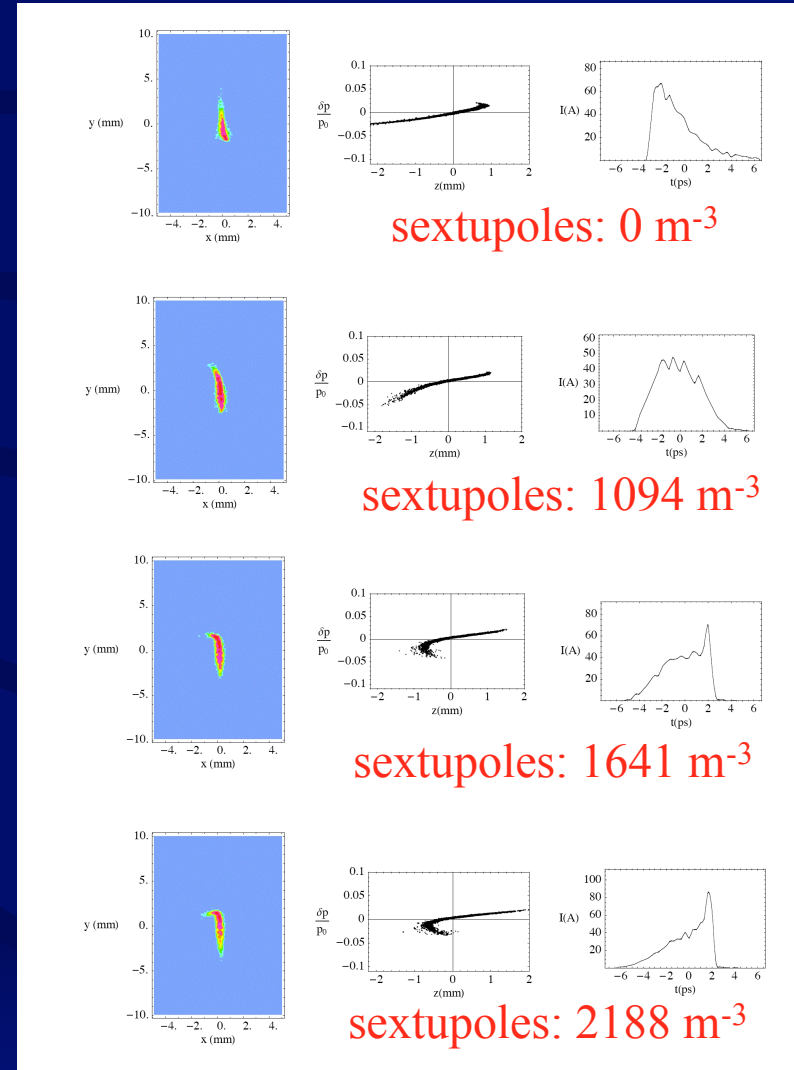
ELEGANT Simulation

“streak” in x,y z phase space current profile



Gaussian beam

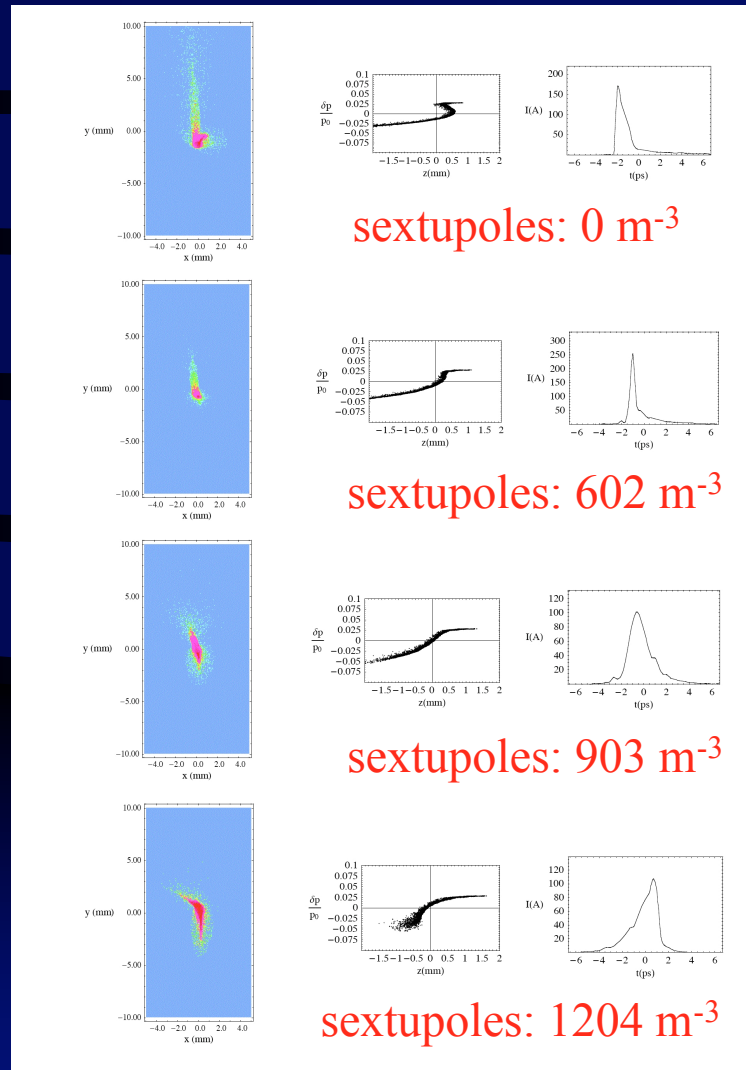
“streak” in x,y z phase space current profile



Asymmetric (front-heavy) beam

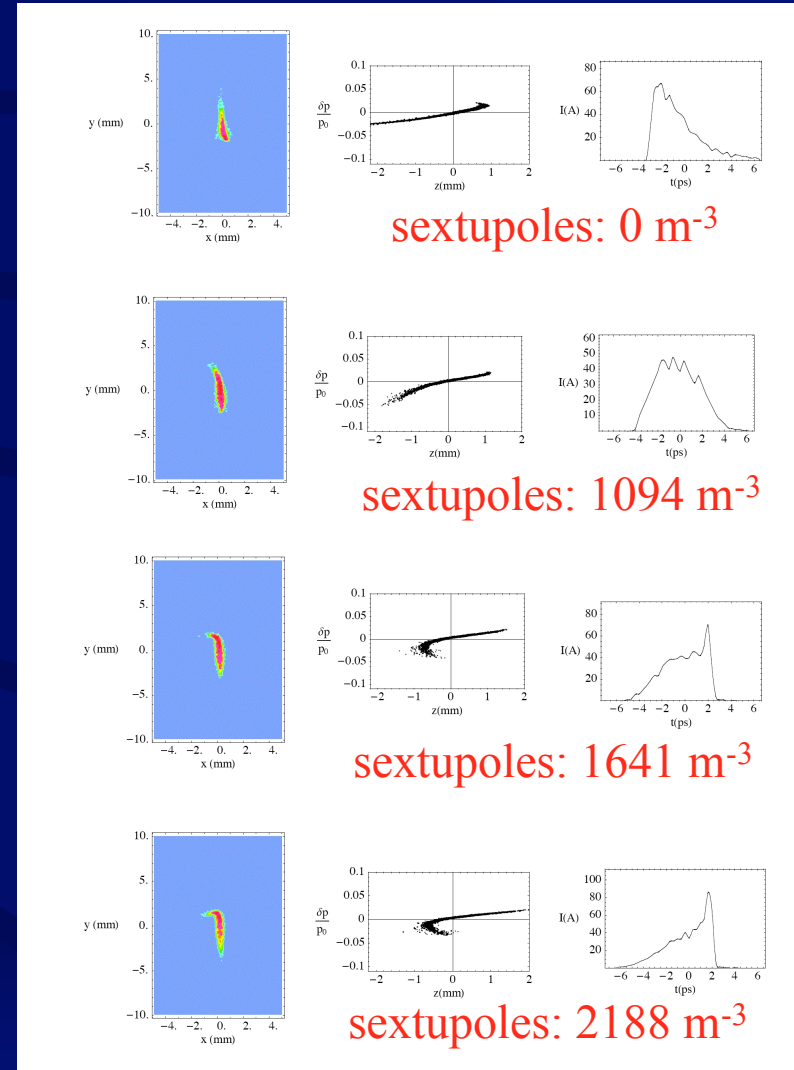
ELEGANT Simulation

“streak” in x,y z phase space current profile



Gaussian beam

“streak” in x,y z phase space current profile

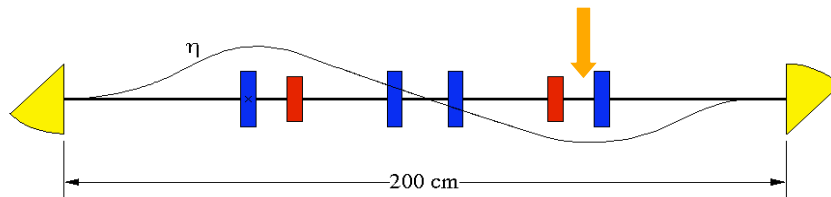


Asymmetric (front-heavy) beam

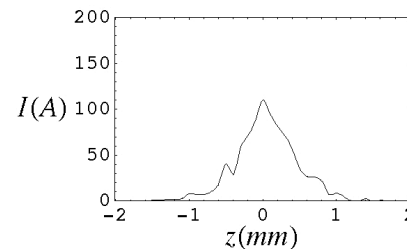
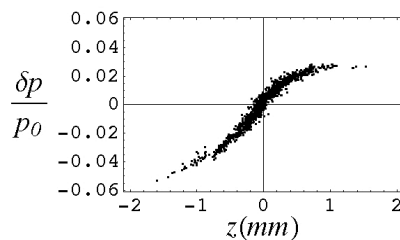
Future Applications

Witness Beam Generation

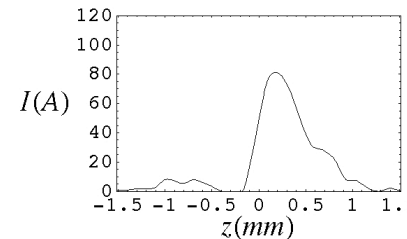
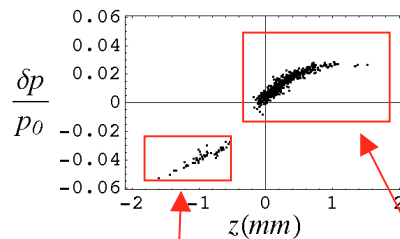
For PWFA application, drive beam needs a witness beam to accelerate.



Region of high dispersion in x
Strong correlation b/w x and z
Insert mask in x to sever beam in z



No mask inserted
Undercorrected with sextupoles to
elongate profile



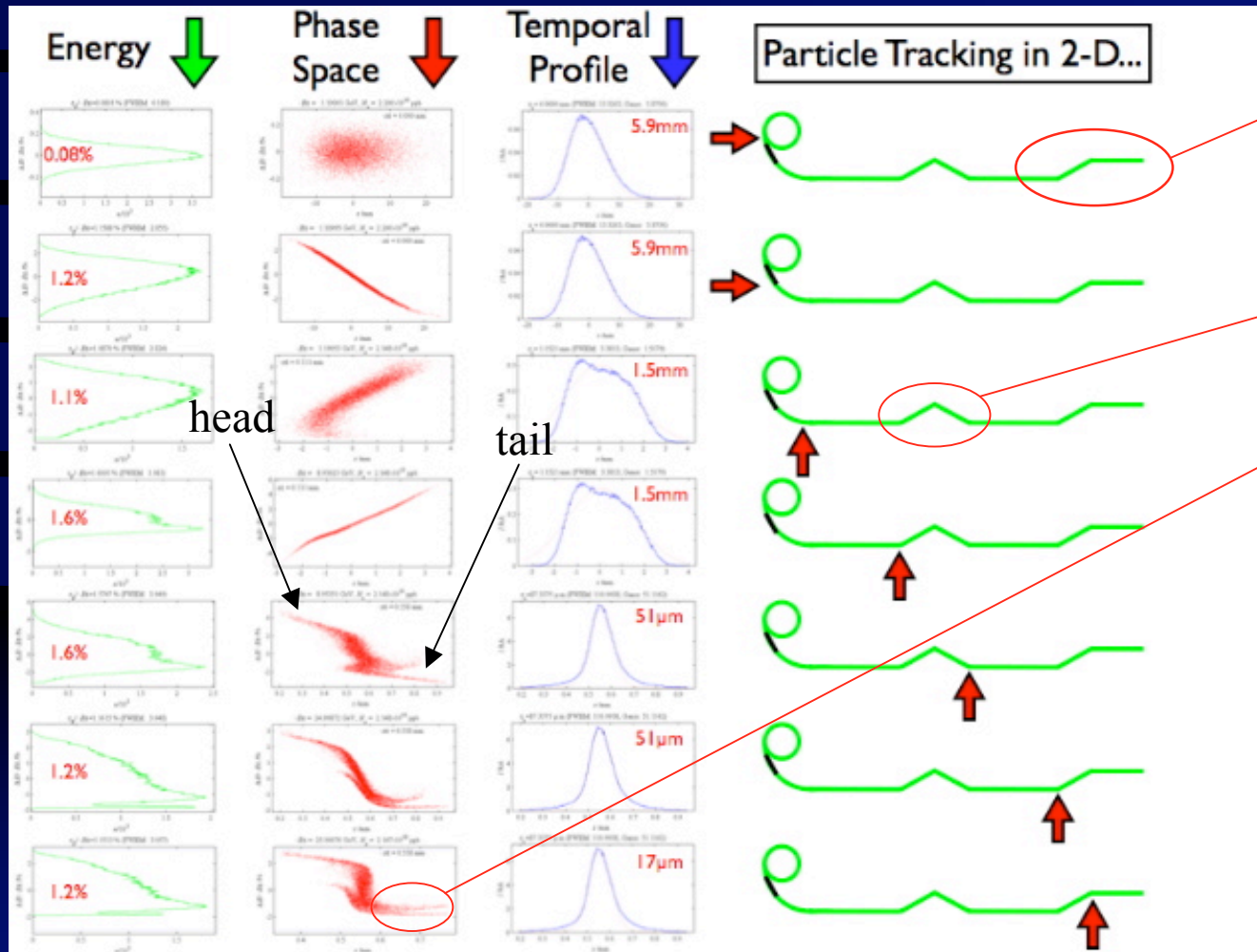
With 1cm mask inserted at above
location

witness beam

ramped drive beam

Future Applications

Tailored Profiles for FACET?



dogleg: $R_{56} = -7\text{cm}$

- reduce chicane compression to increase bunch length to $\sim 2 k_p^{-1}$
- collimate to remove low-energy tail

Conclusions

- Proposal:
 - ramped beams: improved transformer ratio ($R > 2$) for PWFA applications
 - feasible using dogleg compression with sextupoles
 - deflecting cavity diagnostic (500 fs resolution)
- Deflecting cavity
 - final cavity design finalized in 3-phase process w/ 2 prototypes
 - cavity testing indicates that it operates within the design specifications
 - high power RF testing: no breakdown problems observed
- Experimental tests:
 - unchirped (uncompressed) beam has asymmetric structure
 - chirped beam w/residual dispersion = semi-tomographic reconstruction
 - evidence for ramp-shaped electron beams
- Other Experiments:
 - deflector use for measuring optimized charge distributions
 - dogleg high-brightness focus studies
 - witness bunch generation

Acknowledgements

Laser-Plasma Group

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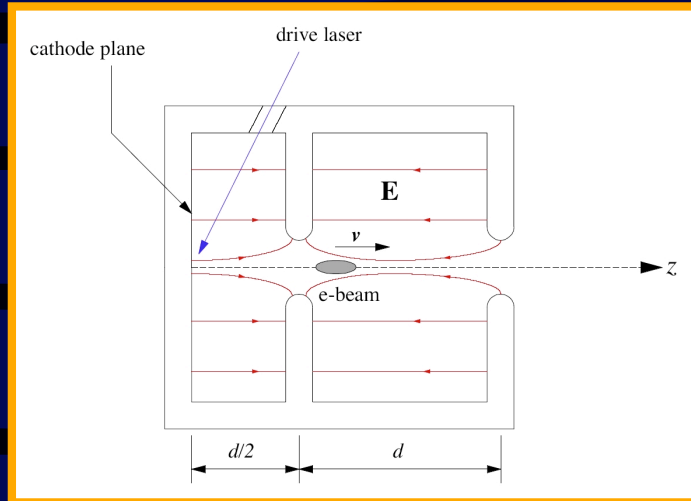
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Mike Dunning
.... and many more

Introduction: The RF Photoinjector



- Acceleration from rest to relativistic energies (~ 1 to 10 MeV)
- Temporal structure of electron beam reflects that of laser pulse on the cathode.
- Capable of producing low-emittance beams.
- Emittance: figure of merit; measure of area occupied by beam distribution in transverse phase space.

$$\epsilon_{x,N} = \frac{1}{mc} \sqrt{\langle p_x^2 x^2 \rangle - \langle p_x x \rangle^2} \approx \sqrt{\epsilon_{x,th}^2 + \epsilon_{x,rf}^2 + \epsilon_{x,sc}^2}$$

thermal emittance: $\epsilon_{x,th} \propto \sigma_x$

RF emittance: $\epsilon_{x,rf} \propto \sigma_x^2 \sigma_z^2$

space charge emittance: $\epsilon_{x,sc} \propto \frac{1}{\sigma_z}$

- Trade-off between rf and sc components
- Implies optimal pulse length σ_z .
- Generally determined by photoinjector codes (e.g. PARMELA, HOMDYN)
- Typical $\sigma_z \sim 10$ degrees of RF phase
- For S-Band (2.856 GHz) $10 \text{ deg} \sim 10 \text{ ps}$

Beam Brightness

$$B_{\perp} = \frac{I}{\epsilon_{x,N}^2}$$

$$I \propto \frac{Q}{\sigma_z}$$

$$Q_{\max} \approx \frac{2}{5} \pi \sigma_x^2 \epsilon_0 E_0$$

- “brightness”: measure of density of particles in transverse phase space.
- emittance constrained by photoinjector: $\epsilon_N > 1 \mu\text{m}$
- σ_z constrained (~ 10 deg of RF phase) to minimize ϵ_N
- Q constrained by cathode image charge limit

Rosenzweig, Colby, AIP Conference Proc. 334, p. 724-737 (1995).

- optimal photoinjector brightness
- to obtain higher brightness beams, require compression techniques

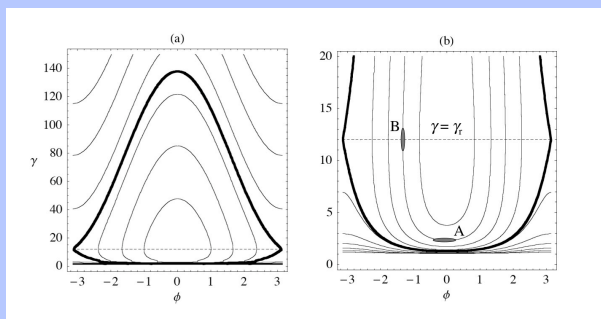
$$B_{\perp,opt} = 16(2\pi)^{9/2} \alpha k \frac{I_A [1 + \frac{3}{5} A]^2}{\sigma_z A^2}$$

$$\alpha=1.5; \lambda=10\text{cm}; A=1; I_A=16\text{kA}; \sigma_z=3\text{mm}$$

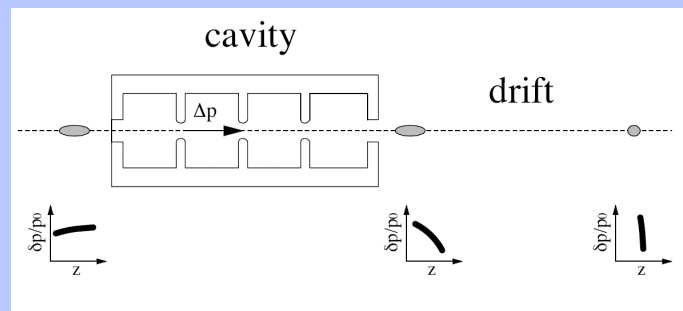
$$B_{\perp,opt} = 80 \text{ A}/\mu\text{m}^2$$

Bunch Compression Techniques

RF Techniques

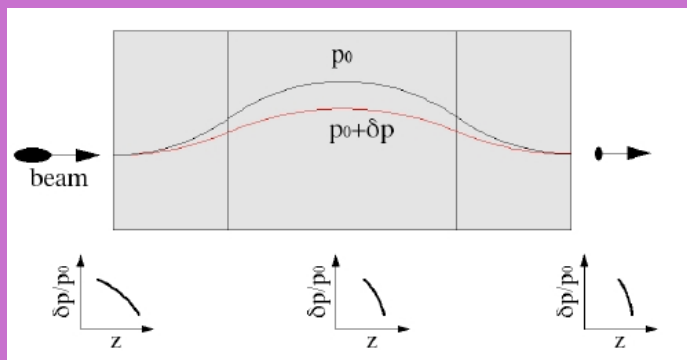


Phase Space Rotation

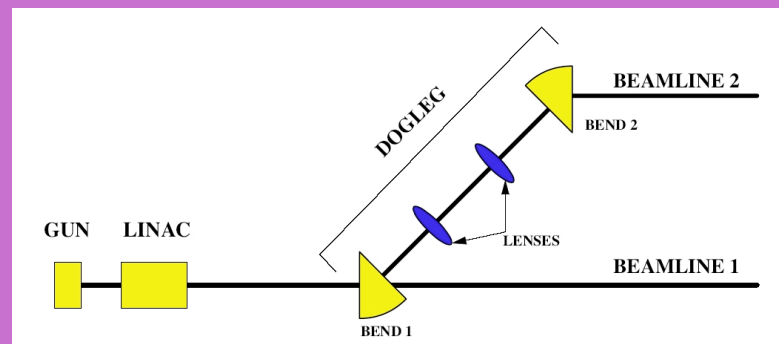


Ballistic Compression

Magnetic Techniques



Chicane



Other Nonisochronous Devices
(e.g. dogleg compressor)

Applications for High Brightness Beams

Free electron laser

- high gain regime
- minimize the gain length L_g

$$L_g = \frac{\lambda_u}{2\sqrt{3}\pi\rho} \quad \rho \propto \omega_p^{2/3} \propto n_e^{1/3}$$

Inverse Compton Scattering

- beam + laser ($h\nu$) $\rightarrow$ higher $h\nu'$ photons
- shortness of scattered pulse limited by shortest of beam, laser

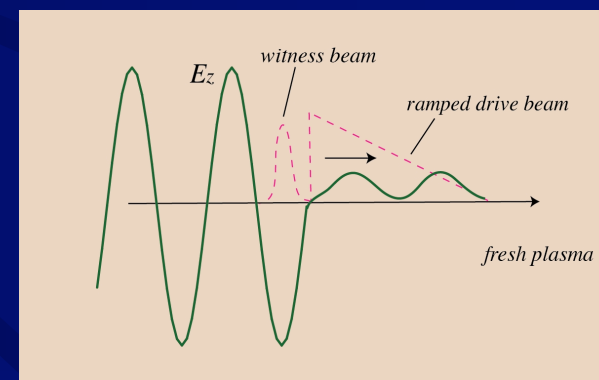
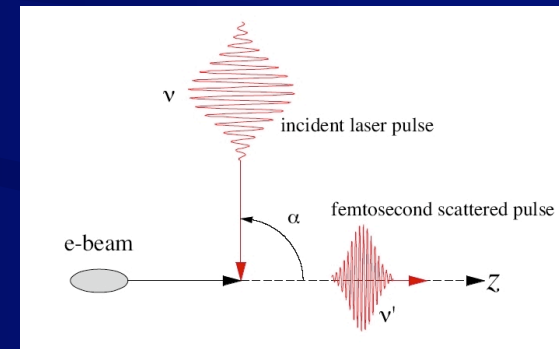
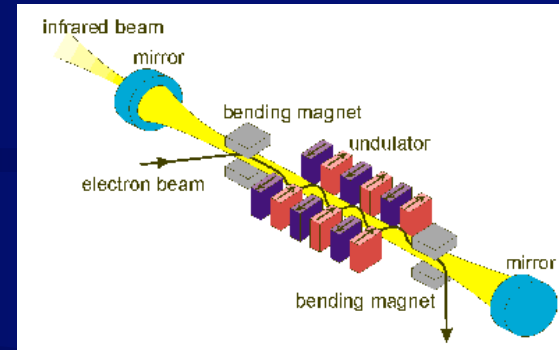
$$N_{ph} = \mathcal{L}\sigma_T \quad \mathcal{L} \propto \frac{N_e}{A_{int}}$$

Plasma Wakefield Accelerator:

- beam + plasma $\rightarrow$ high-gradient wakes
- beam density, time profile important

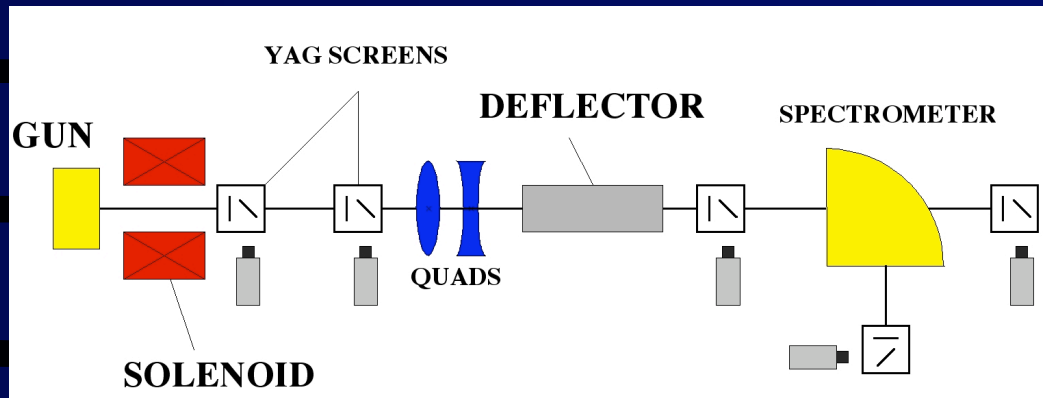
$$E_{\max} = E_0 1.3 \Lambda \ln(1 / \sqrt{\Lambda / 10}) \quad n_{\text{beam}} \gg n_0$$

$$\Lambda = (n_{\text{beam}} / n_0) k_p^2 \sigma_r^2$$

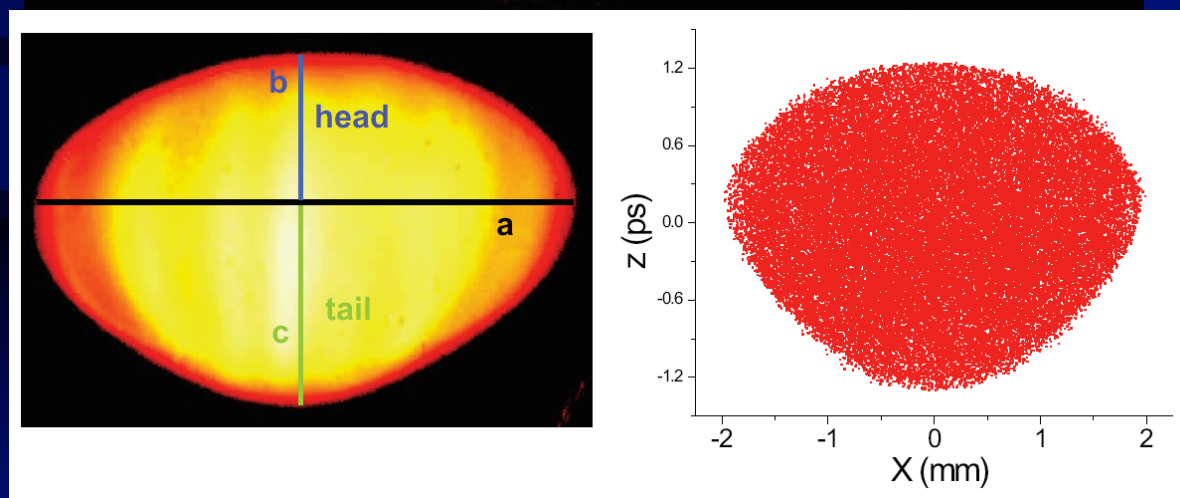


Further Applications: Deflecting Cavity

Dynamically Optimized Beam Profiles



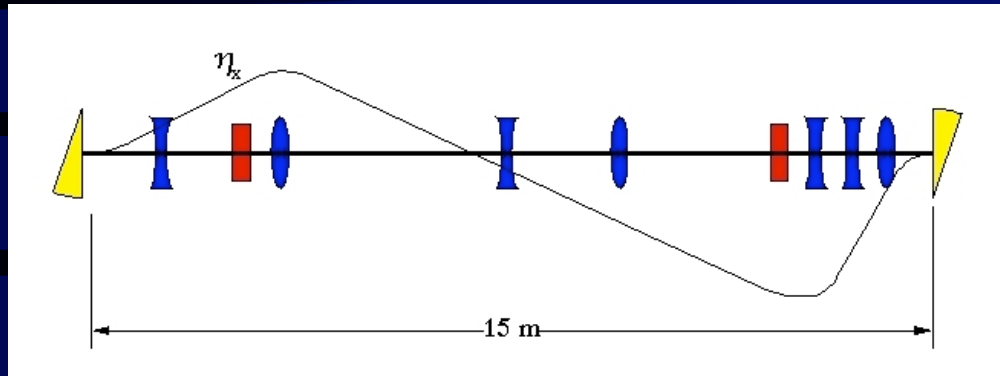
Beam Charge:	0 to 50 pC
Beam energy:	up to 4 MeV
Power Source:	18 MW Klystron
RF Frequency:	2.856 GHz (S-Band)
Cathode laser:	10 μ J at $\lambda = 266$ nm
Laser pulse length:	50 fs RMS



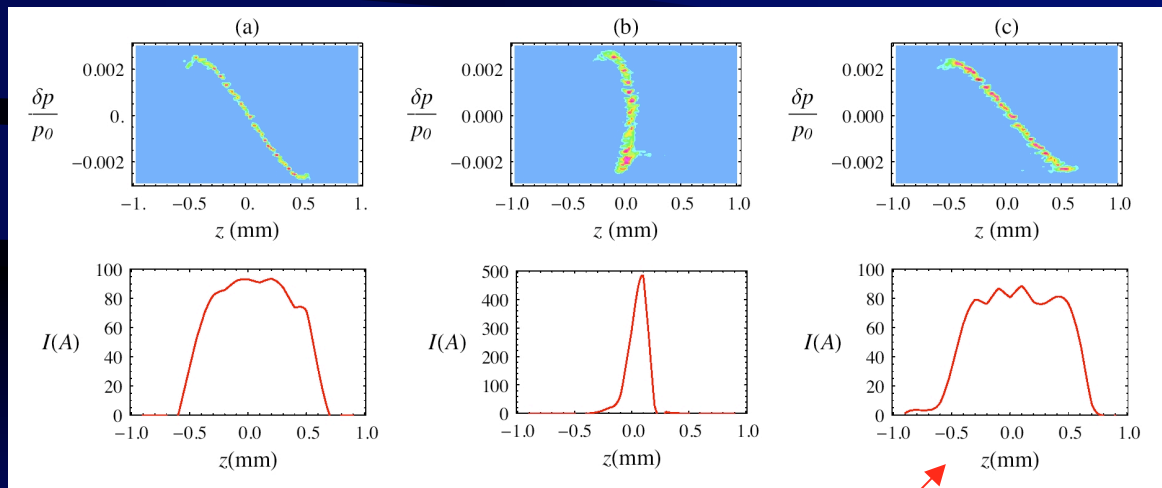
P. Musumeci, J. T. Moody, R. J. England, J. B. Rosenzweig, T. Tran “Experimental generation and characterization of uniformly filled ellipsoidal beam distributions,” (submitted to PRL).

Further Applications: Doglegs

BNL - VISA I



- observed compression when running off-energy by 0.76%
- however, for negative chirp, dogleg should expand, not compress



sextupoles used to remove T_{566}

R. J. England, et. al. "Sextupole correction of the longitudinal transport of relativistic beams in dispersionless translating sections," PR-STAB 8, 012801 (2005).

-0.4 cm -10 m (!)

$$z_f = z_0 + R_{56}\delta + T_{566}\delta^2$$

 off energy by $\Delta = -0.76\%$

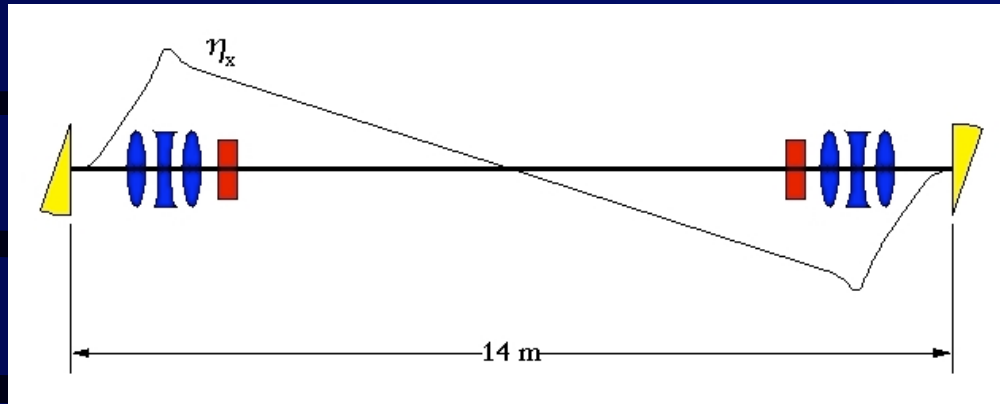
$$\hat{R}_{56} \simeq R_{56} + 2T_{566}\Delta$$

$$z_f = z_0 + \hat{R}_{56}\hat{\delta} + \hat{T}_{566}\hat{\delta}^2$$

 $+18 \text{ cm}$

Further Applications: Doglegs

SLAC - ORION Low-E Hall Dogleg



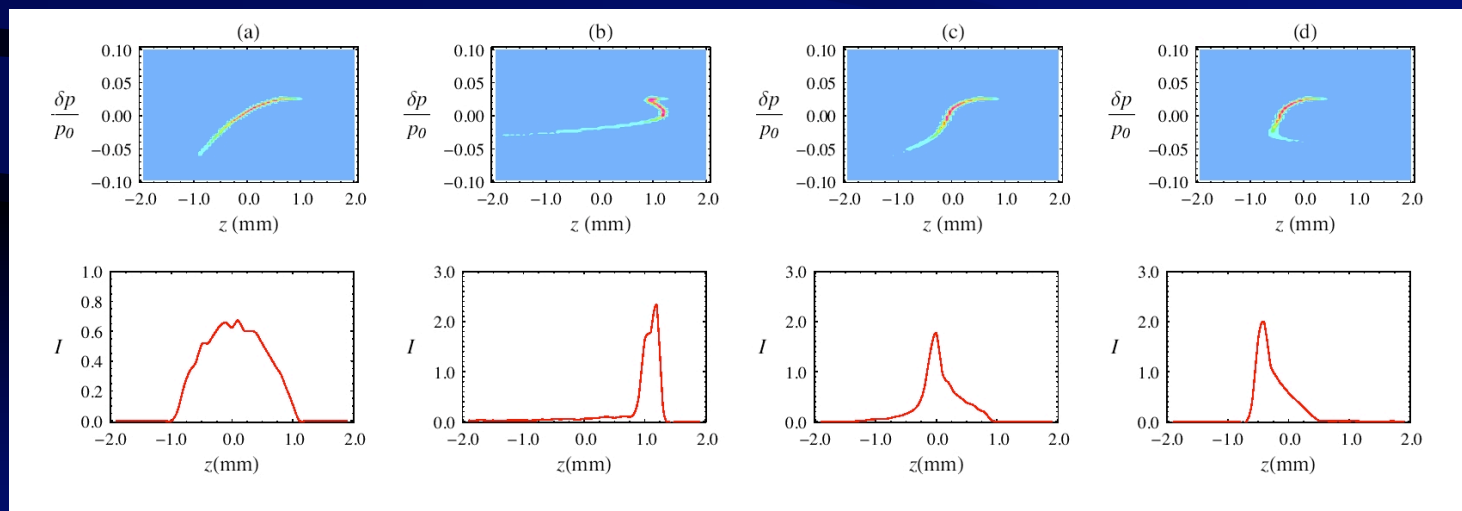
- studies for PWFA / general transport
- large energy spread requires octupole correction

initial beam

no sextupoles

with sextupoles

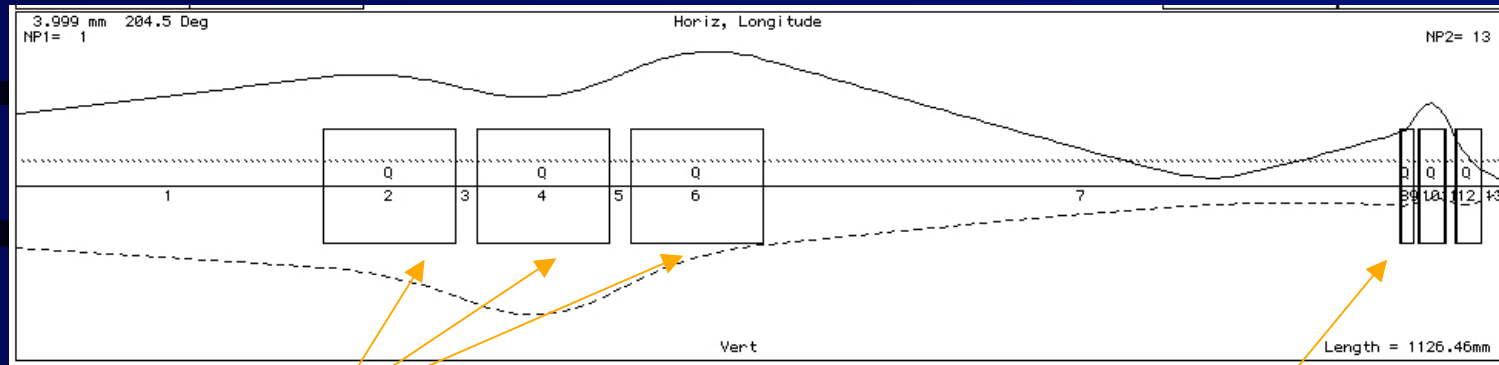
sext's + oct's



R. J. England, et. al. "Sextupole correction of the longitudinal transport of relativistic beams in dispersionless translating sections," PR-STAB **8**, 012801 (2005).

Future Experiments

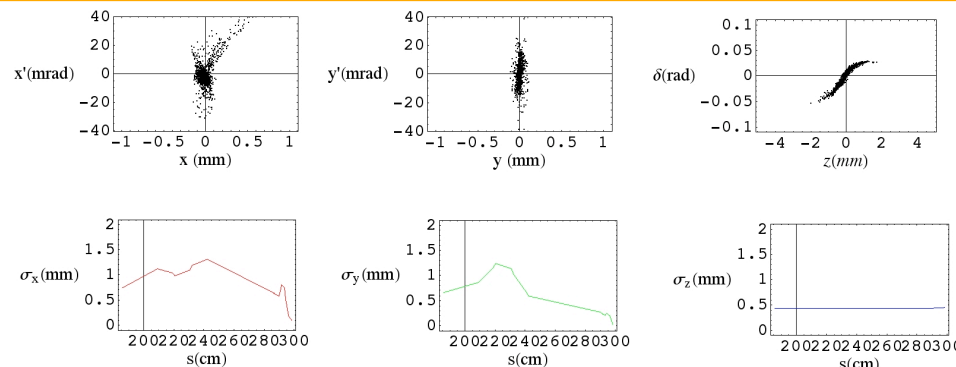
PMQ Focusing



standard iron quads

PowerTrace 1.08 Simulation

PMQs (110 T/m)



- Hybrid Permanent Magnet and Iron
- Grey cubes are Alnico; $M=1.175$ T
- Field gradient: 13 MeV; $Q = 300$ pC
- $\sigma_x = 130 \mu\text{m}$; $\sigma_y = 11 \mu\text{m}$; $\sigma_z = 0.5 \text{ mm}$
- Bore diameter = 18 mm
- Benefits: cheaper, better field profile
- Downsides: small bore; in-vacuum

$$B \approx \frac{2Q / \sigma_t}{\epsilon_{N,x} \epsilon_{N,y}} = 412 \text{ mA} / \mu\text{m}^2$$