

A METHOD FOR CALCULATING NEAR-OPTIMUM ION-EXTRACTOR PROFILES*

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Abstract

A process and program have been developed to interactively calculate the near-optimum electrode profiles for high-perveance ion-extraction systems. A Mathcad™ program determines electrode profiles for high-current (high-perveance) high-quality beams. The program input starts with key parameters: plasma density, estimated mix of ions, extraction voltage, total current, plus desired output beam size and divergence. The computations simulate a spherically convergent extraction system that simultaneously minimizes the aberrations from the exit aperture while directly compensating for both the exit aperture de-focusing lens, and internal space charge in the beam. The program outputs cylindrical (r,z) coordinates of the emission and extractor electrodes, plus displays the beam perveance and output beam size and divergence. This technique was used successfully in multiple projects over the past 25 years. Electrode shapes used in past hardware tests are examined with the successive over relaxation (SOR) code PBGUNS in an accompanying paper.

Background And Need

The ability to extract a high-quality (low-emittance) high-current beam of ions or electrons from an emitting surface or a plasma is critical to nearly all devices that utilize energetic charged-particle beams. Included in this category are all electron-beam RF power sources and all particle accelerators.

There is often a need for higher currents and higher beam quality. This improved performance is most often noted by a combination of increased power, smaller size, improved efficiency, lower activation, simpler cooling, and/or simplified setup and operation. Achieving and maintaining high quality in any beam is an on-going challenge, one that becomes much more difficult as the beam current and total power are increased.

PROCESS

The extractor geometry to be discussed is that of the ‘diode’ extractor shown in Figure 1. Goal of the process is to simulate the ‘perfect’ inward flow of ions (or electrons) between two concentric spheres. This concept was investigated by Langmuir & Blodgett [1] and discussed by Pierce [2]. A very similar process may be used for ‘triode’ or other extractor geometries. But the author has emphasized diode extractors only because they have demonstrated superior performance in the projects he has pursued.

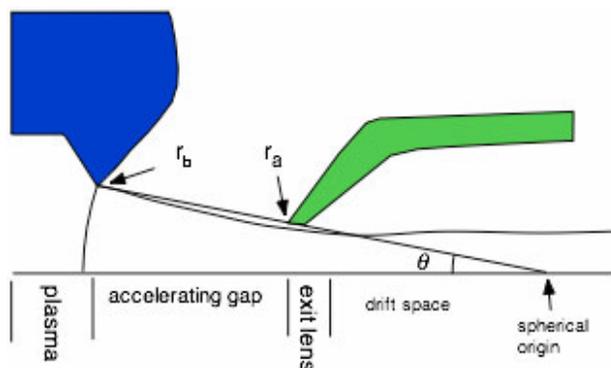


Figure 1: Diode Extractor geometry used in this optimization process.

An extractor geometry made with complete spheres is not practical, but it is possible to approximate the ideal field of complete concentric spheres, but only for a finite solid angle. The technique used in this process entails first determining the exact potential solution for a ‘perfect’ spherical geometry. Then, a few selected Legendre functions are combined to accurately approximate this ideal potential distribution along the outer edge of the particle beam. These functions are selected to provide both a good mathematical fit, and to ensure that the final solution is physically realizable. Then the resulting sum of functions is used to define the electrode shapes needed in the region outside the particle beam. When converged, this process provides electrode shapes that create electric fields inside the extracted beam that very closely approximate the ideal spherical geometry, and facilitate a nearly perfectly laminar flow of charged particles. This technique has proven successful in specifying the “spherical Pierce” geometry that yields a low-aberration extraction system.

The perfect radial potential distribution for a space-charge-limited flow of charged particles filling the region between two concentric spheres may be determined by solving the one-dimensional Poisson equation in spherical coordinates:

$$\Phi''_r = \frac{-2}{r} \Phi'_r - \frac{r_b^2 J_0}{\epsilon_0} \sqrt{\frac{m}{2e}} \frac{1}{(V - \Phi_r)^{0.5}} \frac{1}{r^2} \quad (1)$$

Langmuir and Blodgett [1] are credited with showing that the radial potential solution to this DE may be nicely

approximated by:
$$V(r) = V_a \left(\frac{-\alpha}{-\alpha_a} \right)^{4/3} \quad (2)$$

$$\text{with } \gamma = \ln\left(\frac{r}{r_b}\right) \tag{3}$$

where

$$\alpha = \gamma - 0.3\gamma^2 + 0.075\gamma^3 - 0.001432\gamma^4 + 0.002161\gamma^5 - 0.0002679\gamma^6 + \dots \tag{4}$$

The most-serious imperfection in this process is the extractor aperture through which the charged particles exit the two-electrode system. The exit aperture creates a divergent lens, and internal space charge causes further divergence in the exiting beam. Fortunately the convergent effect of the initial two-electrode extractor can effectively compensate for the first-order effects of both the exit lens and space charge. There remains a small aberration from the edge effects that fortunately normally impacts only the outer 1—2% of the beam.

The Paraxial exit lens effect is given by:

$$\frac{1}{f} = \frac{\Delta r'}{r} = \frac{\Delta E}{4V_0} \quad \text{giving:} \quad \Delta r' = \frac{r\Delta E}{4V_0} \tag{5}$$

There is also a first-order space-charge term, given by:

$$\Delta r' = \frac{\Delta z C P_e}{r} \tag{6}$$

where Δz is the distance over which the full space charge is effective, P_e is the beam perveance (in μPervs), and C is a constant equal to 15.2 mrad/ μPerv . When combined, we have a total exit divergence effect:

$$\Delta r' = \frac{r\Delta E}{4V_0} + 0.0152 \frac{\Delta z}{r} P_e \tag{7}$$

Design Process Details

During program execution, only six parameters need be input: 1) total beam current, 2) effective ion mass, 3) extraction voltage, 4) current density at emission aperture, 5) space-charge neutralization length & 6) convergence half-angle of the diode extractor. The program first computes the beam perveance, resultant beam divergence (after exit lens) and peak field within the extractor gap. The operator then may make corrections to ensure that output beam divergence and peak fields are acceptable. All other calculations are automatic, with the program output including graphic and tabular representations of the emission and extractor electrode profiles.

The potential distribution along the beam edge is calculated by eqn 2 above and a number of points are fitted with a summation of four Legendre Functions. The output of this process establishes the values for the coefficients for these functions. We have shown that the following four Legendre functions (Eqn 8) provide a good fit to the potential, create physically realizable profiles, and avoid instabilities.

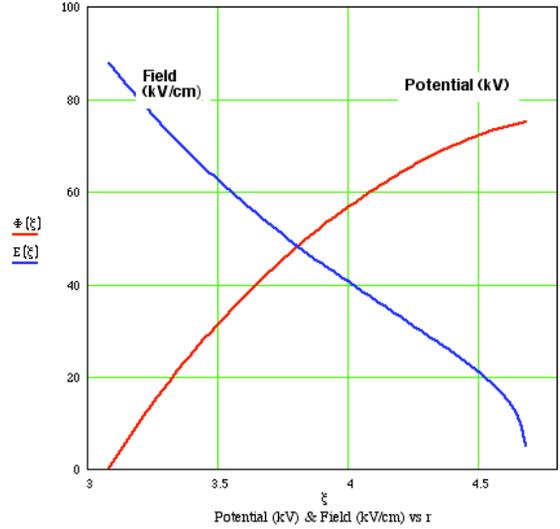


Figure 2: Depiction of the desired potential and electric field along the beam edge and measured from the spherical origin shown in Fig. 1.

$$V(r, \theta) = a_0 + a_1 \cos \theta + \frac{b_2 (3 \cos^2 \theta - 1)}{r^2} + \frac{b_3 (5 \cos^3 \theta - 3 \cos \theta)}{r^3} \tag{8}$$

A non-linear equation solver gives values for the four coefficients a_0, a_1, b_2 & b_3 . The MINERR function in Mathcad™ works well for this, but previous versions have used BASIC code and Mathematica™, with similar success.

The remaining step is the determination of the equipotentials for the two electrodes (which we call the emission and extraction electrodes).

A sample of calculated electrode shapes is shown in figure 3, fabricated as figures of rotation (fig 4).

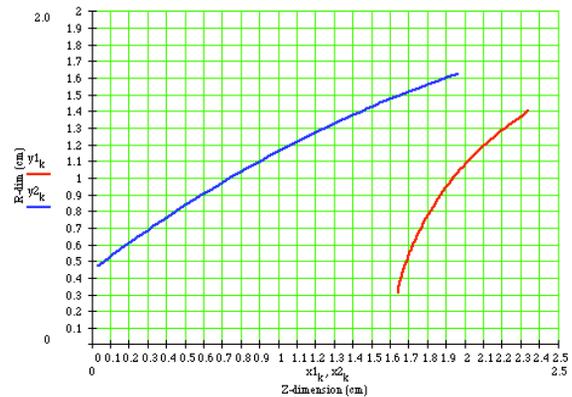


Figure 3: Calculated profiles for the surfaces of the emission and extractor electrodes.

Represented in 3-D, the emission/extractor geometry might become:

