

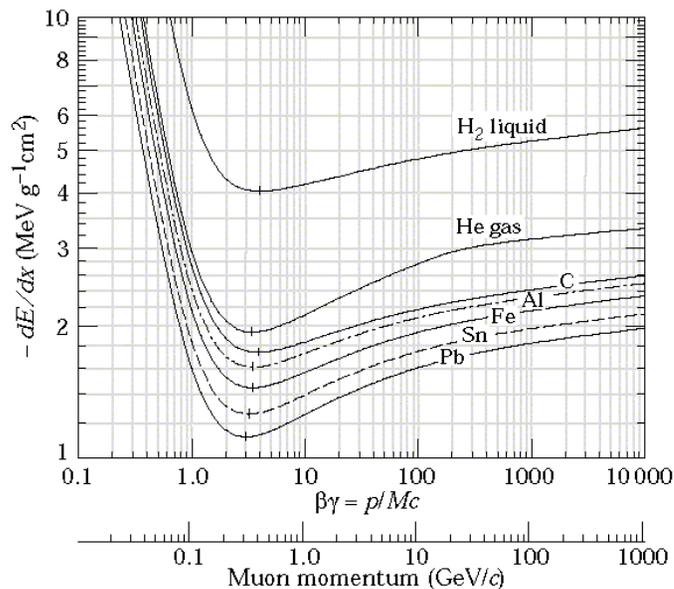
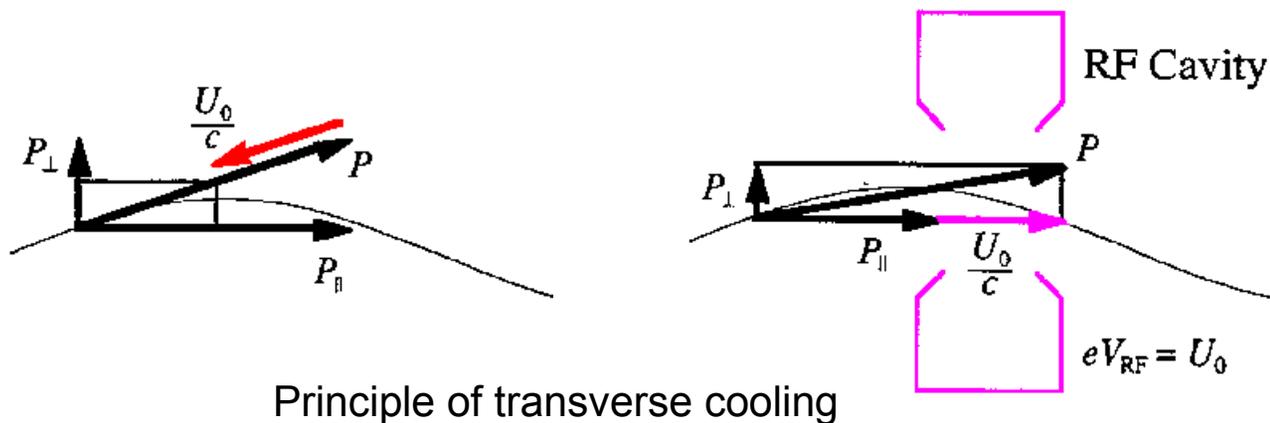


FERMI NATIONAL ACCELERATOR LABORATORY

US DEPARTMENT OF ENERGY

# Circularly Inclined Solenoid Channel for 6D Ionization Cooling of Muons

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There is no longitudinal cooling in the most suitable range 2-300MeV/c.

With higher momentum  $p > 300\text{MeV}/c$  it is difficult to obtain small  $\beta$ -function which is necessary for small equilibrium emittance:

$$\frac{d\varepsilon_N}{ds} = -\frac{1}{\beta^2 E} \frac{dE}{ds} \varepsilon_N + \frac{\beta\gamma}{2} \frac{\beta_{\perp}}{\beta} \frac{d\langle\theta_{\text{rms}}^2\rangle}{ds}$$

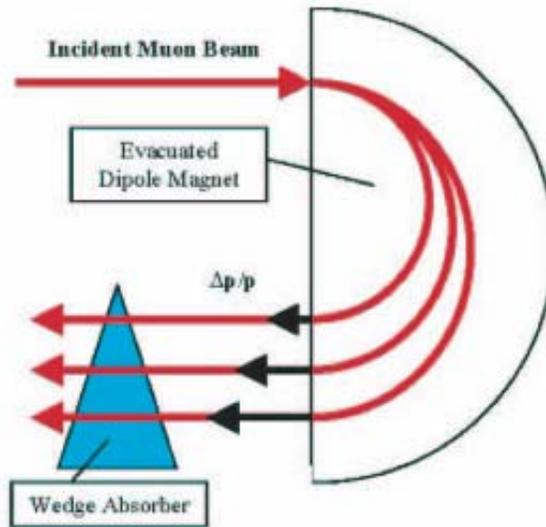


Figure 1. Use of a Wedge Absorber for Emittance Exchange

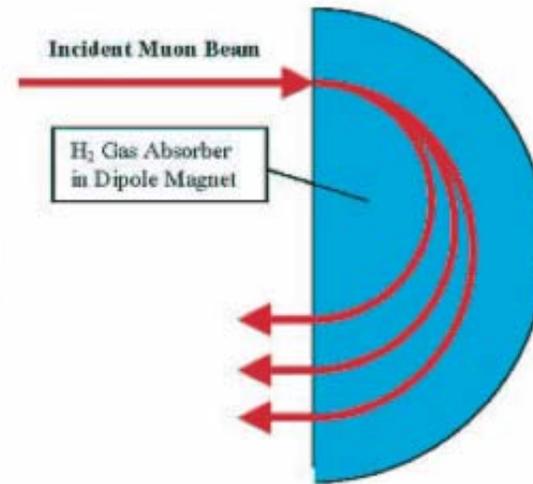
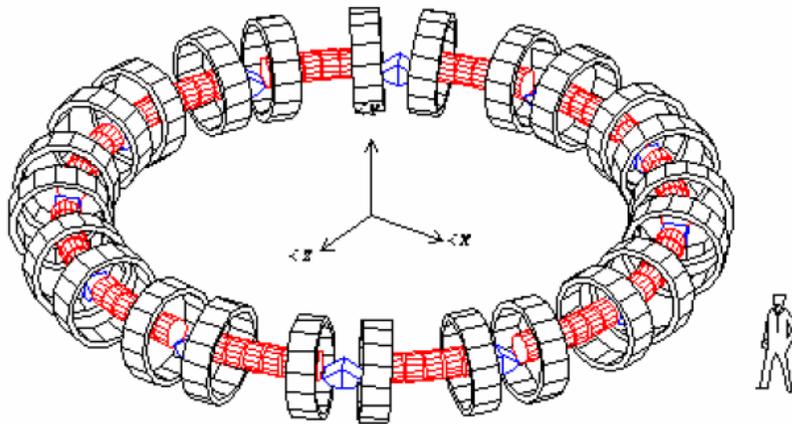


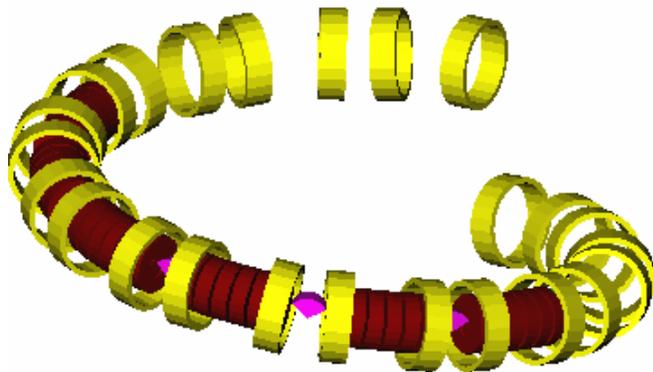
Figure 2. Use of Continuous Gaseous Absorber for Emittance Exchange

Dispersion and/or large positive momentum compaction  
 ⇒ higher momentum muons make longer path in the absorber ⇒ lose more energy ⇒ longitudinal cooling

RFOFO ring

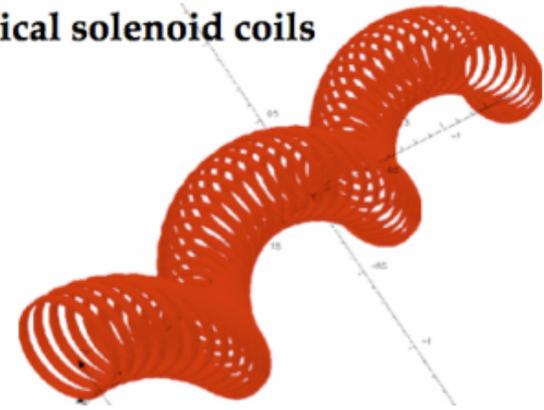


“Guggenheimed”



Helical Cooling Channel

Helical solenoid coils



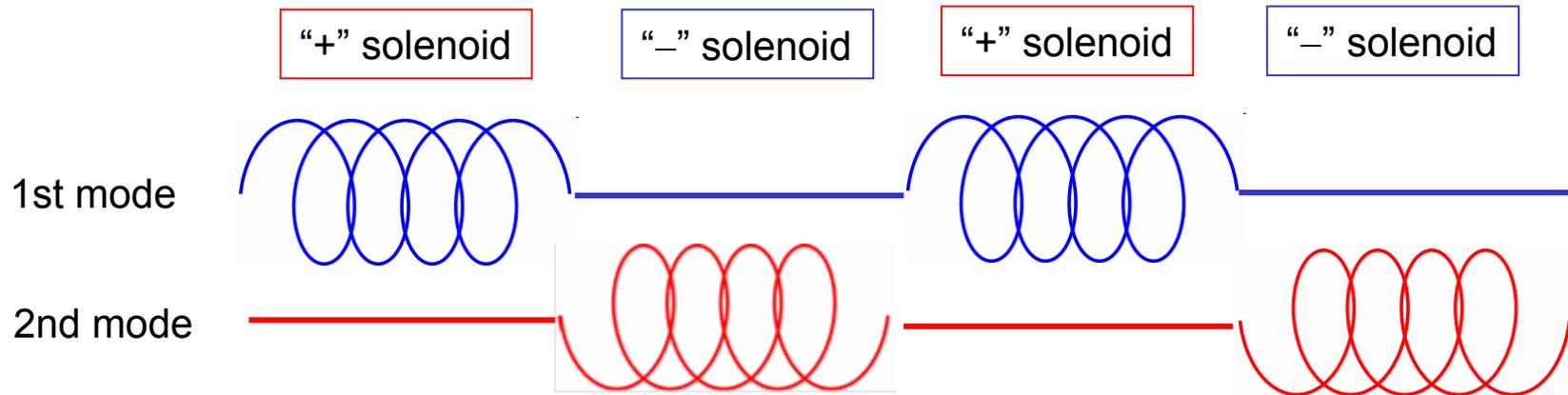
“Guggenheim” : too long  $\Rightarrow$  muon decay;  
problem with RF in magnetic field.

HCC: no viable solution yet for RF inside coils.

Both channels are selective to muon sign, it is  
either  $\mu^+$  or  $\mu^-$

## Alternating solenoid focusing

- necessary for damping both transverse modes in ~ straight channel



## Resonance dispersion generation:

resonant dependence of the periodic orbit on the tune,  $x_{p.o.} \sim 1/\sin\pi Q_x$ ,  
 + large chromaticity  $\Rightarrow$

$$D_x = \frac{dx_{p.o.}}{d\delta_p} = -\pi Q'_x x_{p.o.} \cot(\pi Q_x)$$

Naturally  $Q'_x < 0 \Rightarrow D_x \cdot x_{p.o.} > 0 \Rightarrow$  positive momentum compaction

The initial proposal - planar snake - reported at PAC07. Further analysis revealed certain difficulties:

- Insufficient transverse dynamic aperture

$$\frac{\partial B_z}{\partial z} \Rightarrow \frac{\partial B_r}{\partial r}, \quad \frac{\partial^2 B_z}{\partial z^2} \Rightarrow \frac{\partial^2 B_z}{\partial r^2}, \quad \frac{\partial^3 B_z}{\partial z^3} \Rightarrow \frac{\partial^3 B_r}{\partial r^3} \quad \text{etc.} \quad \Rightarrow \text{strong aberrations}$$

Cures:

- increase number of cells per period 4→6 (smaller phase advance per cell)
- increase solenoid radius (smoother  $B_z$ )

This helped, but brought about another problem:

- Unequal damping rates of transverse modes

Cure: mix the transverse modes better by

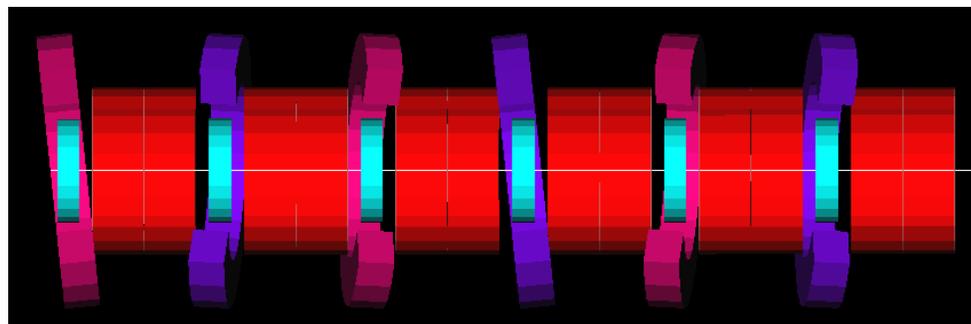
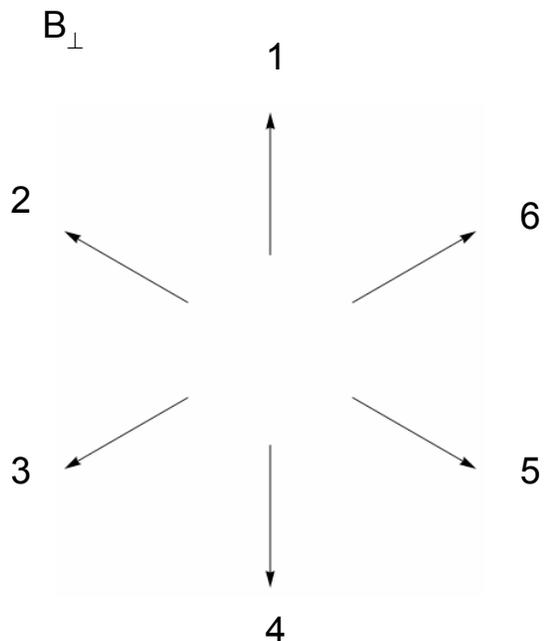
- adding a constant magnetic field component (R.Palmer)
- making the snake helical (the subject of this report)

- Choose phase advance/period (6-cell period here)  $> 2\pi \times \text{integer}$  to obtain  $\alpha_p > 0$
- Create rotating  $B_{\perp}$  field by tilting (or displacing) solenoids in rotating planes

$$x \cdot \cos(\phi_k) + y \cdot \sin(\phi_k) = 0, \quad k=1, 2, \dots$$

Example for 6-cell period:

Solenoid #	1	2	3	4	5	6
Polarity	+	-	+	-	+	-
Roll angle $\phi_k$	0	$2\pi/3$	$4\pi/3$	0	$2\pi/3$	$4\pi/3$



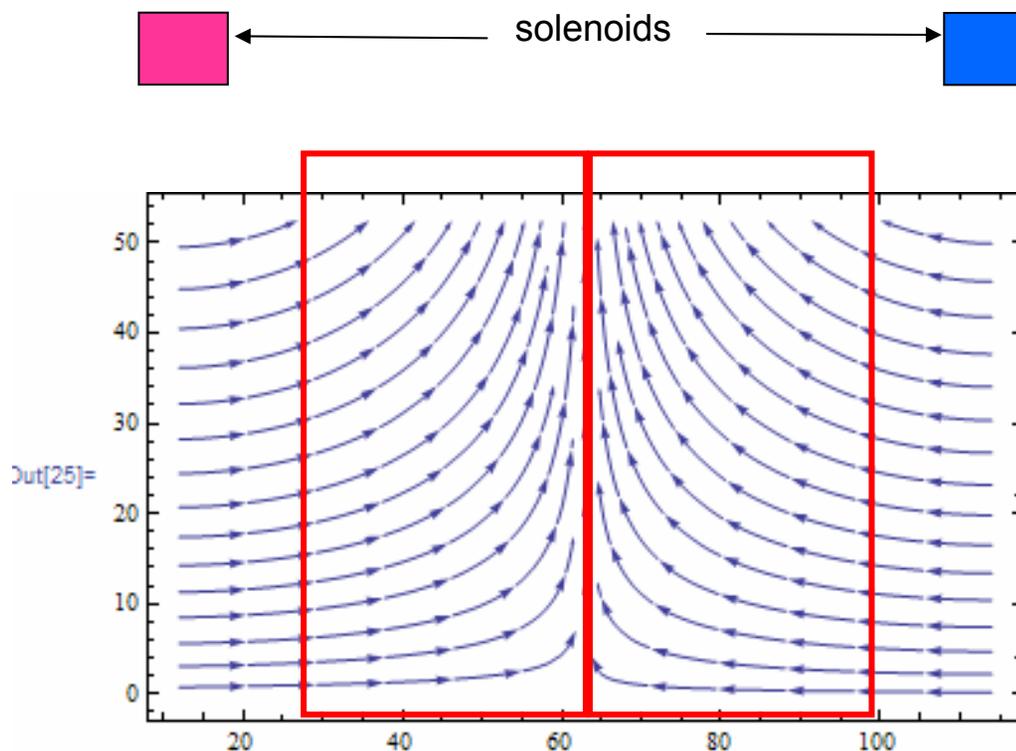
Channel parameters:

200 MHz pillbox RF 2x36cm,  $E_{\text{max}}=16\text{MV/m}$

Solenoids:  $L=24\text{cm}$ ,  $R_{\text{in}}=60\text{cm}$ ,  $R_{\text{out}}=92\text{cm}$ ,

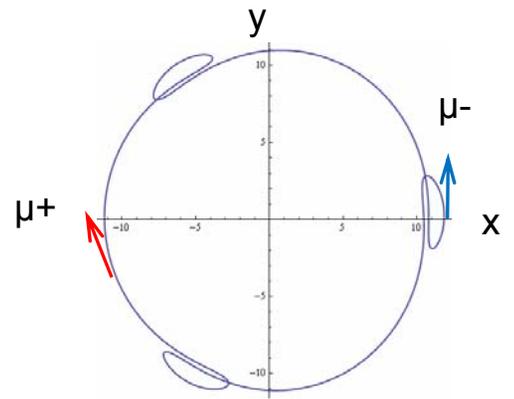
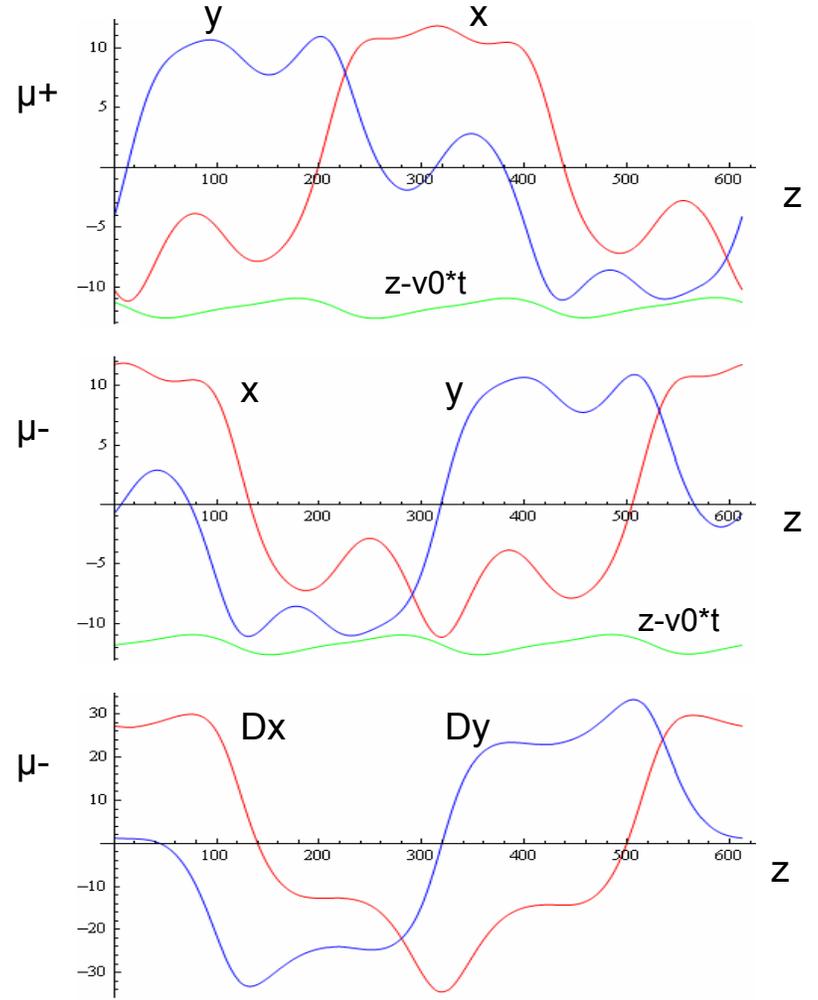
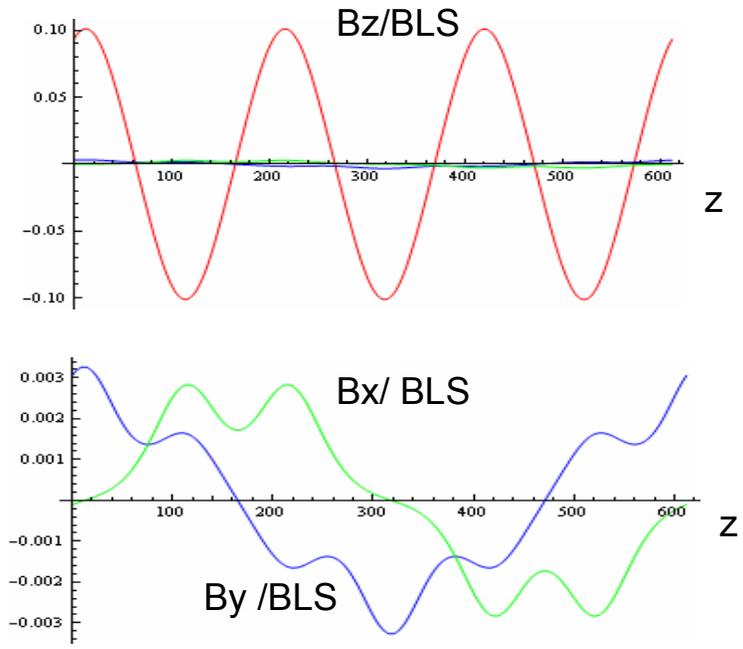
Absorbers: LH2, total width (on-axis)  
6x15cm,

Total length of 6-cell period 6.12m

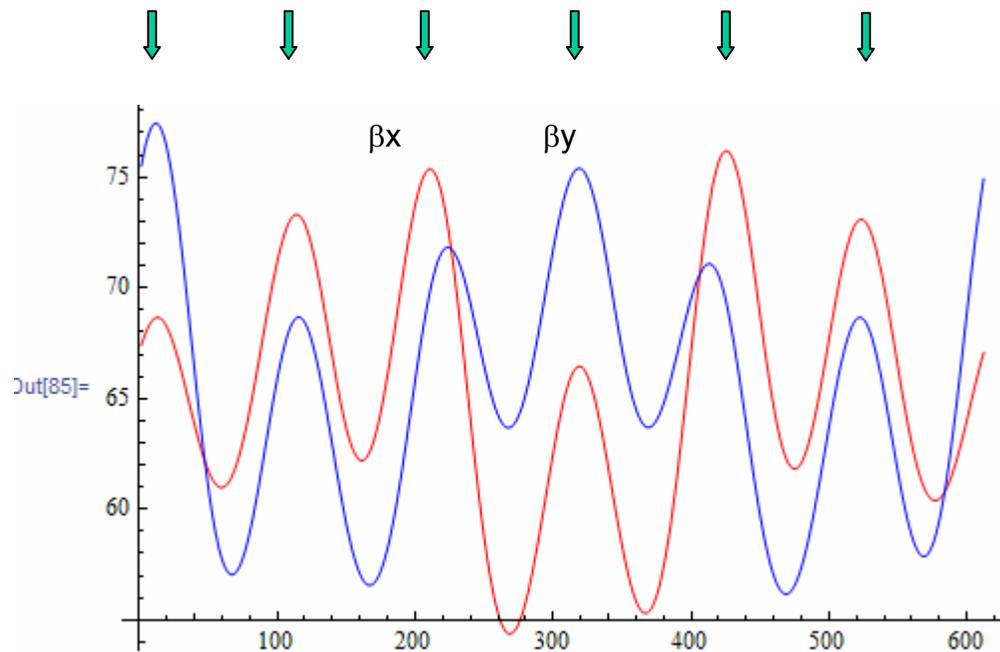


“Poor man” magnetic insulation for  $N_{\text{cavities}} \leq 2$   
Length of 1 cavity is not enough for cooling rate

20mrad pitch angle, BLS=25.2 for  $p=200\text{MeV}/c$



Helical FOFO snake – good for cooling both  $\mu^+$  and  $\mu^-$ !



at the absorber locations

$\langle \beta \rangle \sim 70\text{cm}$

- compare with MICE's

$\langle \beta \rangle \sim 45\text{cm}$

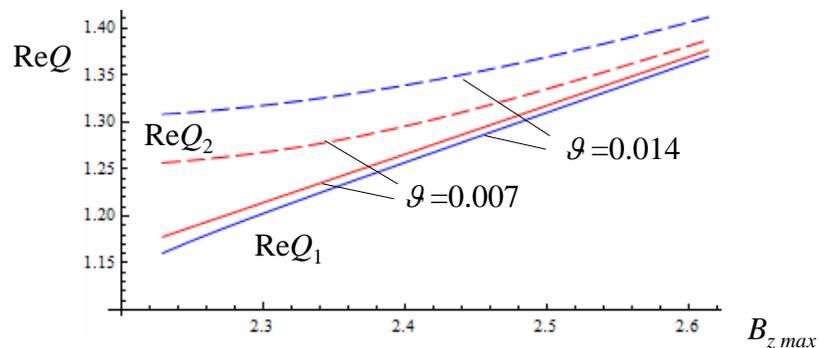
**The best results with 7mrad pitch angle, no absorber wedge angle:**

mode	I	II	III
tune	1.239+0.012i	1.279+0.007i	0.181+0.002i
$\varepsilon_{\text{eq}}$ (mm)	3.2	4.5	6.9

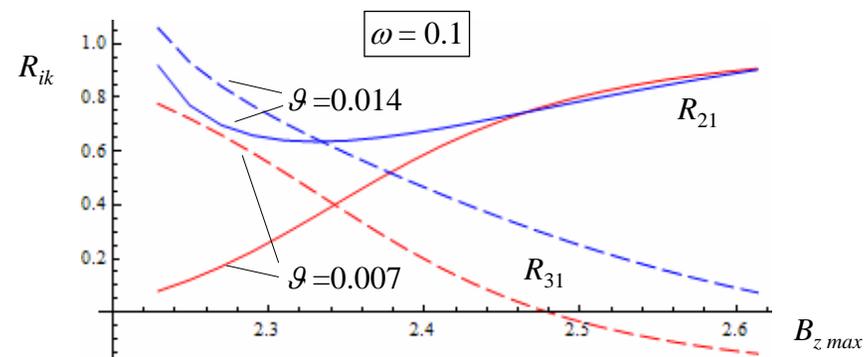
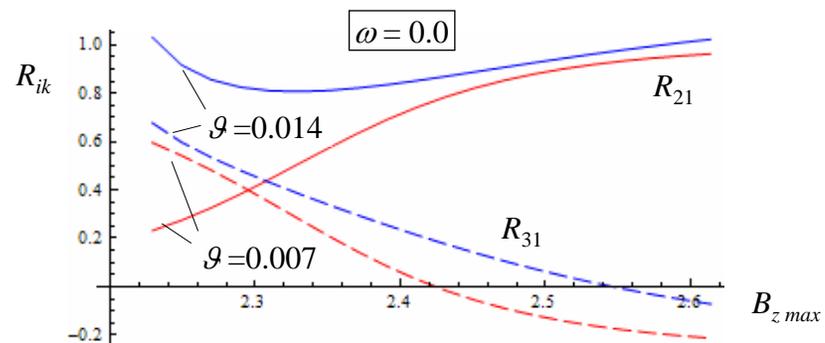
$\text{Im}Q=0.007 \Rightarrow$  cooling rate  $d \log \varepsilon / dz = 2 \times 2\pi / L \text{Im}Q = 1/70\text{m}$

There is difficulty in equalization of damping rates of the transverse modes

$B_{\text{max}}=2.3\text{T} \Rightarrow j=58\text{A/mm}^2, I_{\text{tot}}=4.4\text{MA/solenoid}$



Transverse tunes vs  $B_{z \max}$  for pitch angle  $\vartheta = 7\text{mrad}$  (red) and  $\vartheta = 14\text{mrad}$  (blue).



Damping rate ratios  $R_{ik} = \text{Im}Q_i / \text{Im}Q_k$  for plane absorbers (top) and with wedge angle  $\omega = 0.1$  (bottom)

### Tracking simulations:

- “True” action variables of 1771 particles evenly distributed in tetrahedron

$$(J_I + J_{II})/2.6 + J_{III}/4 < 1 \text{ (cm)}$$

- Phases chosen at random
- No decay nor stochastic processes
- Correlation between energy and oscillation amplitudes introduced into initial conditions to avoid immediate longitudinal blowup

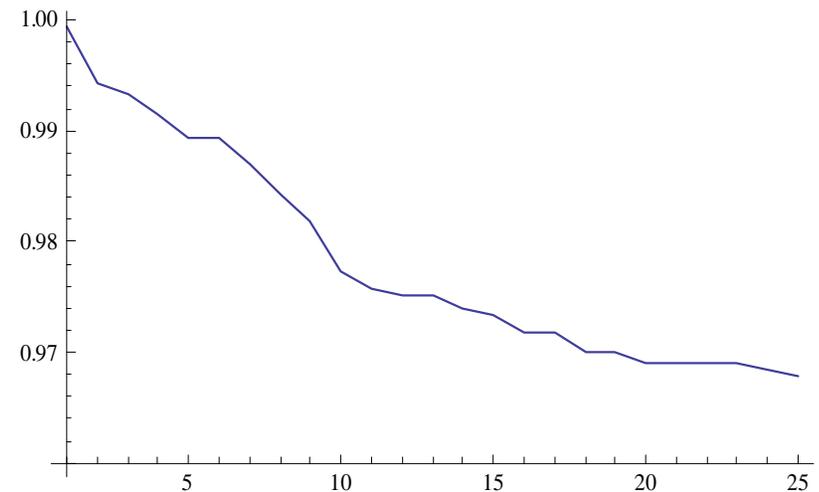
$$\delta_p = \delta_{p0} + 0.042J_1 + 0.013J_2 + 0.039J_3$$

*Note:*

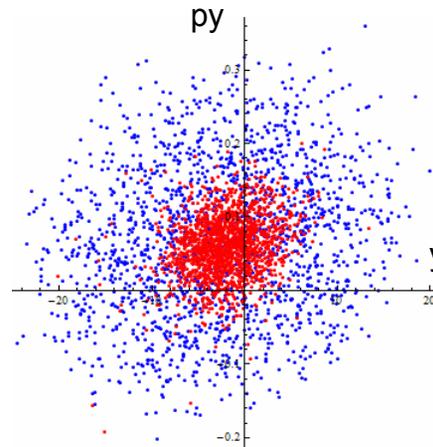
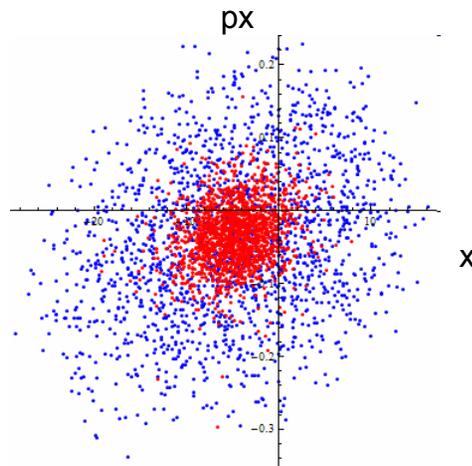
Courant-Snyder invariant = 2J,

to compare with normalized emittance multiply by  $\beta\gamma \approx 2$ :

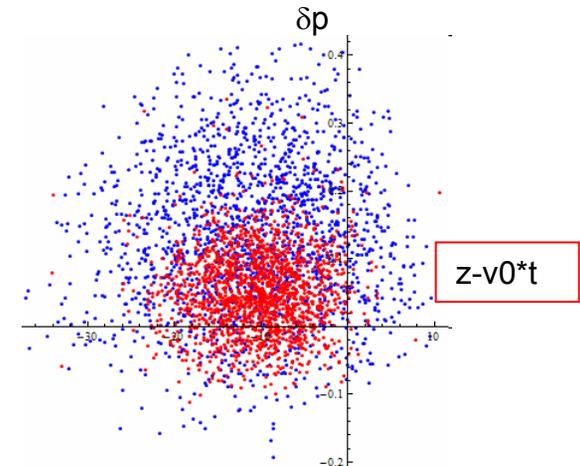
$\beta\gamma \times \text{CSImax} \approx 10\text{cm}$  or  $2.2\sigma$  for  $\varepsilon N = 2\text{cm}$



Survival after 25 periods (153m) 97%



blue - initial, red - final



“Emittances” (cm)

6D

Trans. average

Longitudinal

initial

10.3

1.99

3.75

final

0.07

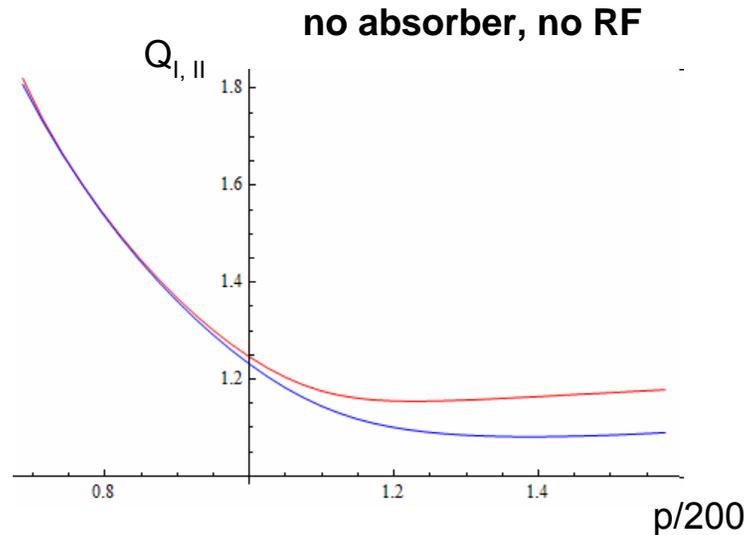
0.29

1.46

$$\delta_p = \frac{\gamma - \gamma_0}{\beta_0^2 \gamma_0},$$

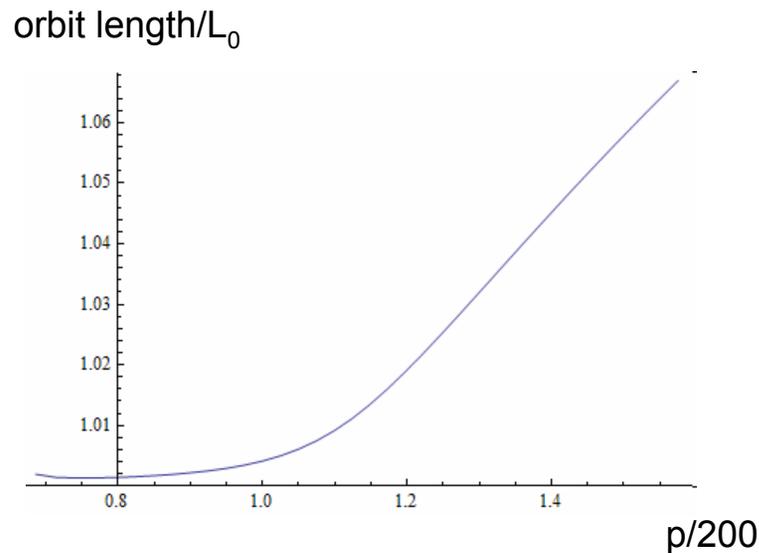
$$\frac{p}{p_0} = 1 + \delta_p - \frac{\delta_p^2}{2\gamma_0^2} + \dots$$

**Why momentum acceptance is so large (>60%)  
in the resonance case?**



**Nice surprise:**

Large 2nd order chromaticity due to nonlinear field components keeps both tunes from crossing the integer !



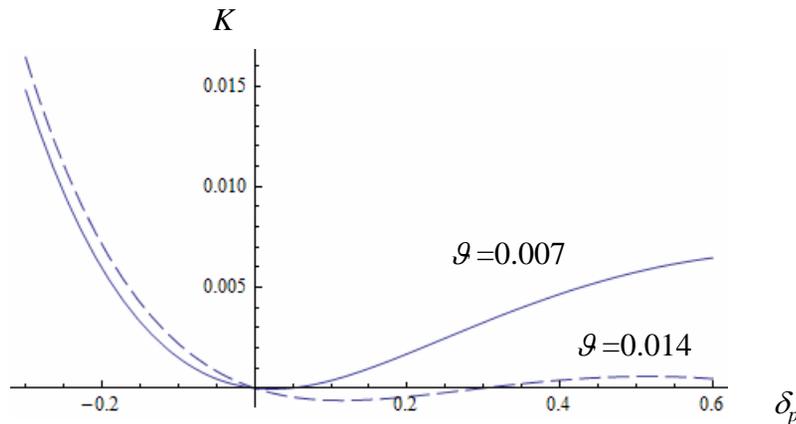
Momentum compaction factor:

$$\alpha_p \approx 0.1 < 1/\gamma_0^2 \approx 0.22$$

- in contrast to classical HCC with homogeneous absorber where

$$\alpha_p > 2/\gamma_0^2$$

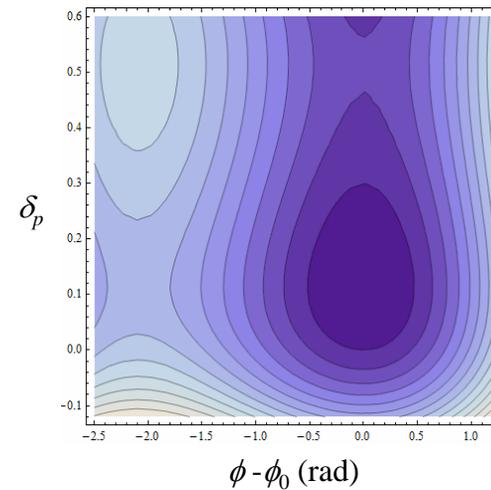
to ensure longitudinal damping.



Kinetic energy in quasi-static approximation

$$K(\delta_p) = \int_{-\delta_p}^{\delta_p} \left[ 1 - \lambda(\xi) \frac{\beta_0}{\beta(\xi)} \right] d\xi$$

where  $\lambda = \text{orbit length}/L_0$



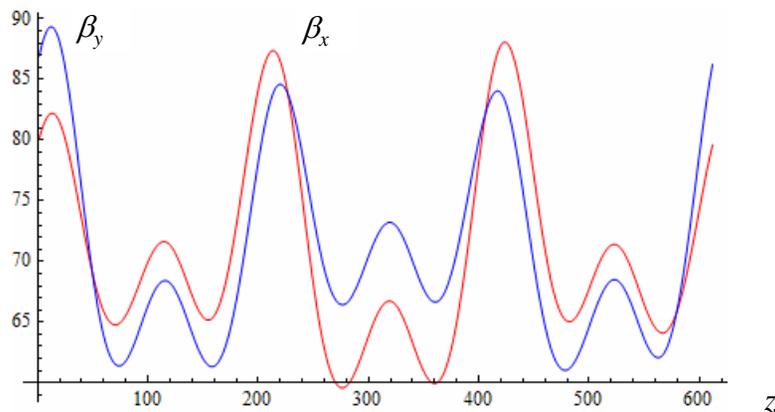
Contour plot of the longitudinal Hamiltonian for  $\mathcal{G} = 0.014$  - kinematic nonlinearity limits momentum acceptance, not insufficient RF bucket depth.

This makes problematic the achieving higher longitudinal damping rate with larger  $\mathcal{G}$ .

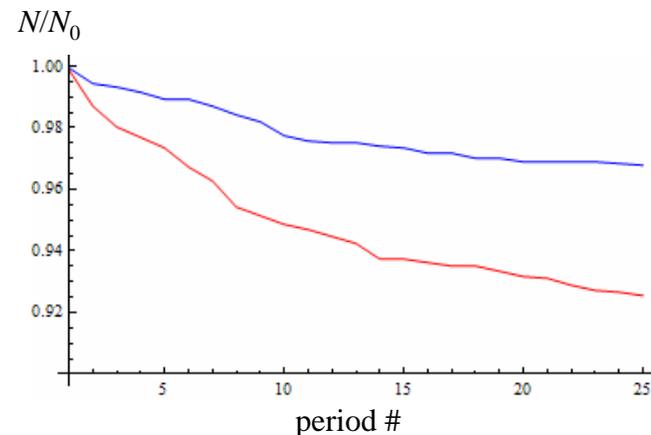
R. Palmer proposed to add  $\sim$  constant solenoidal field to better mix the transverse modes and equalize their damping rates. This can be achieved by powering e.g. negative solenoids with slightly lower current. With just 1.6% difference in currents:

mode	I	II	III
tune	1.211+0.0100i	1.301+0.0108i	0.196+0.0003i
$\varepsilon_{eq}$ (mm)	3.8	3.2	36.5

(fast transverse cooling w/o longitudinal blowup was the intent):



Unfortunately, there is no appreciable gain in equilibrium emittance due to large  $\beta$ -wave excited with unequal solenoid field



Transmission also suffers (red vs. blue)

The proposed “helical FOFO snake:

- provides appreciable longitudinal cooling
- cools  $\mu^+$  and  $\mu^-$  simultaneously
- has sufficiently large acceptance / equilibrium emittance ratio  $\sim 20$  (transmission with account of stochastic effects is yet to be simulated)
- 2nd stage with reduced overall dimensions (using 1-cell RF) will achieve equilibrium emittances to  $\sim 3\text{mm}$  for 200MHz channel.

Works is underway to:

- simulate the channel with “standard” codes like ICOOL and G4BL
- “taper” the channel to optimize cooling rate and transmission
- to combine it with D. Neuffer’s buncher/rotator (see his report at this conference)